Reporting a Misunderstanding in Definition of Functionally Graded Porosity

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Abstract

This paper reports a misunderstanding in relating Young's modulus to density of functionally graded porous structural members through the well-known power law. A list of recently published references which are subjected to this is presented. It is found out that this misunderstanding may cause undeniable errors in the mechanical responses in both static and dynamic loading conditions that here we quantified by implementing finite element analysis.

Keywords: Porous materials; Functionally Graded Porosity; Open-Cell porosity; Beams, Plates and Shells.

1. Introduction

In the last decades, the field of solid mechanics has been faced a huge variety of modern human-made materials and consequently, many researchers and designers focused to understand how these materials can be implemented to enhance the performance of structural members in new demanded loading situations. Among them fiber-reinforced composites [1] and functionally graded materials [2], especially metal-ceramic combination, attracted one of the most attention. The former was selected as the main candidate for the purpose of light-weight design and the latter was developed to solve challenges in high temperature conditions. Cellular solids [3], also known as porous materials, despite of their remarkable potential in optimizing both functionality and light-weighing were neglected for some decades from a structural point of view. However, thanks to unprecedented developments in manufacturing processes, especially digital manufacturing, they are becoming important alternatives in structural design and parallelly the number of studies on their structural behavior are growing.

Recently, the idea of graded porosity, is introduced to maximize the capability of cellular solids in terms of optimized design for structural members like beams, plates, panels, and shells under a variety of loading conditions [4–6]. Unlike uniform cellular solids, for graded porous ones the amount of porosity is varied as a function of position generally in three dimensions, classifying them as non-homogeneous materials. The present paper is organized to report a misunderstanding widely observed in the recent literature for defining graded porosity and it is clarified how much it may affect the accuracy of obtained results.

2. Explaining the misunderstanding in graded porosity

Consider a bulk material of density ρ_0 , Young's modulus E_0 , and Poisson's ratio of v_0 as the parent material of cellular solid. By introducing pores, the density $\rho(x, y, z)$ in every point within the media can be altered as:

$$\rho(x, y, z) = \rho_0 (1 - e_m F(x, y, z))$$
(1)

where (x, y, z) denotes the position, e_m is the porosity parameter, and $0 \le F(x, y, z) \le 1$ is the gradual function. The maximum of porosity, e_m , happens where F(x, y, z)=1, and its minimum is

located where F(x, y, z)=0 with no pore. Here, we focus on open-cell cellular solids whose Young's modulus is related to its density through a power law of index *n* as [3]:

$$E(x, y, z) = E_0 \left(\frac{\rho(x, y, z)}{\rho_0}\right)^n = E_0 \left(1 - e_m F(x, y, z)\right)^n$$
(2)

For the case of uniform distribution of pores, the gradual function is canceled in Eqs. (1) and (2), and the uniform density, $\overline{\rho}$, and uniform Young's modulus, \overline{E} , are calculated as:

$$\bar{\rho} = \rho_0 (1 - e_m) \tag{3a}$$

$$\bar{E} = E_0 (1 - e_m)^n = E_0 (1 - e_0) \tag{3b}$$

In Eq. (3b), e_0 directly reflects the fraction of reduction in Young's modulus because of pores and is dependent to the porosity parameter, e_m , as:

$$e_0 = 1 - (1 - e_m)^n$$
, $e_m = 1 - \sqrt[n]{1 - e_0}$ (4)

The misunderstanding happens if someone mixes up Eqs. (3b), (4) and (2) and evaluates the Young's modulus of graded porosity wrongly as:

$$\hat{E}(x, y, z) = E_0 (1 - e_0 F(x, y, z))$$
(5)

A list of 86 recent references, all assuming n=2, in the literature which wrongly used Eq. (5) instead of Eqs. (2), categorizing in three different structural geometries i.e. beams, plates, and shells under a variety of loading conditions named static loading, vibration, dynamic load, stability, and wave propagation, are presented in Table 1.

3. Estimation of Error

The error (%) in evaluation of Young's modulus in every point of a graded porous material is calculated as:

Error in Young's Modulus (%) =
$$\frac{\hat{E}(x, y, z) - E(x, y, z)}{E(x, y, z)} \times 100$$
(6)

The error for uniform distribution of porosity is zero, however, for graded porosity is related to gradual function, F(z), porosity parameter, e_m , and power law index, n. As, all the references listed in Table 1 used the power law index n=2, suggested by [3], the same value is considered in the present report.

To demonstrate a numerical insight to this error, a structural member of thickness h where the porosity varies across its thickness is considered. Note that z axis is set along the thickness of member and its origin locates at the mid-plane. As Fig.1 shows, three general types of porosity variation have been chosen called Pyramid, P-type Sandglass, S-type, and diamond, D-type. For the P-type, maximum and minimum of the porosity locate at the bottom and top surfaces of the member, resulting in an unsymmetrical distribution with respect to the mid-plane. For the D-type, as a symmetric distribution, the minimum of the porosity is at the surfaces while the maximum locates at the mid-plane and the S-type is vice versa.

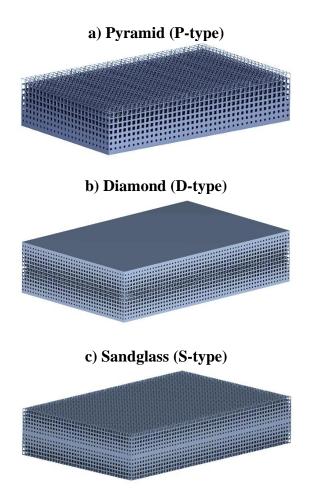


Fig 1. Types of variation of porosity along the thickness a) Pyramid, b) Diamond, and c) Sandglass.

	Static Load	Vibration	Dynamic Load	Stability	Wave Propagation
Beam	[7] [8] [9] [10] [11] [12]	[7] [13] [14] [15] [16] [17] [18] [19] [8] [20] [9] [21] [22] [23] [24] [25] [10] [26] [27]	[28] [29] [22] [30] [31]	[16] [32] [9] [33] [11] [12]	[34] [35] [36]
Plate	[37] [38] [39] [40]	[41] [37] [42] [43] [44] [45] [46] [47] [48] [49] [50] [51] [52] [53] [54] [55] [56] [57] [58] [59] [60] [61] [62] [63]	[64] [65]	[47] [48] [39] [66]	[67]
Shell	[68] [69] [70]	[71] [72] [73] [74] [75] [76] [77] [78] [79] [80] [81] [82] [83] [84] [85] [86]	[87] [88] [89]	[79] [90] [91]	[92]

Then, four different mathematical functions have been assumed: linear, parabolic, cubic, and cosine. Accordingly, totally 12 definitions for graded function, F(z), are obtained which are listed in Table 2.

Table 2: Definition of gradual function, $F(z)$, across the thickness of structural member.								
	P-type	D-type	S-type					
^a Polynomial	$\left(\frac{1}{2} - \frac{z}{h}\right)^m$	$1 - \left \frac{2z}{h}\right ^m$	$\left \frac{2z}{h}\right ^m$					
Cosine	$\cos\left(\frac{\pi z}{2h} + \frac{\pi}{4}\right)$	$\cos\left(\frac{\pi z}{h}\right)$	$\left(1-\cos\left(\frac{\pi z}{h}\right)\right)$					

^{*a*} m=1: linear, m=2: parabolic, m=3: cubic

Considering the Cosine function, which is widely addressed by the references in Table 1, Fig. 2 and Fig. 3 depict the variation of error in the calculation of Young's modulus along the thickness for highand low-range of porosity parameter, e_m , respectively. High range porosity is popular in lightweighting optimal designs, while the low range is usually found in soil mechanics analysis where the porosity hardly exceeds to 0.5. Regardless the value e_m , the error is zero at the surfaces for the Ptype, and at the surfaces and mid-plane for the D- and S- types where Eqs. (2) and (5) are equal. Besides, one should notice that for different types of porosity distributions, P-type, D-type, and Stype, the location of maximum value of error shifts with respect to the mid-plane which can affect the rigidity of member significantly.

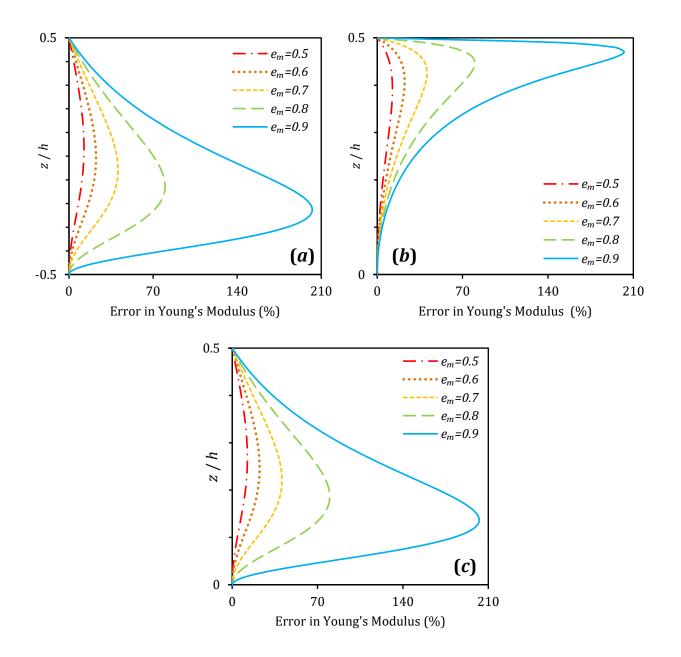


Fig 2. The errors in evaluating Young's modulus via Eq. (5) through the thickness of structural members for cosine distribution and high-range porosity, a) P-type, b) S-type, and c) D-type. The error values for S- and D- types have shown for the half of thickness due to symmetry.

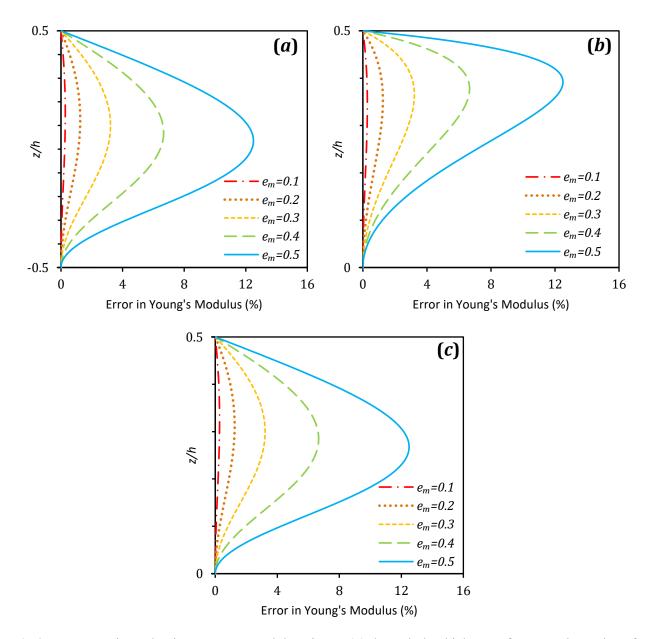


Fig 3. The errors in evaluating Young's modulus via Eq. (5) through the thickness of structural members for cosine distribution and low-range porosity, a) P-type, b) S-type, and c) D-type. The error values for S- and D- types have shown for the half of thickness due to symmetry.

The maximum values of the error for evaluating Young's modulus across the thickness can be calculated by substituting Eqs. (2) and (5) into Eq. (6). After some calculations and simplifications, one can show that the maximum value of error across the thickness is only a function of porosity parameter, e_m , as follows:

Maximum Error in Young's Modulus (%) =
$$\frac{e_m^2}{4(1-e_m)} \times 100$$
 (7)

Interestingly, this maximum is the same for all 12 different porosity graded functions introduced in Table 2, although the location of this maximum is not, as shown in Figs. (2) and (3). The variation of maximum error of Eq. (7) with respect to the porosity parameter, e_m , is plotted in Fig. (4). It is revealed that for a relatively low value of porosity, e_m =0.5, the maximum error is 12.5% while it significantly increases up to 202.5% for a high value of porosity, e_m =0.9. In general, it is concluded that the underlined misunderstanding causes an irrefutable error in the evaluation of Young's modulus for the case of graded porosity along the thickness of structural members. One should pay enough attention that the positive values of error for all cases mean Young's modulus evaluated by Eq. (5) is an overestimation which results in an unreal stiffness for structural members.

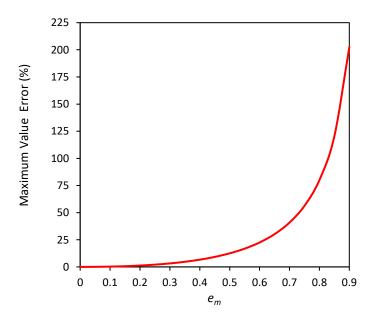


Fig 4. Maximum value of error (%) in evaluation of Young's modulus of graded porosity by Eq. (7).

The next question is how much the reported error in the evaluation of Young's modulus may affect the response of graded porous structural members. Based on the well-known theories of beams, plates, and shells [93], the influence of Young's modulus is reflected in governing equations of structural members as stretching, A, stretching-bending, B, and bending, D, stiffnesses, obtaining by integration across the thickness of members, h, as:

$$(A, \hat{A}) = \int_{-h/2}^{+h/2} \frac{(E, \hat{E})}{1 - \nu_0^2} dz$$
(8a)

$$(B,\hat{B}) = \int_{-h/2}^{+h/2} \frac{(E,\hat{E})}{1-\nu_0^2} z \, dz \tag{8b}$$

$$(D,\widehat{D}) = \int_{-h/2}^{+h/2} \frac{(E,\widehat{E})}{1-\nu_0^2} z^2 dz$$
(8c)

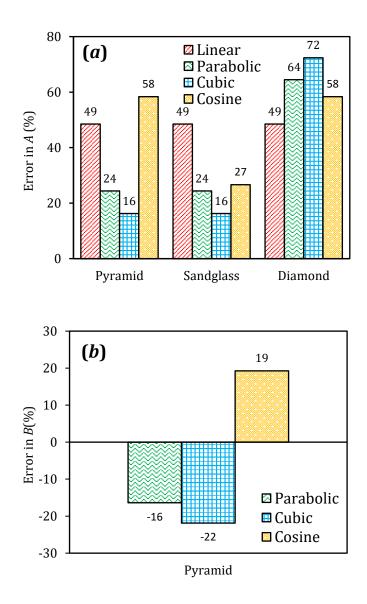
where \hat{A} , \hat{B} , and \hat{D} are the mistaken stiffnesses using \hat{E} defined by Eq. (5) and ν_0 is the Poisson's ratio of the bulk material which is considered to be constant and is neglected for beam-like members. Due to symmetry, stretching-bending stiffness is zero for S- and D- types. One can calculates the corresponding errors in stiffnesses as:

Error in Stiffness (%) =
$$\frac{(\hat{A}, \hat{B}, \hat{D}) - (A, B, D)}{(A, B, D)} \times 100$$
(9)

It is noted that this error is independent of the thickness of the member and the Poisson's ratio as they are canceled from the fraction.

Fig. 5 shows the error in stiffnesses, *A*, *B*, and *D* because of the reported misunderstanding in the variation of Young's modulus along the thickness of a structural member for the porosity distributions introduced in Table 2. To emphasize the error, the situation where the value of porosity gradually varies across the thickness between 0 to the maximum of 0.99, is considered. Although $e_m = 0.99$ may not practically reachable for uniform porosity, however, it is meaningful for graded porosity as an ideal upper bond. The values of error presented on the top of bars in Fig. 5, show remarkable errors

for all graded porosities. For stretching and bending stiffnesses, the errors are positives which means members are considered wrongly stiffer for both in-plane and out-of-plane deformations. Nevertheless, stretching-bending stiffness shows no error for Linear, negative errors for both Parabolic and Cubic, and positive error for Cosine gradual functions. A negative error means the reported misunderstanding underestimates the coupling between in-plane and out-of-plane deformations while positive one overestimates it.



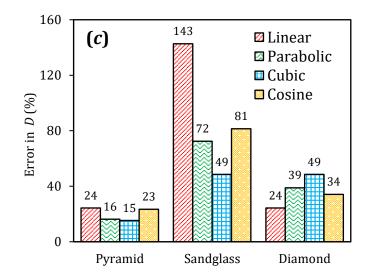


Fig 5. Error in stiffnesses for e_m =0.99. a) Stretching Stiffness, *A*, b) Stretching-Bending Stiffness, *B*, and c) Bending Stiffness, *D*.

The mentioned errors in Eq. (9) and demonstrated in Fig. (5) directly affect the responses of structural members under all static and dynamic loading conditions. For the structural members with low slenderness ratios, usually called thin structures, the relationship between responses and stiffnesses is straightforward. As an example, deflection, w, transverse natural frequencies, ω , and critical buckling load, P_{cr} of a thin member having symmetric material distribution with respect to its midplane (S-type and D-type) are simply related to the bending stiffness. Hence, it is possible to find out how much \hat{w} , $\hat{\omega}$, and \hat{P}_{cr} , predicted based on Eq. (5), are wrong:

$$w \propto \frac{1}{D} \to w = (\widehat{D}/D)\widehat{w}$$
 (10a)

$$\omega \propto \sqrt{D} \to \omega = \sqrt{D/\hat{D}}\,\hat{\omega}$$
 (10b)

$$P_{cr} \propto D \rightarrow P_{cr} = (D/\widehat{D})\widehat{P}_{cr}$$
 (10c)

4. Finite Element Simulation Case studies

To make a better insight to the errors even for thick structural members, bending and free vibration analysis of a square plate with simply supported boundary conditions and graded porosity along its thickness is numerically examined based on the finite element analysis using COMSOL Multiphysics commercial software. To achieve this, a square shell model with the length side of *a* is modeled and different slenderness ratios are defined by fixing the side length and setting a proper value of thickness, *h*. The model is meshed by implementing 20 elements through the thickness, and 30×30 in-plane elements, a total of 18000 quadratic elements. The well-known first-order shear deformable theory (FSDT) is implemented to also assure the accuracy of results for thick plates. In the case of bending analysis, a uniform distributed load, q_0 , is applied on the surface of the plate. The Poisson's ratio is considered ν_0 =0.3. As an isotropic but non-homogeneous material, density is defined as a function along the thickness by Eq. (1) with the bulk value of ρ_0 =7800 kg/m³, while Young's modulus is defined either by Eq. (2) or Eq. (5) considering the bulk value of E_0 =200 GPa to calculate the resulting errors in maximum deflections and fundamental natural frequencies for all 12 different porosity distribution functions of Table 2, thanks to the simple functional definition of properties in COMSOL Multiphysics. Finally, stationary and eigenfrequency analyses are performed and the errors are evaluated as follow:

Error in defelction, frequency (%) =
$$\frac{(\widehat{w}, \widehat{\omega}) - (w, \omega)}{(w, \omega)} \times 100$$
 (11)

First, the validation study of the present finite element model is performed. The length side and the thickness are set to a=1000 mm and h=1 mm to model a thin enough plate and a uniform pressure of $q_0=1$ Pa is applied. Table 3 compares FEM results to those obtained from the exact analytical solution for deflection of thin simply supported square plates in [93] as:

$$w_{max} = 0.00406 \frac{q_0 a^4}{D} \tag{12}$$

It is noted that the comparison is only possible for symmetric distribution of properties with respect to the mid-plane. In the other word, Eq. (12) is not valid for P-type porosity distribution as stretchingbending stiffness, *B*, is not zero. An excellent agreement between FEM and exact solution is observed that validate the accuracy of numerical model. Besides, Table 3 reveals remarkable errors in the deflection of thin plates caused by using Eq. (5). It is observed that for S- and D-types where Eq. (10) is valid, the value of error is consistent with the value of error in bending stiffness, D, presented in Fig. 5(c) where the maximum error is for Sandglass-Linear type and the lowest one belongs to Diamond-Linear porosity distribution. The high value of error for Pyramid-Cosine type reflects all the errors in A, B, and D stiffnesses as a stretching-bending coupling happens.

			<i>h/a</i> =0.001				
Туре		w _{max} [Exact]	w _{max} [FB	Error [%]			
		Eq. (12)	Eq. (2)	Eq. (5)	Eq. (11)		
	Linear	2.1509	2.1522	0.8870	-58.7864		
	Parabolic	0.9553	0.9559	0.5544	-42.0023		
Sandglass	Cubic	0.6584	0.6588	0.4436	-32.6655		
	Cosine	1.1117	1.1124	0.6130	-44.8939		
	Linear	0.3676	0.3678	0.2957	-19.6030		
	Parabolic	0.5131	0.5134	0.3697	-27.9899		
Diamond	Cubic	0.6584	0.6588	0.4436	-32.6655		
	Cosine	0.4659	0.4662	0.3475	-25.4612		
	Linear	-	1.4348	0.6652	-53.6381		
	Parabolic	-	0.7553	0.4669	-38.1835		
Pyramid	Cubic	-	0.5655	0.3960	-29.9735		
	Cosine	-	3.5867	1.3097	-63.4845		

Table 3: Validation of the finite element model by comparing the maximum deflection of thin, porous square plate, w_{max} [mm], to the exact solution of thin plates in Eq. (11) and evaluation of the errors [%] caused by using Eq. (5) instead of Eq. (2). Note: $a=1000 \text{ mm}, h=1 \text{ mm}, E_0=200 \text{ GPa}, v_0=0.3, e_m=0.99, q_0=1 \text{ Pa}.$

After validation, the values of errors in the deflection and the fundamental frequency of the porous square plates are calculated for various values of slenderness ratio in Table 4 and Table 5, respectively. It is seen that the errors are affected significantly by the types of porosity distribution while they are not much sensitive to the slenderness ratio. It is noted that although the values of deflection and natural frequencies are dependent to the material properties of the bulk, however, for a linear analysis, the values of error are independent as the bulk properties are cancelled in Eq. (11).

One should remember that Eq. (5) overestimates Young's modulus and wrongly considers the plate stiffer. It means using Eq. (5) underestimates and overestimates the deflections and the frequencies, respectively, which justifies the error percentage sign in Table 4 and Table 5.

Туре		<i>h/a</i> =0.01, <i>q</i> ₀ =1 kPa			<i>h/a</i> =0.1,	<i>h/a</i> =0.1, <i>q</i> ₀ =1 MPa			<i>h/a</i> =0.2, <i>q</i> ₀ =10 MPa		
		Eq. (2)	Eq. (5)	Error [%]	Eq. (2)	Eq. (5)	Error [%]	Eq. (2)	Eq. (5)	Error [%]	
	Linear	2.1525	0.8872	-58.79	2.1863	0.9099	-58.38	2.8609	1.2236	-57.23	
	Parabolic	0.9561	0.5546	-41.99	0.9773	0.5717	-41.51	1.3021	0.7792	-40.16	
Sandglass	Cubic	0.6589	0.4437	-32.66	0.6766	0.4589	-32.18	0.9125	0.6311	-30.85	
	Cosine	1.1126	0.6132	-44.89	1.1352	0.6310	-44.41	1.5048	0.8565	-43.08	
	Linear	0.3682	0.2960	-19.61	0.4020	0.3187	-20.71	0.6305	0.4846	-23.14	
Diamond	Parabolic	0.5140	0.3700	-28.02	0.5701	0.4041	-29.11	0.9252	0.6344	-31.43	
	Cubic	0.6596	0.4440	-32.68	0.7380	0.4895	-33.67	1.2196	0.7842	-35.70	
	Cosine	0.4667	0.3478	-25.46	0.5162	0.3791	-26.56	0.8331	0.5925	-28.88	
	Linear	1.4351	0.6654	-53.63	1.4689	0.6882	-53.15	1.9642	0.9464	-51.82	
Pyramid	Parabolic	0.7555	0.4671	-38.18	0.7767	0.4841	-37.67	1.0513	0.6698	-36.29	
	Cubic	0.5656	0.3962	-29.96	0.5833	0.4114	-29.48	0.7959	0.5717	-28.17	
	Cosine	3.5872	1.3100	-63.481	3.6368	1.3413	-63.12	4.7337	1.7952	-62.08	

Table 4: Finite element evaluation of the errors [%] in maximum deflection, w_{max} [mm] of moderately thick, porous square plate caused by using Eq. (5) instead of Eq. (2) for various slenderness ratios. Note: a=1000 mm, h=10, 100, and 200 mm, $E_0=200$ GPa, $v_0=0.3$, $e_m=0.99$, $q_0=1$ kPa, 1 MPa, and 10 MPa.

Table 5: Finite element evaluation of the errors [%] in fundamental natural frequency, ω [Hz] of moderately thick, porous square plate caused by using Eq. (5) instead of Eq. (2) for various slenderness ratios. Note: $a=1000 \text{ mm}, h=10, 100, \text{ and } 200 \text{ mm}, E_0=200 \text{ GPa}, v_0=0.3, e_m=0.99.$

Туре		<i>h/a</i> =0.01			<i>h/a</i> =0.1			<i>h/a</i> =0.2		
		Eq. (2)	Eq. (5)	Error [%]	Eq. (2)	Eq. (5)	Error [%]	Eq. (2)	Eq. (5)	Error [%]
Sandglass	Linear	21.745	33.870	55.76	214.75	332.77	54.96	414.61	633.71	52.85
	Parabolic	28.326	37.191	31.30	278.59	364.16	30.72	532.06	687.63	29.24
	Cubic	32.195	39.233	21.86	315.74	383.29	21.39	598.63	719.79	20.24
	Cosine	26.861	36.183	34.71	264.50	354.67	34.09	506.74	671.50	32.51
Diamond	Linear	52.572	58.636	11.54	496.33	557.54	12.33	866.21	987.12	13.96
	Parabolic	54.225	63.911	17.86	506.83	602.12	18.80	868.47	1047.5	20.61
	Cubic	55.003	67.036	21.88	511.21	627.80	22.81	867.49	1080.3	24.53
	Cosine	54.570	63.208	15.83	510.90	596.29	16.71	878.03	1039.9	18.44
Pyramid	Linear	26.63	39.11	46.85	261.41	382.04	46.15	496.85	717.32	44.37
	Parabolic	31.87	40.53	27.18	312.06	395.30	26.67	589.61	739.55	25.43
	Cubic	34.75	41.52	19.49	375.98	421.79	12.18	700.73	782.81	11.71
	Cosine	19.69	32.58	65.48	288.59	405.99	40.68	543.78	756.46	39.11

To investigate the effect of the value of porosity, Fig. 6 and Fig. 7 respectively demonstrate the errors in deflection and the fundamental natural frequency of the simply supported square plate versus variation of porosity parameter, e_m , for all the graded porosity introduced in Table 2. Two slenderness ratios, h/a=0.1 and 0.2 are assumed to see how the errors are correlated to this ratio. As expected, increasing the porosity parameter, e_m , rises the errors and it is observed that the errors have low dependency on the slenderness ratio.

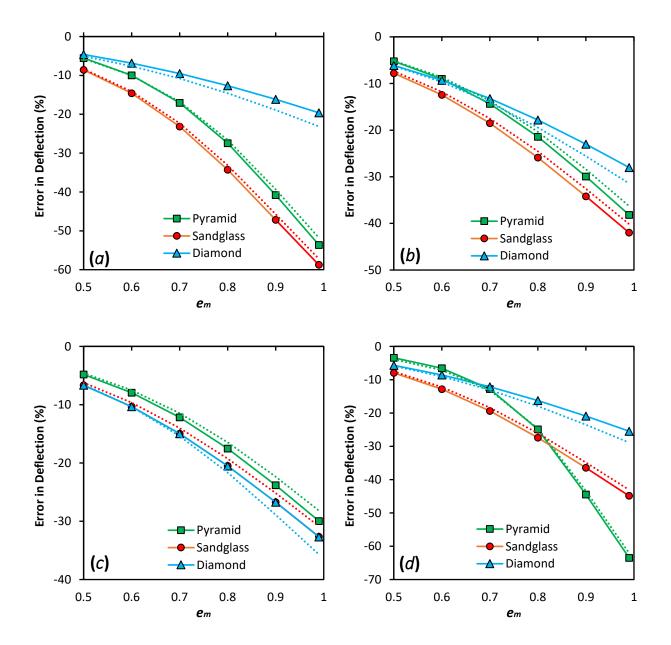


Fig 6. FEM evaluation of the errors in deflection of simply supported square plates with graded porosity. a) Linear, b) Parabolic, c) Cubic, and d) Cosine. Continues lines: h/a=0.01, Dot lines: h/a=0.2.

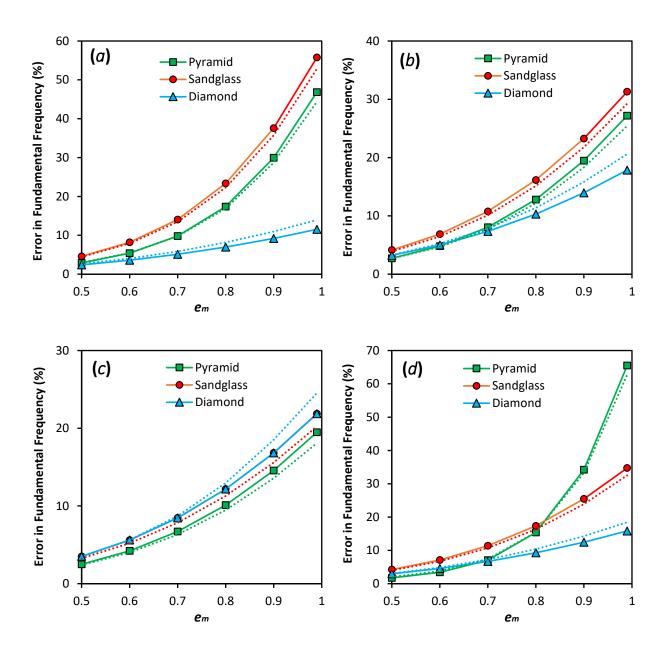


Fig 7. FEM evaluation of the errors in fundamental natural frequency of simply supported square plates with graded porosity. a) Linear, b) Parabolic, c) Cubic, and d) Cosine. Continues lines: h/a=0.01, Dot lines: h/a=0.2.

5. Conclusion

Reviewing the recent literature on beam-, plate-, and shell-like structural members highlights notable attention to the application of modern engineering materials, in particular, the concept of functionally graded porosity due to its advantage in light-weighting optimal designs thanks to remarkable recent developments in digital manufacturing. However, an exhaustive review reveals a misunderstanding in the definition of the graded density-stiffness relationship repeated by many recent publications. The aim of the present paper is to clarify this issue and quantify the value of error that may happen in the evaluation of the stiffness of structural members. Various types of graded porosity distribution across the thickness of structural members are assumed and it is shown that unfortunately, this misunderstanding causes an undeniable stiffness overestimation that significantly affects the predicted behavior of these members in both static and dynamic loading conditions, especially for high values of maximum porosity parameter. As a case study, finite element analysis has been implemented to quantify the value of error in the deflection and frequency response of graded porous square plates.

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Data Availability

The raw/processed data required to reproduce these findings cannot be shared at this time due to technical or time limitations.

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