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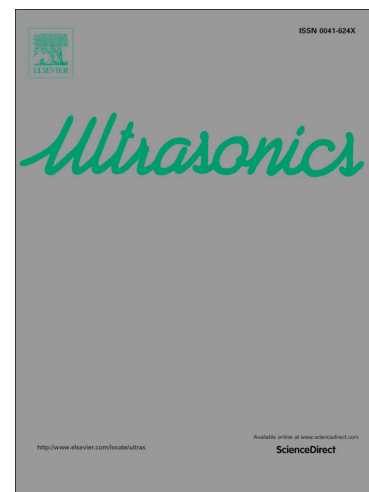
Ultrasound Multifrequency Strategy To Estimate the Lung Surface Roughness, *In Silico* and *In Vitro* Results

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Ultrasound Multifrequency Strategy To Estimate the Lung Surface Roughness, *In Silico* and *In Vitro* Results

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Highlights

- Proposing a multifrequency approach to estimate the lung surface roughness.
- Proposing an experimental model to mimic different roughness at the lung surface.
- Evaluating consistency between *in vitro* and *in silico* experiments.

Lung ultrasound (LUS) is an important imaging modality to assess the state of the lung surface. Nevertheless, LUS is limited to the visual evaluation of imaging artifacts, especially the vertical ones. These artifacts are observed in pathologies characterized by a reduction of dimensions of air-spaces (alveoli). In contrast, there exist pathologies, such as chronic obstructive pulmonary disease (COPD), in which an enlargement of air-spaces can occur, which causes the lung surface to behave essentially as a perfect reflector, thus not allowing ultrasound penetration. This characteristic high reflectivity could be exploited to characterize the lung surface. Specifically, air-spaces of different sizes could cause the lung surface to have a different roughness, whose estimation could provide a way to assess the state of the lung surface. In this study, we present a quantitative multifrequency approach aiming at estimating the lung surface's roughness by measuring image intensity variations along the lung surface as a function of frequency. This approach was tested both *in silico* and *in vitro*,

and it showed promising results. For the *in vitro* experiments, radiofrequency (RF) data were acquired from a novel experimental model. The results showed consistency between *in silico* and *in vitro* experiments.

Keywords: In silico, in vitro, multifrequency analysis, Quantitative Lung Ultrasound (LUS).

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1 I. INTRODUCTION

2 Lung ultrasound (LUS) is nowadays widely adopted in clinical practice to assess the state of the
3 lung surface.^{1,2} Specifically, the main characteristics of LUS (i.e., portability, real-time imaging,
4 and non-ionizing radiations) render it particularly suitable for patients' monitoring.^{3,4}

5 Nevertheless, being LUS mainly based on the visual evaluation of imaging artifacts, it remains
6 subjective and qualitative.^{1,2,4,5} To improve LUS specificity, researchers have recently started to
7 develop quantitative LUS approaches aiming at estimating the state of the lung surface.⁶⁻¹³

8 However, these approaches can be used to assess the state of lung surface only in pathologies
9 characterized by a reduction of dimensions of air-spaces (alveoli), and thus an increased
10 permeability of the lung with respect to ultrasound waves.^{2,14} Therefore, these techniques cannot
11 be used for patients affected by lung pathologies characterized by an enlargement of air-spaces.¹⁴

12 Being the third leading cause of death worldwide (causing 3.23 million deaths in 2019),^{15,16}
13 chronic obstructive pulmonary disease (COPD) can be considered the most representative
14 example of lung disease characterized by an enlargement of air spaces. Indeed, the peripheral air-
15 space dimensions in COPD patients are generally above 490 μm , whereas they are mainly
16 between 340 and 440 μm in healthy subjects.¹⁷ Similarly to what happens in a healthy lung, the
17 enlargement of peripheral air-spaces' dimensions causes the lung surface to behave essentially as
18 a perfect reflector, thus not allowing ultrasound penetration.¹⁴ This characteristic high reflectivity
19 could be exploited to characterize the lung surface. Specifically, we hypothesize that air-spaces of
20 different sizes could cause the lung surface to exhibit a different roughness,¹⁸ whose estimation
21 could provide a way to assess the state of the lung surface.¹⁴ To clarify, as the increase of air-
22 spaces' dimensions can occur also peripherally (thus, at the lung surface),^{17,18} this increase is
23 implicitly translated into a variation of lung surface (air) roughness. Fig. 1 shows a pictorial

24 representation of lung surface roughness variation following air-spaces' enlargement. This
25 enlargement of peripheral air-spaces causing the lung surface to have a different roughness can
26 be clearly observed from lung histologies of COPD patients, as shown in ¹⁸.

27 In this study, we present a quantitative multifrequency approach aiming at estimating the lung
28 surface's roughness. In particular, this estimation is performed by measuring image intensity
29 variations along the lung surface as a function of frequency.¹⁴ Specifically, when roughness is
30 introduced at the surface and a wave is transmitted in a direction perpendicular to the surface, it
31 is expected that a part of the wave is back-scattered in directions different from the transmitted
32 one, thus causing a decrease of intensity. This deviation from the transmission direction should
33 depend on the relation between the wavelength (thus, the frequency) in the propagation medium
34 (i.e., soft tissues) and the roughness. Therefore, we aim at understanding whether it is possible to
35 observe intensity variations of the waves back-scattered from a rough surface, map this variation,
36 and, finally, evaluate the possibility to extract information on the roughness dimensions from the
37 obtained map. We set our experiments in specific ranges of roughness and frequency, which
38 correspond to the range of interest of peripheral air-spaces' dimensions¹⁷ and the frequencies
39 typically used in medical ultrasound, respectively. This approach was tested both *in silico*
40 (numerical simulations), and *in vitro*, where radiofrequency (RF) data were acquired from a novel
41 experimental model.

42 The paper is organized as follows. While materials and data acquisition are presented in Sec. II-
43 A, the process to quantify the surface intensity is described in Sec. II-B. The results are then
44 presented in Sec. III, followed by discussion and conclusions in Sec. IV.

45 II. MATERIALS AND METHODS

46 A. Materials and Data Acquisition

47 *In Silico*

48 The computational 2D (two-dimensional) domain simulated with the k-wave¹⁹ MATLAB
49 toolbox consisted of 2000×4000 pixels along the lateral dimension and depth, respectively. As
50 the numerical grid-size (square pixels) was equal to 10 μm, the physical size of computational
51 domain resulted in 20×40 mm. Specifically, a homogeneous muscle layer was simulated in the
52 first 20 mm of depth, whereas a lung surface with 9 different levels of roughness was added at
53 20 mm.¹⁴ The roughness was simulated by introducing, at the lung surface, semi-circular
54 scatterers (representing alveoli) having a diameter ranging from 200 to 600 μm, with a 50-μm
55 step-size. This roughness was introduced only in the central part of the domain (between -3 and
56 3 mm in the lateral dimension), whereas the remaining parts of the interface (from -10 to -3 mm
57 and from 3 to 10 mm) were kept smooth (see Fig. 2). This was done to clearly visualize intensity
58 variations between the smooth and the rough parts of the mimicked lung surface. Moreover, the
59 smooth areas served as reference points when the reconstructed images were normalized with
60 respect to their maximum (see Sec. II-B).

61 Steel was used to mimic the lung surface, as it can simulate a highly reflective acoustic interface.
62 Specifically, the reflection coefficient of a steel/muscle interface ($R \cong 0.93$) is comparable with
63 the reflection coefficient of an air/muscle interface ($R \cong 0.99$). It is important to highlight how a
64 material able to form a highly reflective acoustic interface is needed as we aim at analyzing the
65 phenomena occurring at the interface and not inside that material. Moreover, steel allowed us to
66 fabricate lung-mimicking phantoms having controllable size at micrometric scale, thus providing
67 us the possibility for a consistent comparison with the *in silico* experiments. Table I shows the
68 simulated acoustic properties for muscle and steel.

69 For each domain, data were acquired with a plane wave imaging strategy by transmitting a 4-μs-
70 time-length pulse (bandwidth equal to 0.5 MHz at -6 dB), and center frequencies from 3 to 10

71 MHz, with a 1-MHz step size (8 images per domain). Table II shows the wavelength values for
72 muscle (λ_0) and steel (λ_{steel}), as well as their ratio with the numerical grid-size. In transmission
73 phase, the entire array (composed by 64 elements) was excited. The kerf and pitch were 45 and
74 245 μm , respectively; however, for the domain approximation, the actual simulated values were
75 40 and 240 μm , respectively. The array was placed at 150 μm of depth from the beginning of the
76 computational domain, and centered with respect to the lateral dimension (i.e., laterally
77 extending from approximately -7.7 mm to 7.7 mm). To reconstruct each image, a sub-array of 16
78 elements was linearly shifted along the entire array in reception, thus forming images composed
79 by 49 lines along the lateral dimension. No focus was applied both in transmission and reception
80 phases. The time sampling interval dt equals to 316 ps, resulting in a sampling frequency of $1/dt$
81 $\cong 3.1645$ GHz. No time gain compensation (TGC) was applied, and a speed of sound of 1580
82 m/s was assumed for the time-space conversion (along depth).

83 The choice of a plane wave imaging strategy was made to save computational time (thus,
84 performing a higher amount of simulations). Similarly, plane wave transmissions were used *in*
85 *vitro*.

86 ***In Vitro***

87 The ULA-OP programmable platform²⁰ and an LA533 (Esaote, Florence, Italy) linear array
88 probe (having pitch and element size along the lateral dimension equal to 245 and 220 μm ,
89 respectively²¹) were exploited to acquire RF data with different center frequencies. Specifically,
90 the data were acquired by transmitting pulses having the same bandwidth and center frequency
91 of the *in silico* experiments. The utilized probe has a -6 dB bandwidth from 3.8 to 12 MHz and a
92 -12 dB bandwidth from 3.2 to 13.2 MHz.¹¹ The maximum of the transducer transfer function is
93 at 8 MHz. A 50 MHz sampling frequency was used ($dt = 20$ ns). A sub-aperture of 64 elements
94 was employed in transmission and reception, and the images were reconstructed with dynamic

95 beamforming (dynamic focus in reception). Each final image, which was reconstructed by
96 linearly shifting this sub-aperture over the entire array (192 elements), consists of 129 lines. To
97 avoid saturation phenomena, the driving signal amplitude was maintained to 10% of the
98 maximum amplitude allowed by the ULA-OP system.¹² To clarify, the driving signal is the
99 electrical signal utilized to excite each element of the transmit aperture (the maximum output
100 voltage is 24 Vpp²⁰). No time gain compensation (TGC) was applied, and a speed of sound of
101 1480 m/s was assumed for the time-space conversion (along depth).

102 To *in vitro* mimic the same levels of roughness utilized *in silico*, a phantom consisting of austenitic
103 stainless steel AISI (American Iron and Steel Institute) 316L (Euronorm number = 1.4404) was
104 produced by a Concept Laser Mlab (General Electric Additive, Boston, US), i.e., a 3D metal
105 LPBF (laser powder bed fusion) printer (maximum power = 100 Watts, and laser spot size = 45
106 μm). Specifically, the phantom consisted of 9 stripes along its length made with varying
107 roughness. Each rough stripe was composed of semi-cylinders arranged consecutively along the
108 phantom width. The 9 rough stripes were separated by 8 smooth stripes (no roughness), each
109 having the same length (1 cm). Therefore, as shown in Fig. 3, the total phantom length was 17
110 cm (9 cm + 8 cm), whereas its width was 5 cm (to allow the probe to laterally cover the entire
111 area).

112 The steel phantom was immersed in a water tank and positioned on a steel plate (Fig. 4), which
113 was used to align the phantom consistently with the probe displacement, guided by an automatic
114 positioning system (GAMPT, Merseburg, Germany). We have defined the axis parallel to the
115 phantom width as lateral direction, and the axis parallel to the phantom length as elevation
116 direction (Fig. 4, left). The phantom was placed at 2 cm of depth from the probe (Fig. 4, top
117 right).

118 The data were sequentially acquired. Specifically, the probe was placed above the first rough

119 stripe (Fig. 4, top right), and all the data corresponding to that roughness level were acquired
 120 with different center frequencies. Then, the probe was automatically moved along the elevation
 121 direction by 2 cm to acquire the data corresponding to the second roughness level. The process
 122 was repeated until reaching the last rough stripe.

123 **B. Quantification of Surface Intensity**

124 *In Silico*

125 To estimate the surface intensity, a two-step procedure was applied to the reconstructed images.
 126 In step 1, to extract the envelope, the Hilbert transform was applied. Each reconstructed image
 127 was then normalized with respect to its maximum value. In step 2, we displayed the images in
 128 logarithmic scale with a 35-dB dynamic range, and defined a region of interest (ROI) in which
 129 we computed the total intensity (I_{TOT}).^{9,11,12} This ROI was defined as the area where the
 130 roughness was introduced, i.e., between -3 and 3 mm in the lateral dimension, and extending in
 131 depth from 18.5 mm to 21.5 mm (see Fig. 5, first row). The depth range was set by considering
 132 the spatial length of the transmitted pulse, i.e., $4 \mu s \times 1580 m/s / 2 \cong 3 mm$, and the surface
 133 depth (20 mm). Only values above -35 dB in the ROI (empirical threshold) were used to
 134 compute I_{TOT} .^{9,11}

135 To evaluate the surface intensity as a function of frequency, we exploited a scaled version of
 136 I_{TOT} (normalized I_{TOT}). Specifically, for each roughness level, we normalized the eight I_{TOT}
 137 values (obtained with frequencies varying from 3 to 10 MHz) with respect to their maximum
 138 (see Fig. 5, bottom, red line).^{9,11,12}

139 *In Vitro*

140 The surface intensity was estimated *in vitro* with a procedure consistent with what was done *in*
 141 *silico*. The main differences are associated with step 1 and with the ROI definition. Specifically, as

142 first operation, to be consistent with the *in silico* experiments, we evaluated only the 49 central
 143 lines (from approximately -6 to 6 mm along lateral dimension) of the reconstructed images.
 144 Considering only the central lines allowed us also to prevent undesired contributions coming
 145 from the edges of the steel phantom. Then, we applied a sixth-order bandpass Butterworth filter
 146 having a 1-MHz bandwidth and centered at the different center frequencies.¹² We successively
 147 applied the Hilbert transform to each filtered image, thus extracting the envelope.¹² To
 148 normalize the images consistently with what was done *in silico*, we acquired data from a smooth
 149 steel surface by using the same acquisition settings (e.g., from 3 to 10 MHz of center frequency
 150 and 0.5 MHz of bandwidth). Then, after having applied the same processing steps used for the
 151 rough surface (above mentioned), we extracted the 8 maximum values (one for each center
 152 frequency) from the smooth surface data. Finally, the rough surface images were normalized
 153 with respect to these maxima, and displayed in logarithmic scale with a 35-dynamic range. After
 154 these processing operations, the surface intensity was computed by means of the I_{TOT}
 155 parameter.^{9,11} Specifically, this parameter was computed in a ROI extending over the entire
 156 lateral dimension (from approximately -6 to 6 mm) and depths from 18.5 mm to 21.5 mm (see
 157 Fig. 5, second row). The empirical threshold was set to -35 dB as done *in silico*.

158 The normalized I_{TOT} (see Fig. 5, bottom, blue line) was then computed as done *in silico*.

159 III. RESULTS

160 Fig. 6 shows the normalized I_{TOT} for the 9 different scatterers' diameter as a function of
 161 frequency for *in silico* (**normalized I_{TOT}^{IS}** ; see Fig. 6, red graphs) and *in vitro* (**normalized I_{TOT}^{IV}** ;
 162 see Fig. 6, blue graphs) experiments. It is observable how a strong agreement between
 163 **normalized I_{TOT}^{IS}** and **normalized I_{TOT}^{IV}** exists, and becomes stronger when the diameter
 164 increases.

165 To more precisely assess the consistency between *in silico* and *in vitro* results, Fig. 7 and Fig. 8
 166 show the normalized I_{TOT} values for *in silico* (Fig. 7) and *in vitro* (Fig. 8) experiments as a function
 167 of λ and D. To consistently compare Fig. 7 and Fig. 8, we set in the *x-axis* the wavelength of
 168 soft-tissues (λ_{ST}), computed considering a speed of sound of 1530 m/s, which is between the
 169 speed of sound in water (1480 m/s; *in vitro* experiments) and in muscle (1580 m/s; *in silico*
 170 experiments). It is observable how the normalized I_{TOT} values seem to be consistent, and how it
 171 is possible to draw a linear model from both figures. Specifically, to model the relation between
 172 λ_{ST} and D, a linear regression (LR) model fitting specific normalized I_{TOT} values > -3 dB (red
 173 circles in Figs. 7 and 8) was obtained by means of the *fitlm* MATLAB function for both *in silico*
 174 (Fig. 7) and *in vitro* (Fig. 8) experiments. To further clarify, the -3 -dB arbitrary threshold was
 175 applied to the normalized I_{TOT} values depicted in Fig. 6. Then, only a subgroup of values greater
 176 than -3 dB (red circles in Figs. 7 and 8) were used to fit the LR models depicted in Figs. 7 and 8
 177 (purple line). Specifically, for each D (from 200 to 600 μm), we considered the first peak of
 178 normalized I_{TOT} (values greater than -3 dB) starting from the smallest λ_{ST} (highest frequency).
 179 However, for the *in vitro* experiments, the fit was performed by considering the peaks observed
 180 from 400 to 600 μm , as this approach seems not to work on these data for lower values of D
 181 (below 400 μm). The linear models obtained from *in silico* and *in vitro* experiments are $D [\mu\text{m}] = -$
 182 $93.9794 + 2.2142 \times \lambda_{ST} [\mu\text{m}]$ and $D [\mu\text{m}] = -34.0824 + 1.992 \times \lambda_{ST} [\mu\text{m}]$, respectively. By
 183 considering the *in vitro* results (Fig. 8), no peaks (values of normalized $I_{TOT} > -3$ dB) were
 184 observed for higher frequencies (8, 9 and 10 MHz) and, thus, smaller λ_{ST} (191, 170, and 153 μm).
 185 In contrast, *in silico* results showed peaks at those frequencies (Fig. 7). Overall, 34 (9 of which
 186 used for LR) and 24 (5 of which used for LR) peaks were detected *in silico* and *in vitro*,
 187 respectively.

188 Finally, we performed further numerical simulations following the same procedure adopted for

189 steel but simulating air (speed of sound and volumetric mass density equal to 300 m/s and 1.23
190 kg/m³, respectively¹³) instead of steel. The results are shown in Fig. 9, where the obtained linear
191 model is also presented (Fig. 9, bottom). The fit was performed by considering the peaks
192 observed from 400 to 600 μm (as done for the experimental model made by steel), and the
193 obtained linear model is $D [\mu\text{m}] = 18.1648 + 2.6560 \times \lambda_{\text{ST}} [\mu\text{m}]$. The obtained linear model
194 seems to be consistent with the model obtained in Fig. 7 (simulation of steel), with the main
195 difference associated with a positive offset of about 112 μm .

196 IV. DISCUSSION AND CONCLUSIONS

197 LUS is an imaging modality used by clinicians to evaluate the state of the lung surface in real
198 time.^{1,2} However, LUS has a poor specificity as it is mainly based on the visual evaluation of
199 imaging artifacts.^{1,2,4,5}

200 Even though recent attempts to develop LUS quantitative techniques exist,^{6–13} these approaches
201 are applicable only to pathologies characterized by a reduction of air-spaces' dimensions.^{2,14}
202 Indeed, these techniques were developed to characterize a lung having an increased permeability
203 with respect to ultrasound waves.^{2,14} No quantitative approaches have been designed for lung
204 pathologies characterized by an enlargement of air-spaces (impermeable lung).¹⁴

205 For this reason, in this article we have proposed a quantitative multifrequency approach to
206 estimate the lung surface's roughness, which can allow the indirect estimation of the air-spaces'
207 dimensions.¹⁴ This approach was tested both *in silico* and *in vitro* (using a novel experimental
208 model).

209 As shown in Fig. 6, it is clear how by increasing D from 250 to 600 μm the variability between
210 the *in silico* and *in vitro* results tends to decrease, especially for 550 and 600 μm . This can be
211 explained by the ability of the steel model to consistently mimic the numerically simulated 2D
212 domain when D is larger. Specifically, as shown in Fig. 3 and Fig. 4 (bottom right), micrometric

213 imperfections in the printing process more strongly affect a roughness characterized by smaller
214 values of D . Therefore, these imperfections could lead to stronger inconsistencies between *in*
215 *silico* and *in vitro* results when D is smaller.

216 However, as the peripheral air-space dimensions are generally above $340\ \mu\text{m}$,¹⁷ the experimental
217 model could be considered a reliable tool to mimic real air-spaces dimensions, which generate
218 different roughness levels at the lung surface. Specifically, as above mentioned, the *in vitro* results
219 are strongly consistent with *in silico* results for values of D above $500\ \mu\text{m}$, which correspond to
220 the peripheral air-spaces dimensions of COPD patients.¹⁷ This shows how this novel
221 experimental model could be exploited to reliably mimic levels of roughness observable at the
222 lung surface of patients affected by pathologies characterized by increased air-spaces'
223 dimensions. Moreover, we have shown how a simple linear model could be utilized to assess
224 surface roughness by measuring image intensity variations as a function of frequency.

225 Even though promising results were presented in this study, the presence of specific limitations
226 should be highlighted. The first limitation consists in the presence of shear waves and mode
227 conversions in steel, which were not considered *in silico*, but could potentially have played a role
228 during the *in vitro* experiments. Moreover, contrary to what happens in steel (shear velocity equal
229 to $3100\ \text{m/s}$ ²³), shear waves should not play a significant role in lung tissue. Another limitation
230 could be associated with the periodic and simplified geometry that we analyzed in this study with
231 respect to a more heterogeneous geometry observable in real lungs. The small discrepancies
232 between the imaging strategies utilized *in silico* and *in vitro* could also have an impact on the
233 obtained results, even though both strategies are based on unfocused transmissions. In addition,
234 it is important to highlight how the threshold to normalized I_{TOT} used to fit the linear models
235 was arbitrarily set to $-3\ \text{dB}$. Finally, even though the variation of intensities at the lung surface
236 could be caused by different roughness levels, it could be caused also by other lung

237 abnormalities, such as sub-pleural consolidations. All these aspects should be evaluated when
238 translating this multifrequency approach *in vivo*.

239 In this study, for simplicity, only unfocused (plane wave) transmissions were employed. In future
240 studies, the impact of focused beams will also be investigated. As other future studies, we aim at
241 assessing how the presence of a heterogeneous medium in the first 20 mm of depth can impact
242 on the roughness estimation. Moreover, the impact of the ultrasound beam's angle of incidence
243 on the roughness estimation will be analyzed. After this, we also plan to validate this quantitative
244 multifrequency approach *in vivo* by acquiring and analyzing RF data from COPD patients.

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249 Union nor the granting authority can be held responsible for them.

250 **AUTHOR DECLARATIONS**

251 **Conflict of Interest**

252 The authors declare that they have no known competing financial interests or personal
253 relationships that could have appeared to influence the work reported in this paper.

254 **DATA AVAILABILITY**

255 Data are available at the following link [https://drive.google.com/drive/folders/1TA-](https://drive.google.com/drive/folders/1TA-VoHpShULrIr9Y4EnjRHo9QJ9tYZdp?usp=sharing)
256 [VoHpShULrIr9Y4EnjRHo9QJ9tYZdp?usp=sharing](https://drive.google.com/drive/folders/1TA-VoHpShULrIr9Y4EnjRHo9QJ9tYZdp?usp=sharing).

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326 TABLES AND FIGURES

TABLE I. Simulated acoustic properties for muscle and steel.

Medium	Speed of sound [m/s]	Volumetric mass density [kg/m ³]	Acoustic impedance [MRayl]
Muscle ²²	$c_0=1580$	$\rho_0=1041$	$Z_0= c_0 \times \rho_0 \cong 1.645$
Steel ²³	$c_{\text{steel}}=5940$	$\rho_{\text{steel}}=7860$	$Z_{\text{steel}}= c_{\text{steel}} \times \rho_{\text{steel}} \cong 46.69$

TABLE II. Wavelength values for muscle and steel and their ratio with the numerical grid-size.

	Frequency [MHz]							
	3	4	5	6	7	8	9	10
λ_0 [μm]	527	395	316	263	226	198	176	158
$\lambda_0/\text{grid-size}$	52.7	39.5	31.6	26.3	22.6	19.8	17.6	15.8
λ_{steel} [μm]	1980	1485	1188	990	849	743	660	594
$\lambda_{\text{steel}}/\text{grid-size}$	198	148.5	118.8	99	84.9	74.3	66	59.4

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