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# Perspectives for anti deuteron search in cosmic rays with an helium calorimeter

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**Abstract.** The observation of sub-GeV antideuteron in the cosmic ray flux could be a very strong signature of dark matter annihilation in our galaxy. The goal of the Anti Deuteron Helium Detector (ADHD) project is to study the signatures offered by an high pressure Helium calorimeter for the identification of antideuterons in space. The performance of an high pressure (200 bar) Helium scintillator prototype, capable of fast timing, will be shown.

## 1. Introduction

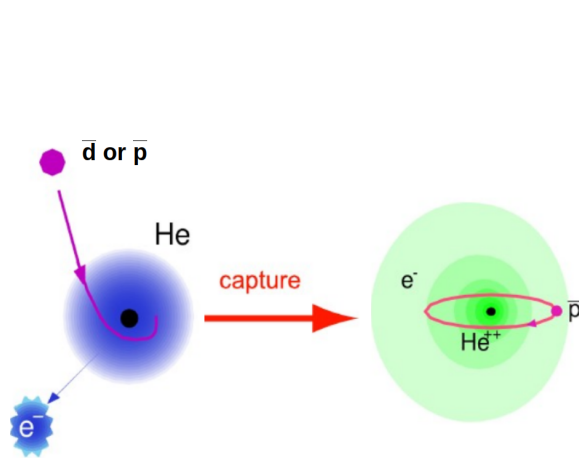
Typical lifetime for stopping antideuterons  $\bar{d}$  (like stopping antiprotons  $\bar{p}$ ) in the matter is of the order of  $\sim$  ps. However, since 1991, the existence of long-living ( $\sim \mu$ s) metastable states for stopping  $\bar{p}$  in helium target was measured [1, 2]. These metastable states in helium have also been measured for other heavy negative particles [3, 4], such as  $\pi^-$  and  $K^-$ . The theoretical description of the effect [5, 6, 7, 8, 9] is predicting that the metastable state lifetimes increase as the reduced mass squared, i.e. a slightly larger delay is expected for  $\bar{d}$  capture in helium with respect to the measured  $\bar{p}$  case. The antiprotonic-helium metastable states are well understood and their existence is already used for other fundamental physics measurements like the antiproton to electron mass ratio [10]. The phenomenology for the formation of metastable states in helium can be summarized following the scheme of Fig. 1.

## 2. Antideuteron identification with the ADHD experiment.

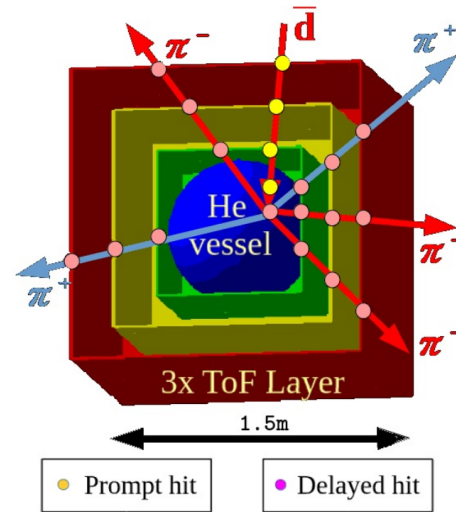
A possible configuration for an helium-based  $\bar{d}$  detector is depicted in Fig. 2. The response of a similar detector to  $\bar{d}$ ,  $\bar{p}$  and to the main cosmic ray components (p, He, C,  $e^-$ ) has been simulated with the Geant4 10.5 package [11]. The inner part is a  $\sim$  20 kg scintillating helium calorimeter (HeCal) where the 400 bar gas is filling a 300L ( $\varnothing = 90$  cm) spherical thermoplastic vessel. Vessel wall thickness of  $\sim$ 3 cm (vessel mass  $\sim$  100kg) ensures a burst pressure larger than 800 bar. A similar vessel is already considered for He transportation in space [12].

Helium gas is a fast UV scintillator, having a light yield similar to other fast plastic/liquid scintillators and capable of  $\sim$  ns timing performances [13]. The HeCal is surrounded by three layers, made by 4 mm thick plastic scintillator bars, providing velocity measurement ( $\beta$ ) by Time of Flight (ToF) and charge measurement ( $Z$ ) by ionization energy loss measurement ( $dE/dX$ ). It is assumed that with current technology such a ToF detector is capable of measuring  $\beta$  with 5% resolution and energy loss with 10% resolution. Considering the energy loss in the ToF and in the vessel, a minimum kinetic energy of  $\sim$  60 MeV/n is necessary for  $\bar{d}$  to reach the He target.





**Figure 1.** Stopping  $\bar{p}$  and  $\bar{n}$  (but also  $\pi^-$  and  $K^-$ ) can be captured by He and trapped in ( $\mu$ s living) metastable states.



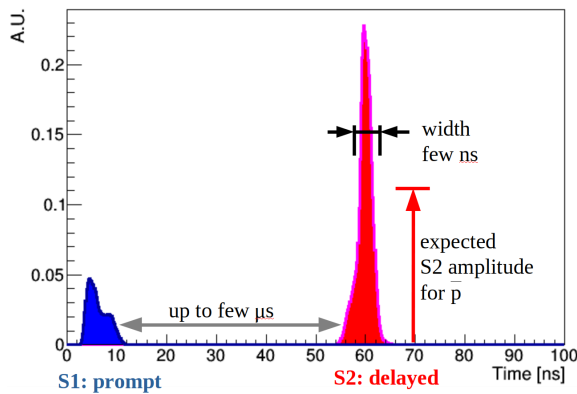
**Figure 2.** The ADHD detector is made by a 400 bar He gas calorimeter (HeCal) surrounded by three scintillator layers (ToF).

On the other hand,  $\bar{d}$  with kinetic energy larger than 140 MeV/n would typically cross the 400 bar He active region without stopping inside.

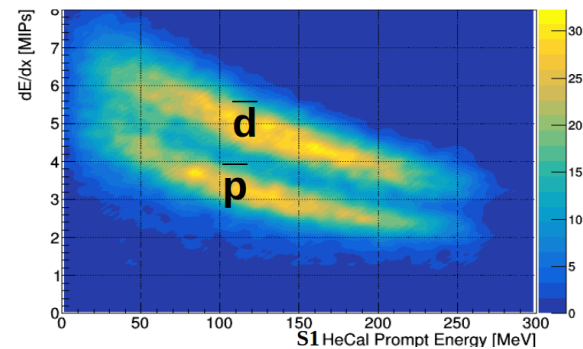
This defines the 60-140 MeV/n energy window explorable by ADHD in this configuration.

Fig. 2 also shows the typical event topology for a stopping  $\bar{d}$  within the He gas. The antiparticle initially produces three prompt hits (yellow) in the ToF and one prompt energy release (S1) in the HeCal; these prompt hits occurs within 10 ns (and are also produced by any other ionizing stopping particle). Then, only for  $\bar{p}$  and  $\bar{d}$ , the antiparticle can be captured in the He metastable states and after a time delay going from several tens of ns to few  $\mu$ s the annihilation occurs (pink delayed hits in Fig. 2). Typical  $\pi^\pm$  multiplicity is  $3.0 \pm 0.2$  for each anti-nucleon annihilation at rest [14], therefore twice the number of delayed out-going tracks is expected for  $\bar{d}$  w.r.t.  $\bar{p}$ . For the same reason, also the delayed signal (S2) in the HeCal for  $\bar{d}$  is expected to have a double amplitude w.r.t. the delayed signal for  $\bar{p}$ . The characteristic temporal structure of S1/S2 signals as measured by HeCal for  $\bar{d}$  is shown in Fig. 3. The time gap from S1(prompt) to S2(delayed) is related to the metastability of He capture and is statistically distributed with  $\tau_c \sim O(1 \mu s)$ . The S1 signal is related to the energy loss in the scintillating He, the amplitude measures the residual particle kinetic energy after the energy losses due to ToF and Vessel crossing. The duration of S1 signal is of the order of few ns that is the slowdown time for the particle in He.

The S2 signal is related to the number, velocity and trajectory of charged pions produced in the annihilation. The amplitude of S2 is not related to the particle kinetic energy, but measures the number of anti-nucleons. Typical S2 amplitude is higher than S1 and the duration of S2 signal is of the order of  $\sim$  ns that is the transit time of pions to move from the helium to the vessel. This peculiar S1/S2 signature, together with the measurement of  $\beta$  and  $dE/dX$  (see fig. 4) for the in-going particle, is very distinctive for  $\bar{d}$  and  $\bar{p}$  acting as a powerful rejection for other CR components. Combining prompt and delayed event information, ADHD can identify a single  $\bar{d}$  over 1000 background  $\bar{p}$  in the 60-140 MeV/n range. Considering the expected  $\bar{p}$  flux, this technique would be able to test the presence of  $\bar{d}$  in cosmic rays down to a flux of  $2-3 \times 10^{-6} (\text{m}^2 \text{s sr GeV/n})^{-1}$  with less than 1  $\bar{p}$  as background. It is interesting to note that the ADHD expected sensitivity is similar to the ones expected by other future experiments.



**Figure 3.** Typical HeCal signature for  $\bar{d}$  in ADHD detector.



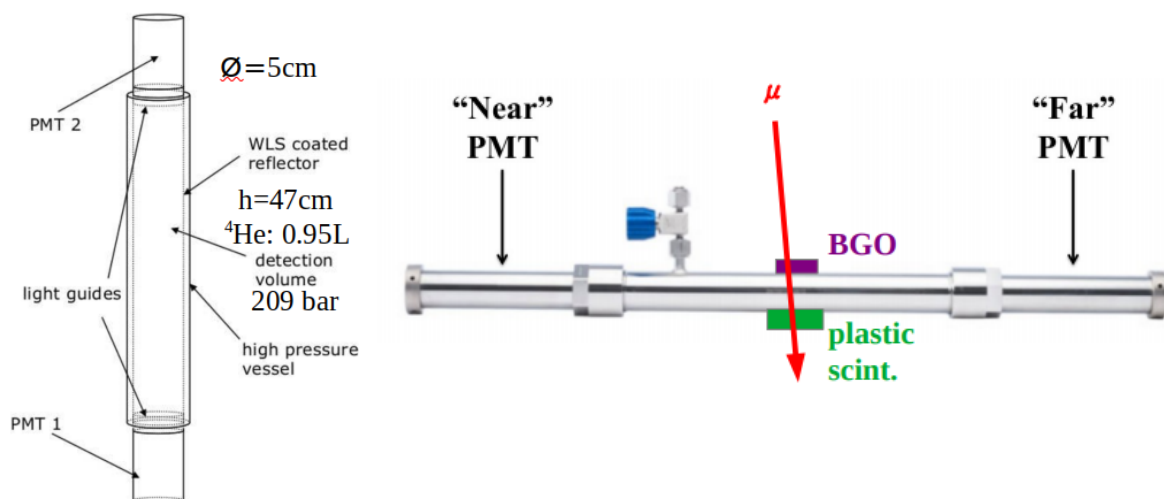
**Figure 4.**  $\bar{p}$  background is rejected thanks to the comparison of prompt track energy loss (and  $\beta$ ) as measured by ToF with the total kinetic energy as measured by HeCal (S1).

### 3. Measurements on HeCal prototype Arktis B-470.

The timing and energy resolution of HeCal detector are key parameters for the ADHD project. Some measurements on a prototype based on the fast neutron detector B-470 Arktis Radiation Detectors [15, 16] have been conducted in INFN-TIFPA laboratories in order to test the response of pressurized He gas as scintillator.

The Arktis B-470 detector consists of a stainless steel cylindrical chamber filled with pressurized  $^4\text{He}$  gas close to 200 bar and two Hamamatsu-R580 photomultiplier tubes (PMTs) installed at the two ends of the chamber (see Fig. 5). The inner wall of the chamber is lined with a wavelength shifter to convert ultra-violet He scintillation to the optical wavelengths for the PMTs. Response of this detector to fast neutrons are extensively studied, however for the aim of ADHD project the response of He scintillation to charged (stopping) particles is necessary.

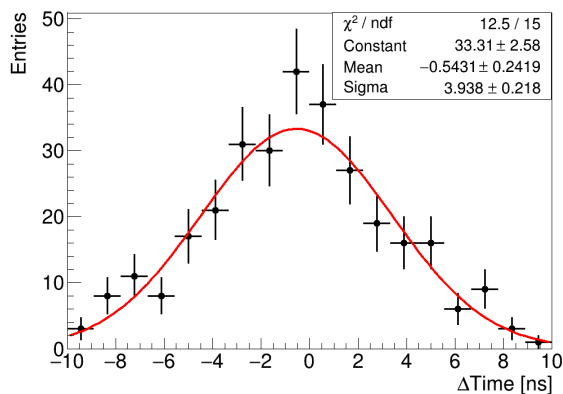
A preliminary test with muons,  $\mu$ , from cosmic rays was performed. The B-470 detector was operated in coincidence with a ( $10 \times 10 \times 30 \text{ mm}^3$ ) plastic scintillator and a small ( $5 \times 5 \times 20 \text{ mm}^3$ ) BGO scintillator, placed in a telescopic configuration to detect crossing  $\mu$ . The data was acquired



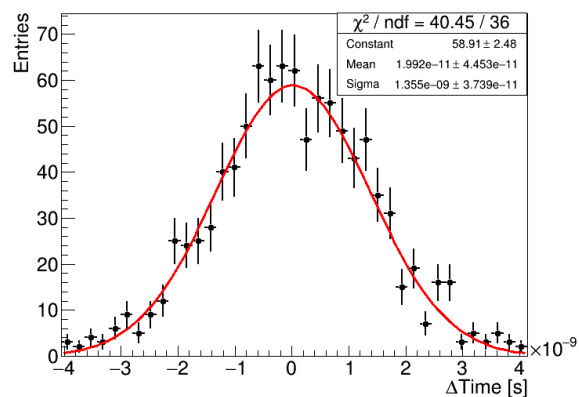
**Figure 5.** The Arktis B-470 detector used as a prototype of ADHD/HeCal.

with a LeCroy HDO9104-MS oscilloscope by sampling the detector waveforms at 20 Gs/s. The “minimum” energy deposition in the helium volume was obtained for  $\mu$  transversely crossing the detector diameter (depositing  $\sim 0.3$  MeV in 200 bar helium), conversely the maximum energy deposition was obtained for  $\mu$  crossing the whole detector (vertically placed, depositing  $\sim 3$  MeV in 200 bar helium).

The time difference between the plastic scintillator trigger and a single side of the B-470 detector is shown in Fig. 6 and Fig. 7 for the two configurations, respectively.



**Figure 6.** Time resolution of the Arktis B-470 detector measured with  $\mu$  transversely crossing the detector ( $\sim 300$  keV deposited).

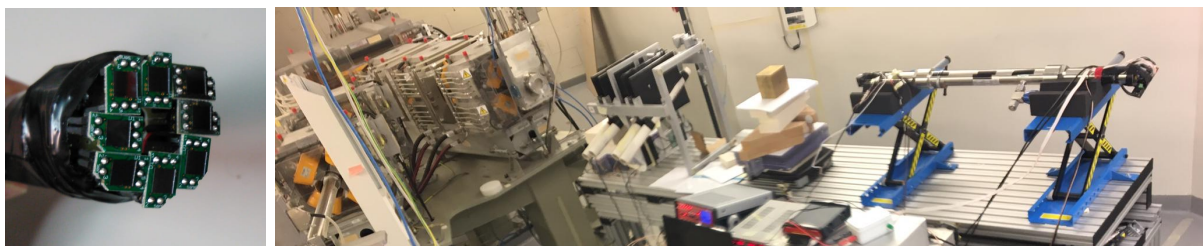


**Figure 7.** Time resolution of the Arktis B-470 detector measured with  $\mu$  crossing the whole detector length ( $\sim 3$  MeV deposited).

In particular the obtained (single side) time resolutions of  $3.9 \pm 0.2$  ns for  $\sim 0.3$  MeV and  $1.36 \pm 0.04$  ns for  $\sim 3$  MeV are in reasonable agreement with the time resolution of 2.4 ns achieved in [13] testing the same type of detector with a collimated  $^{137}\text{Cs}$  source. This preliminary measurement confirms the expected improvement of time resolution with the energy thanks to the increasing photon statistics, thus suggesting the capability to achieve a time resolution better than 1/2 ns for HeCal detecting  $\bar{d}$  in the 60-140 MeV/n energy range.

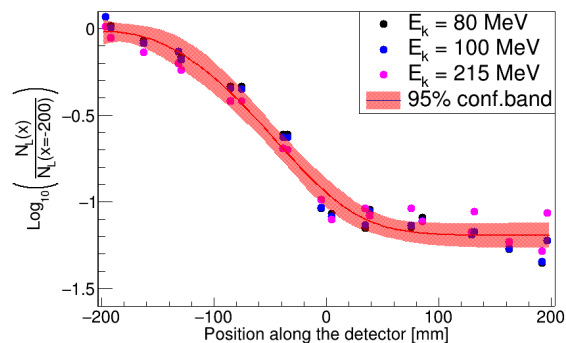
#### 4. Perspectives and conclusions

A test on the proton beam line in the experimental room of the Trento Proton Therapy Facility (Italy) [17] was pursued in the beginning of 2021. To allow the detection of the protons in the calorimeter avoiding PMT crossing, one PMT of the Arktis B-470 detector was replaced with an array of 8x Silicon Photo-Multipliers (see Fig. 8). The SiPM circular array is shielded by 20cm

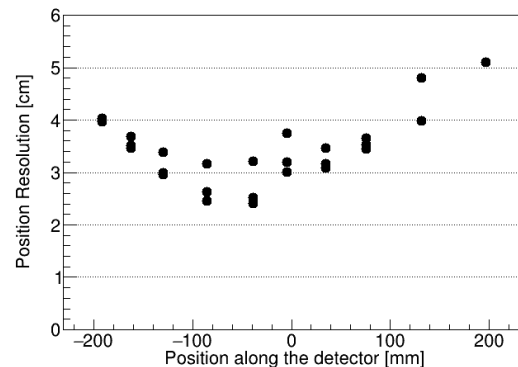


**Figure 8.** An array of 8x SiPM replaces one PMT of the Arktis B-470 detector (left picture). The detector prototype during the test at the Trento Proton Therapy facility (right picture).

of iron and a central hole of 1 cm allows the proton to enter in the Helium target by passing only through the optical window. The beam test is necessary to characterize the B-470 detector with protons in the energy range 70-230 MeV that is of interest for  $\bar{d}$  detection with ADHD. Analysis of data collected during the proton beam test is ongoing, here we briefly report the preliminary measurement of B-470 space resolution obtained with (80-100-215 MeV) protons transversely crossing the detector. The probability of photon detection at each side was simply modeled as:  $P(d) = P_0 e^{-(d/d_0)^2} + P_\infty$  where  $d$  is the distance from the photon detector, the photon attenuation length is  $d_0 = 123 \pm 2$  mm, whereas  $P_\infty = 1 - P_0 = 0.063 \pm 0.003$ . Figure 9 summarizes the photon attenuation measurements along the detector. Due to good symmetry observed in the results of SiPM array response, these data are also reported in fig. 9 along the PMT ones by simply flipping the sign of the measured position with respect to the detector center. By inverting the model of the photon absorption, it is possible to infer the beam position from the measured signal asymmetry. Figure 10 shows that the inferred position resolution is of the order of few cm. Position resolution of the SiPM (right) side is slightly worse than the PMT (left) side due to the smaller detection surface of the 8xSiPM array ( $2.9 \text{ cm}^2$  vs  $9 \text{ cm}^2$ ).



**Figure 9.** Photon attenuation measured by testing the Arktis B-470 detector with Protons transversely crossing the detector.



**Figure 10.** Space resolution achieved by the Arktis B-470 detector at the Proton beam.

#### 4.1. Acknowledgments

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