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# A compressive sensing-based computational method for the inversion of wide-band ground penetrating radar data

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**Abstract.** This work presents an innovative computational approach for the inversion of wide-band ground penetrating radar (*GPR*) data. The retrieval of the dielectric characteristics of sparse scatterers buried in a lossy soil is performed by combining a multi-task Bayesian compressive sensing (*MT-BCS*) solver and a frequency hopping (*FH*) strategy. The developed methodology is able to benefit from the regularization capabilities of the *MT-BCS* as well as to exploit the multi-chromatic informative content of *GPR* measurements. A set of numerical results is reported in order to assess the effectiveness of the proposed *GPR* inverse scattering technique, as well as to compare it to a simpler single-task implementation.

## 1. Introduction

During the last decades, ground penetrating radar (*GPR*) gained particular attention as a valid prospecting technology in many applicative scenarios dealing with the non-invasive investigation of buried domains, such as, for example, civil engineering, military operations, cultural heritage monitoring, and archaeology [1]-[4]. Thanks to such a success, several inverse scattering methodologies, both deterministic [5]-[8] and stochastic [9],[10], have been proposed to retrieve accurate and easy-to-interpret images of what lies below the interface from *GPR* measurements.

Differently from other prospecting tools, a key feature that characterizes *GPR* is the availability of wide-band data. As a matter of fact, such a frequency diversity can be seen as a powerful source of additional information in order to effectively tackle the *ill-posedness* and *non-linearity* issues of the subsurface imaging problem.

Within this framework, different techniques have been proposed in order to process multi-chromatic data. These approaches can be mainly classified in (i) multi-frequency (*MF*) [5],[7],[10] and (ii) frequency hopping (*FH*) [6] strategies. On the one hand, *MF* techniques are based on the simultaneous processing of the collected wide-band information [10], while on the other hand *FH* ones process each frequency sample in a cascaded manner, starting from the lowest frequency up to the highest one, and exploiting the acquired information at each step in order to initialize the successive one [6]. Given that, even if *MF* strategies could in principle benefit from a larger amount of information to perform the inversion, they must typically solve a more complex problem with a significantly larger number of unknowns with respect to *FH*-based methodologies.

Moreover, it is worth remarking that in many practical scenarios the unknown buried objects are sparse [5]. Given that, compressive sensing (*CS*)-based techniques [11]-[13] could be seen as an



interesting candidate in order to exploit such an *a-priori* information about the nature of the sought solution, by further regularizing the *GPR* inverse scattering problem. Following this line of reasoning, this work presents an innovative inversion methodology that effectively integrates a multi-task Bayesian *CS* (*MT-BCS*) solver [11] within a *FH* strategy in order to retrieve accurate reconstructions of the buried scenario. The *MT-BCS* solution at each stage of the frequency hopping loop is initialized in an innovative way, by exploiting the *acquired* information about the detected scatterers at the lowest frequencies. A set of numerical experiments is presented and discussed, in order to assess the effectiveness of the developed *FH-MT-BCS* technique, and to verify its superiority with respect to a single-task *BCS* (*ST-BCS*)-based implementation.

## 2. *GPR-IS* problem formulation and *FH-MT-BCS* solution approach

Let us consider a two-dimensional scenario consisting of two homogeneous half-spaces separated by a planar interface at  $y = 0$ . The upper half-space is free-space, while the lower half-space is occupied by a lossy medium having complex permittivity  $\varepsilon_{b,eq}(f) = \varepsilon_{rb}\varepsilon_0 - j\sigma_b/(2\pi f)$ . A buried investigation domain  $D_{inv}$  is illuminated by a set of  $V$   $z$ -oriented ideal line sources placed at fixed height above the interface and excited by a wideband current signal. Under the hypothesis that one or multiple targets are present in  $D_{inv}$ , the electromagnetic interaction between the field radiated by the sources and the half-space scenario generates a measured time-domain total field equal to

$$e_v^{tot}(\mathbf{r}_m, t) = e_v^{scat}(\mathbf{r}_m, t) + e_v^{inc}(\mathbf{r}_m, t); \quad v = 1, \dots, V; \quad m = 1, \dots, M \quad (1)$$

where  $\mathbf{r}_m$  is the location of the  $m$ -th receiver above the interface ( $m = 1, \dots, M$ ,  $M$  being the number of probes), while  $e_v^{inc}(\mathbf{r}_m, t)$  and  $e_v^{scat}(\mathbf{r}_m, t)$  denote the incident and scattered fields, respectively. Once  $e_v^{scat}(\mathbf{r}_m, t)$  has been isolated from  $e_v^{tot}(\mathbf{r}_m, t)$  [8], it is Fourier-transformed to the frequency domain, and a set of  $L$  uniformly spaced samples are extracted from the computed *GPR* spectrum within the 3dB bandwidth of the excitation signal ( $f_l \in [f_{min}, f_{max}]$ ,  $l = 1, \dots, L$ ). Accordingly, assuming a contrast source formulation, the following integral equation holds true at each  $l$ -th frequency

$$E_v^{scat}(\mathbf{r}_m, f_l) = \int_{D_{inv}} G^{ext}(\mathbf{r}_m, \mathbf{r}', f_l) J_v(\mathbf{r}', f_l) d\mathbf{r}'; \quad v = 1, \dots, V; \quad m = 1, \dots, M; \quad l = 1, \dots, L \quad (2)$$

where  $G^{ext}(\mathbf{r}_m, \mathbf{r}', f_l)$  is the external Green's function for the buried scenario [10], while  $J_v(\mathbf{r}, f_l) = E_v^{tot}(\mathbf{r}, f_l) \tau(\mathbf{r}, f_l)$  is the unknown equivalent current modelling the presence of the unknown objects in  $D_{inv}$ ,  $\tau(\mathbf{r}, f_l)$  being the contrast function defined at frequency  $f_l$  as follows

$$\tau(\mathbf{r}, f_l) = (\varepsilon_r(\mathbf{r}) - \varepsilon_{rb}) - j \frac{(\sigma(\mathbf{r}) - \sigma_b)}{2\pi f_l}; \quad l = 1, \dots, L. \quad (3)$$

Assuming an additional noise component on the scattered field and partitioning  $D_{inv}$  into  $N$  sub-domains, (2) can be arranged in matrix form as follows

$$\tilde{\mathbf{E}}_{v,l}^{scat} = \mathbf{G}_l^{ext} \mathbf{J}_{v,l} + \mathbf{n}_{v,l}; \quad v = 1, \dots, V; \quad l = 1, \dots, L \quad (4)$$

where  $\mathbf{G}_l^{ext}$  is the  $M \times N$  Green's matrix,  $\tilde{\mathbf{E}}_{v,l}^{scat} = \left\{ \Re[\tilde{E}_v^{scat}(\mathbf{r}_m, f_l)], \Im[\tilde{E}_v^{scat}(\mathbf{r}_m, f_l)] \right\}_{m=1, \dots, M}$ ,  $\mathbf{J}_{v,l} = \left\{ \Re[J_v(\mathbf{r}_n, f_l)], \Im[J_v(\mathbf{r}_n, f_l)] \right\}_{n=1, \dots, N}^T$ , while the noise vector is  $\mathbf{n}_{v,l} = \left\{ \Re[n_v(\mathbf{r}_m, f_l)], \Im[n_v(\mathbf{r}_m, f_l)] \right\}_{m=1, \dots, M}$ . Accordingly, the solution of the inverse problem at the  $l$ -th frequency step is formulated within the Bayesian framework as follows

$$\hat{\mathbf{J}}_{v,l} = \arg \left\{ \max_{\mathbf{J}_{v,l}} P(\mathbf{J}_{v,l} | \tilde{\mathbf{E}}_{v,l}^{scat}) \right\}; \quad v = 1, \dots, V \quad (5)$$

where  $P(\mathbf{J}_{v,l} | \tilde{\mathbf{E}}_{v,l}^{scat})$  is the posterior probability and  $\mathbf{J}_{v,l}$  are correlated among the different views ( $v = 1, \dots, V$ ). The solution of (5) is then found by means of a customized *MT-BCS* solver as

$$[\hat{\mathbf{J}}_{v,l}]^{ppt} = \left[ \text{diag}(\hat{\mathbf{a}}_l) + (\mathbf{G}_l^{ext})^* \mathbf{G}_l^{ext} \right]^{-1} (\mathbf{G}_l^{ext})^* \tilde{\mathbf{E}}_{v,l}^{scat}; \quad v = 1, \dots, V \quad (6)$$

where  $(\cdot)^*$  is the transpose conjugate operator and  $\hat{\mathbf{a}}_l = \{\hat{a}_{l,n}, n = 1, \dots, N\}$  is the shared vector of hyper-parameters estimated through a Relevant Vector Machine (*RVM*) solver.

The estimated contrast function  $\hat{\tau}(\mathbf{r}_n, f_l)$ ,  $n = 1, \dots, N$ , is then obtained from (6) by averaging over the different views ( $v = 1, \dots, V$ ) the ratio between retrieved currents and the corresponding retrieved total fields.

Since standard *BCS* implementations do not allow to exploit progressively *acquired* information about the solution (i.e., to use the reconstruction at the  $l$ -th frequency to initialize/guide the  $(l+1)$ -th step), an innovative "constrained" *RVM* solver has been developed as described in the following. A "filtering and clustering" procedure [6] is applied to the retrieved contrast in order to determine the region of interest (*RoI*)  $D_l \subset D_{inv}$  ( $D_l = D_{inv}$ ) in which the buried scatterers have been detected at frequency  $f_l$ . Then, the solution at frequency  $f_{l+1}$  is obtained through a "constrained" *RVM* by solving (6) limiting the search space only to those entries  $\hat{a}_{l+1,n}$  of  $\hat{\mathbf{a}}_{l+1}$  that correspond to the  $N_l < N$  cells belonging to  $D_l$ . More precisely, the *MT-BCS* solution at the  $f_{l+1}$  *FH* stage is aimed at retrieving the reduced vector of hyper-parameters

$$\hat{\mathbf{a}}_{l+1} = \{\hat{a}_{l+1,n}; n = 1, \dots, N_l\} \quad \text{subject to } \mathbf{r}_n \in D_l. \quad (7)$$

The *FH-MT-BCS* is terminated once one of the following conditions is met: (i) the last frequency step has been processed (i.e.,  $l = L$ ) or (ii)  $|A_l - A_{l-1}|/|A_l| \leq \eta$ , where  $A_l$  and  $A_{l-1}$  denote the area of  $D_l$  and  $D_{l-1}$ , respectively, and  $\eta$  is a suitable threshold.

### 3. Numerical results

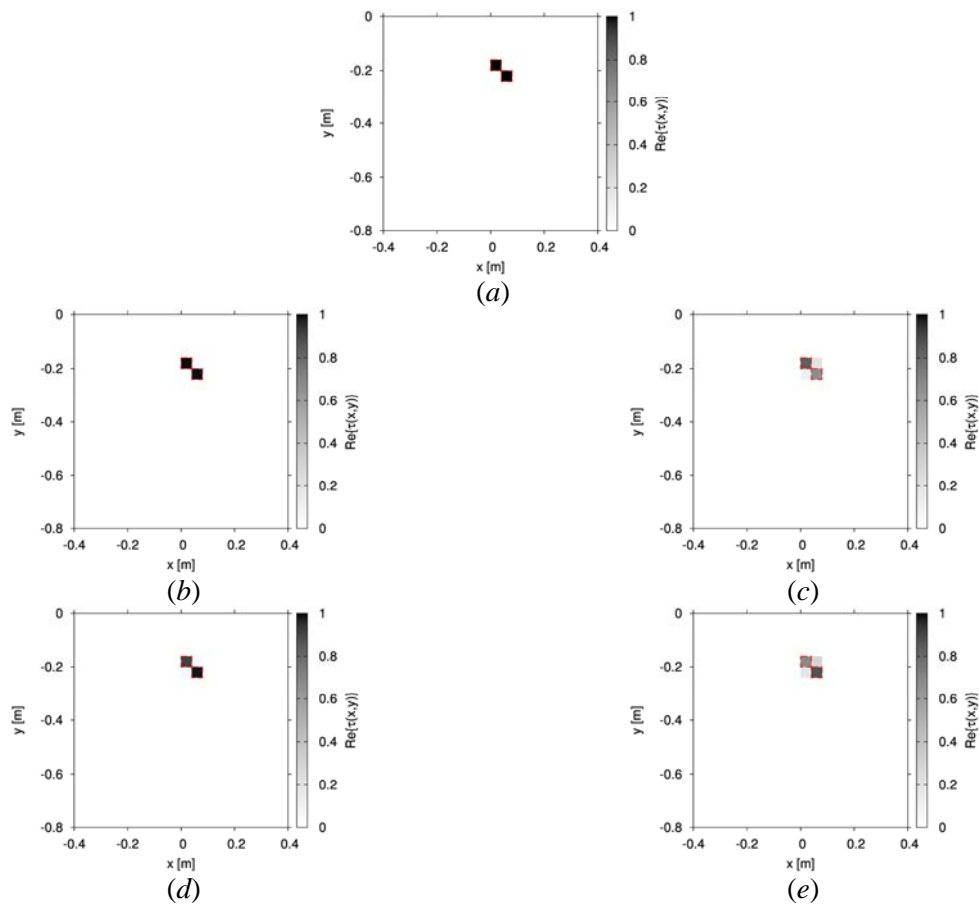
The aim of this section is to assess the performances of the proposed *FH-MT-BCS* algorithm. Towards this end, a set of numerical experiments is reported by generating time-domain synthetic data through the *GprMax2D* simulator [5]. In order to quantify the quality of each reconstruction, the following integral error is computed at a reference frequency  $\bar{f}$  arbitrarily chosen within the  $[f_{\min}, f_{\max}]$  range

$$\chi_{tot} = \frac{1}{N} \sum_{n=1}^N \frac{|\tau(\mathbf{r}, \bar{f}) - \hat{\tau}(\mathbf{r}, \bar{f})|}{|\tau(\mathbf{r}, \bar{f})| + 1}. \quad (8)$$

The investigation domain  $D_{inv}$  is a square of 0.8 m centred at (0,-0.4) m.  $V = 20$   $z$ -oriented infinitely long line sources at fixed height above the interface are used to illuminate it, while  $M = 19$  probes co-located with the sources are considered at each view to collect the scattered data. Moreover,  $D_{inv}$  is filled with a lossy medium with relative permittivity  $\epsilon_{rb} = 4.0$  and conductivity  $\sigma_b = 0.001$  S/m.  $L = 9$  uniformly spaced frequency samples are extracted from the computed *GPR* spectrum within the  $[f_{\min}, f_{\max}] = [200, 600]$  MHz bandwidth, while  $\bar{f} = 400$  MHz.

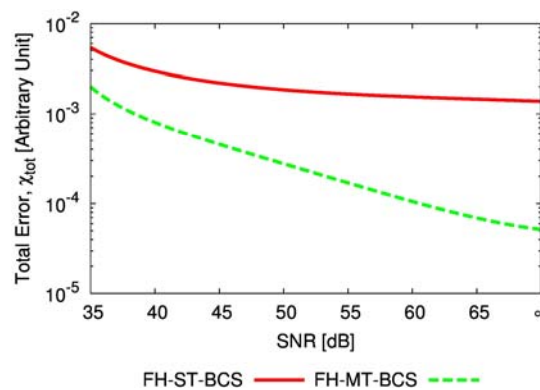
#### 3.1. Performances against noise

The first set of experiments deals with the retrieval of a buried scatterer composed by two diagonal pixels with relative permittivity  $\epsilon_r^{obj} = 5.0$  and conductivity of  $\sigma_b = 0.001$  S/m [Fig. 1(a)]. Figure 1 shows the retrieved dielectric profiles by the *FH-MT-BCS* when considering an additive white Gaussian noise on the measured time-domain total field with a signal-to-noise ratio defined as in [6] and equal to  $SNR = 60$  dB [Fig. 1(b)] and  $SNR = 40$  dB [Fig. 1(d)].



**Figure 1.** Numerical Assessment (Two diagonal pixels scatterer,  $\varepsilon_r^{obj} = 5.0$ ,  $\sigma^{obj} = 0.001$  S/m) – (a) Actual dielectric profile and retrieved solution by (b)(d) the *FH-MT-BCS* and (c)(e) the *FH-ST-BCS* methods for (b)(c)  $SNR = 60$  dB and (d)(e)  $SNR = 40$  dB on time-domain total field.

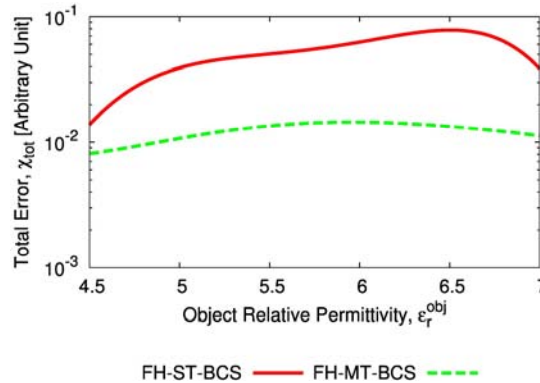
The reconstructions obtained by a single-task *BCS* implementation (*FH-ST-BCS*) are reported, as well [ $SNR = 60$  dB - Fig. 1(c);  $SNR = 40$  dB - Fig. 1(e)], this latter not imposing any correlation among the several views in solving (6). As it can be seen, the *FH-MT-BCS* significantly overcomes the *FH-ST-BCS*, as it is also visible by looking at the computed errors vs. the  $SNR$  reported in Fig. 2.



**Figure 2.** Numerical Assessment (Two diagonal pixels scatterer,  $\varepsilon_r^{obj} = 5.0$ ,  $\sigma^{obj} = 0.001$  S/m) – Behaviour of the reconstruction error versus the  $SNR$  on time-domain total field.

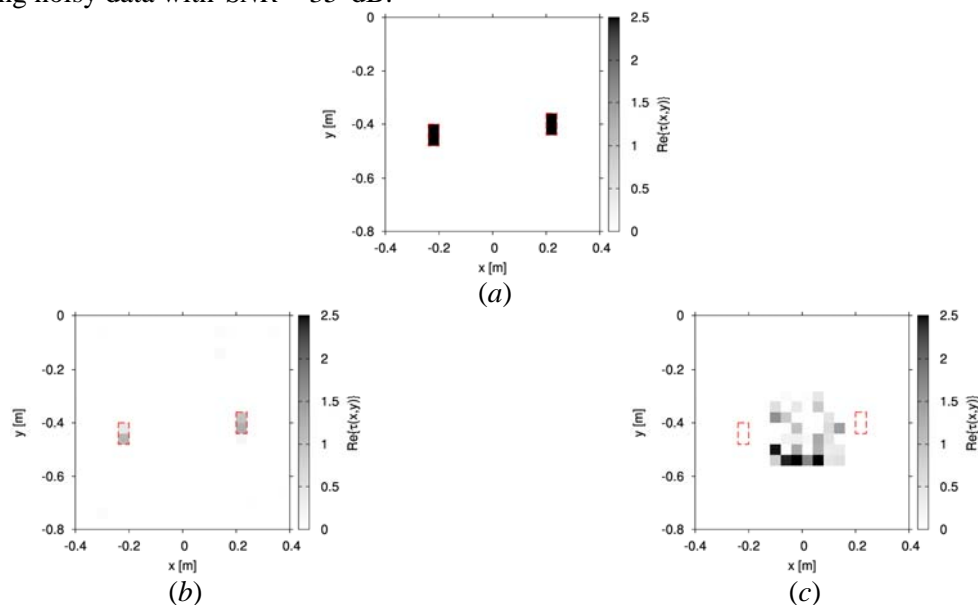
### 3.2. Performances against the relative permittivity of the buried targets

This second set of numerical benchmarks is aimed at assessing the performances of the *FH-MT-BCS* when considering a variation of the relative permittivity of the unknown buried objects in  $D_{inv}$ .



**Figure 3.** Numerical Assessment (Two bars scatterer,  $\sigma^{obj} = 0.001$  S/m,  $SNR = 35$  dB) – Behaviour of the reconstruction error versus the relative permittivity of the object.

Towards this end, Figure 3 compares the reconstruction errors made by the *FH-MT-BCS* and the *FH-ST-BCS* methods when dealing with the reconstruction of two separated sparse targets whose relative permittivity is varied in the range  $\epsilon_r^{obj} \in [4.5, 7.0]$  [ $\sigma^{obj} = 0.001$  S/m - Fig. 4(a)] and considering noisy data with  $SNR = 35$  dB.



**Figure 4.** Numerical Assessment (Two bars scatterer,  $\epsilon_r^{obj} = 6.5$ ,  $\sigma^{obj} = 0.001$  S/m,  $SNR = 35$  dB) – (a) actual dielectric profile and retrieved solutions by the (b) *FH-MT-BCS* and (c) *FH-ST-BCS* techniques.

As it can be observed, lower errors are obtained by the *FH-MT-BCS* in all cases (Fig. 3). As a matter of fact, significantly better reconstructions are provided by the proposed method despite the non-negligible amount of noise, as it can be seen by looking at the retrieved profiles in Fig. 4 for  $\epsilon_r^{obj} = 6.5$ . Even if providing a slight underestimation of the contrast, the *FH-MT-BCS* is able to correctly detect the scatterer support [Fig. 4(b)], while the single-task version yields an unsatisfactory result [Fig. 4(c)].

#### 4. Conclusions

An innovative inverse scattering technique has been presented in order to deal with the retrieval of the dielectric characteristics of a buried investigation domain starting from wide-band *GPR* measurements. The proposed *FH-MT-BCS* methodology efficiently combines the information coming from multi-chromatic data thanks to an innovative initialization strategy of the *RVM* solver, as well as benefits from the regularization capabilities of the *MT-BCS* solver. The main novelty of the work is the innovative integration of the *MT-BCS* within the *FH* strategy in order to progressively exploit the acquired information about the solution at the low-pass reconstructions, thanks to the introduction of a customized "constrained" *RVM* solver. Some numerical results have been presented, in order to validate the effectiveness of the *FH-MT-BCS*, as well as to show its superior performances with respect to a single task-based solution within the same framework. Future work will be devoted at extending the developed algorithm to deal with more realistic scenarios by considering, for example, a non-static behaviour of both object and soil electromagnetic properties, the presence of metallic/PEC scatterers, as well as a full three-dimensional formulation of the subsurface scattering equations.

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