

Modular dynamics in double cones

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Abstract

We investigate the relation between the actions of Tomita-Takesaki modular operators for local von Neumann algebras in the vacuum for free massive and massless bosons in four dimensional Minkowskian spacetime. In particular, we prove a long-standing conjecture that says that the generators of the mentioned actions differ by a pseudo-differential operator of order zero. To get that, one needs a careful analysis of the interplay of the theories in the bulk and at the boundary of double cones. After introducing some technicalities, we prove the crucial result that the vacuum state for massive bosons in the bulk of a double cone restricts to a K.M.S. state at its boundary, and that the restriction of the algebra at the boundary does not depend anymore on the mass. The origin of such result lies in a careful treatment of classical Cauchy and Goursat problems for the Klein-Gordon equation as well as the application of known general mathematical techniques, concerning the interplay of algebraic structures related with the bulk and algebraic structures related with the boundary of the double cone, arising from QFT in curved spacetime. Our procedure gives explicit formulas for the modular group and its generator in terms of integral operators acting on symplectic space of solutions of massive Klein-Gordon Cauchy problem.

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1 Introduction

This paper deals with the task of finding explicit expressions of the (generators of the) Tomita-Takesaki modular groups of automorphisms of von Neumann algebras localized in double cones in 4-dimensional Minkowski spacetime for a free massive scalar field theory in the (cyclic and separating) vacuum state. In the following discussion we shall deal exclusively with Minkowski spacetime and hence we will not mention it explicitly anymore.

Algebraic quantum field theory [Ha92] is a particularly fruitful approach for unveiling structural features of specific (class of) theories and, as such, it has been largely applied to study spin-statistics, PCT theorems and superselection structures. In this respect, one of the milestones has been the connection of Tomita and Takesaki modular theory

[Ta01]¹ to quantum statistical mechanics, via KMS conditions [BR96I, BR96II]. Modular theory has also been connected with geometry-dynamics and thermodynamics, as results from the pioneering papers of Bisognano-Wichmann [BW75, BW76] and Sewell [Se82]. The first gave the interpretation of the modular group of automorphisms of von Neumann algebras (generated by Wightman fields) localized in wedge regions – for a generic scalar field theory in the vacuum Poincaré-invariant state – as the boost transformations leaving the wedges fixed; whereas the second offered a physical interpretation of the Bisognano-Wichmann analysis in terms of the Unruh and Hawking effects. These ideas have been further developed in various directions, which especially emphasize the geometric interpretation of the action of the modular groups. Let us recall some of them.

In the case of massless quantum field theories Buchholz [Bu77] proved that the generator of the modular group, w.r.t. the vacuum state, associated with a von Neumann algebra localized in the (forward) light-cone is related to the scaling transformation. Hislop and Longo [HL95] showed that, for free scalar field theories invariant under the group (suitably defined) of conformal transformations, the modular groups associated with classes of regions as light-cones, wedges and double cones, are all related to the generators of the geometric transformations leaving all respective regions fixed. In particular, they are all related via suitable elements of the conformal group. A purely algebraic version of these results has been presented in [BGL93, BGL95], using the results of a landmark paper by Borchers [Bo92]. Reconstruction of some spacetime and their symmetries can also be done, mainly using modular conjugations, as shown by Buchholz, Summers and collaborators [BDFS, BMS], and, similarly, one can define free theories solely in terms of the modular data [BGL02], basing the approach on the abstraction of the Bisognano-Wichmann correspondence and the representations of the Poincaré group. In the case of conformal field theories on the circle, modular groups (always related to a cyclic and separating vacuum state) continues to have a geometrical flavor when related to single intervals, but, as shown recently in [MLR10] also in case the localization has been weakened to several disjoint intervals. The last paper was inspired by the results of Casini and Huerta [CH09] who in turn were able to put on solid ground the computation of the (resolvent of the) generator of the modular group (for massless free Fermi fields in two dimensions) developed abstractly in the nice paper of Figliolini and Guido [FG89] (see also [BJL02, Sa06]). However, it appears clear that the geometric behavior of the modular groups is tightly related either to the massless condition or to a particular choice of the state. Indeed, in the massive case for localizations different from wedges, for instance on double cones, one should not expect such a geometric behavior. Moreover, even in wedges localizations, it has been proved by Borchers and Yngvason [BY99] that other choice of states, for instance KMS, would certainly break down the geometric connection. This last can only be retrieved in particular situations like near the apex of a double cone or the edge of the wedge. The application of the milestone contributions by Wiesbrock about half-sided modular conditions [Wi93, Wi93b, Wi97]

¹For short lucid introductions we refer to the papers of Guido [Gu08], Lledó [Lle09] and Summers [Su05]. For more extensive considerations we cannot help better than suggesting the reading of Borchers' paper [Bo00].

was crucial to achieve those results.

It is the main goal of this paper to start a direct attack to one of the main questions left open in the last years (see [SW00] for more): What is the explicit form of the generator of the modular group for massive (scalar) theories in the vacuum localized in a double cone? With this paper, we are able to give a precise answer only in the case of free theories, while the ambitious task of dealing with more general theories is left to future investigations.

A well-known conjecture, sometimes referred to as “Fredenhagen’s conjecture,” says that the generator of the modular group for the massive free scalar theory (in the vacuum) for a double cone should differ from the massless one by a pseudodifferential operator of order zero. A possible strategy to proof this conjecture appeared in several places, for instance in [SW00] and [Sa05], and relies upon the idea that, if considering the vacuum for the massless theory as a positive functional (state) in the (Fock) Hilbert space of the massive one, then the modular groups of both theories associated with this state should coincide. Notice that, for the massive case, that vacuum would be the “wrong” vacuum. However, by one of the celebrated theorems of Connes [Co74], the modular automorphisms of the massive case referred to the “wrong” and correct vacua, would differ only by a cocycle, i.e. by the action of a (continuous) family of unitary operators in the vacuum Hilbert space, which is thought not to change the pseudodifferential nature of the generator of the massive case. However, so far, no one was able to turn this ideas into a real proof, not even the authors of the present paper who, instead, took a radically different route.

1.1 Main ideas and structure of the work

After having defined the general geometric scenario within Sec.2, we will tackle the main problem to give an explicit representation of the modular group of the free massive Klein-Gordon theory of a diamond as follows. Instead of relying upon choices of vacua and construction of cocycles (more on this later), our main idea is to compare the massive free scalar theory in the double cone with its shadow at the boundary. This simple idea has been a source of many interesting recent results of QFT in curved spacetime [DMP06, Mo06, Mo08, DMP09, DMP09b], and produced sharp tools for the comparison of theories in the boundary-bulk correspondence.

Generally speaking, one starts by considering the algebra of observables $\mathscr{W}_m(D)$ (the Weyl C^* -algebra of the free Klein-Gordon field with mass $m \geq 0$) localized in the interior of a region D of a globally hyperbolic spacetime, whose boundary contains a 3-dimensional light-like submanifold with the structure of a conical surface V . In the case considered in this paper, D is nothing but a double cone in Minkowski spacetime and V is the lower light-cone forming one-half of the boundary of D . Under some hypotheses (e.g. see [Mo06, DMP09] for asymptotically flat and cosmological backgrounds), exploiting the symplectic structure of the space of solutions of Klein-Gordon equation, it is possible to define a C^* -algebra of observables $\mathscr{W}(V)$ localized on V , together with an embedding map (C^* -algebra isometric homomorphism) $\ell_m^{V,D} : \mathscr{W}_m(D) \rightarrow \mathscr{W}(V)$. Remarkably, the boundary algebra $\mathscr{W}(V)$ does not depend on the mass m anymore, since the information

about m is completely included in the embedding ℓ_m^{VD} as we shall prove in Sec.3.

After having discussed some preparatory issues in Section 3, especially concerning the characteristic Cauchy problem (also known as the *Goursat problem*) of the KG equation with data assigned on V , the further step in our construction is to establish, in Sec.4, the existence of an algebraic state $\lambda : \mathscr{W}(V) \rightarrow \mathbb{C}$ that induces (the restriction to D of) Minkowski vacuum $\omega_m : \mathscr{W}_m(D) \rightarrow \mathbb{C}$ through ℓ_m^{VD} . In other words $\omega_m = \lambda \circ \ell_m^{VD}$. Once again, as a remarkable feature, λ turns out to be independent from m . The idea is then to pass to the von Neumann algebras, $\pi_m(\mathscr{W}_m(D))''$, $\pi_\lambda(\mathscr{W}(V))''$ respectively associated with the GNS representations $(\mathcal{H}_m, \pi_m, \Psi_m)$ of ω_m and $(\mathcal{H}_\lambda, \pi_\lambda, \Psi_\lambda)$ of λ and to study the interplay of the corresponding modular groups, referred to the respective GNS cyclic and separating vectors Ψ_m and Ψ_λ . As we shall prove in Sec.5, though $\mathscr{W}_m(D)$ is smaller than $\mathscr{W}(V)$ – that is ℓ_m^{VD} is not surjective – the algebras coincide when promoted to von Neumann algebras. This is because ℓ_m^{VD} turns out to be implemented by a unitary $Z_m : \mathcal{H}_m \rightarrow \mathcal{H}_\lambda$ which preserves the cyclic vectors, $Z_m \Psi_m = \Psi_\lambda$, identifying the von Neumann algebras: $Z_m \pi_m(\mathscr{W}_m(D))'' Z_m^{-1} = \pi_\lambda(\mathscr{W}(V))''$. This technically complicated result will be achieved by carefully focusing the characteristic Cauchy problem and enlarging the initially defined symplectic space of solutions of KG equation following the analysis performed in Sec.3.

To go on towards the final result, the crucial observation is that the modular group of the boundary has a geometric meaning related to the limiting value to the boundary of the same conformal Killing field studied, in the massless case, by Hislop and Longo in [HL95]. This remark is used in Sec.6 and permits to explicitly construct the group of symplectic isomorphisms of the space of solutions of KG equation that corresponds, through quantization procedure, to the action of the modular group of $(\pi_m(\mathscr{W}_m(D))'', \Psi_m)$ in the one-particle Hilbert space. In the massless case that one-parameter symplectic group will be found to coincide to that presented in [HL95]. Furthermore, the same procedure will give rise to an explicit formula for the infinitesimal generator of the modular group (represented in the symplectic space of KG solutions). The difference of the generator for $m > 0$ and the analog for $m = 0$ (that is a differential operator) turns out to be, as conjectured, a differential operator of class L^0 . More precisely, at the end of Sec.6 we shall prove that the difference of generators belongs to the class $L_{1,1}^0$.

2 Geometric features

Here and in the following (\mathbb{M}, g) is the four dimensional Minkowski spacetime assumed to be time-oriented, g the flat metric with signature $-, +, +, +$ and ∇ is the covariant derivative associated with the metric g . We fix a preferred Minkowskian frame used throughout and adopt the notation $x = (t, \mathbf{x}) \in \mathbb{R} \times \mathbb{R}^3$ for the coordinates of the frame. In the rest of the paper we shall use also the following *light-cone coordinates* on \mathbb{M} :

$$u \doteq \frac{t + \|\mathbf{x}\|}{2} \in \mathbb{R}, \quad v \doteq \frac{t - \|\mathbf{x}\|}{2} \in \mathbb{R}, \quad \omega \doteq \frac{\mathbf{x}}{\|\mathbf{x}\|} \in \mathbb{S}^2, \quad (1)$$

where $u - v \geq 0$, and $\|\mathbf{x}\|$ stands for the euclidean norm in \mathbb{R}^3 .

Another useful tool will be the *signed squared geodesic distance* $\sigma : \mathbb{M} \times \mathbb{M} \rightarrow \mathbb{R}$ of pairs of points in \mathbb{M} with coordinates $x = (t, \mathbf{x})$ and $x' = (t', \mathbf{x}')$:

$$\sigma(x, x') \doteq -(t - t')^2 + \|\mathbf{x} - \mathbf{x}'\|^2. \quad (2)$$

Concerning the *causal structure* of spacetimes, we adopt the definitions as in [Wa84]. In particular, if $K \subset \mathbb{M}$, then $J^+(K)$ and $J^-(K)$ denote, as subsets of \mathbb{M} , the *causal future* and *causal past* of K , respectively. Restricting to some open set $A \subset \mathbb{M}$, and if $K \subset A$, $J^\pm(K; A)$ denote the analogs referred to $(A, g|_A)$ viewed as a spacetime on its own right.

2.1 Relevant properties of conformal Killing vector fields of \mathbb{M}

If Y is a *conformal Killing field* [Sc08] of Minkowski spacetime, we indicate by Υ^Y the *local one-parameter group* of diffeomorphisms generated by Y (sometimes simply called the *flow* of Y). It is the smooth map

$$\Upsilon^Y : \mathcal{O} \rightarrow \mathbb{M},$$

where $\mathcal{O} \subseteq \mathbb{R} \times \mathbb{M}$ an open subset containing the set $\{0\} \times \mathbb{M}$, and which satisfies the relations

- (a) $\Upsilon^Y(0, x) = x$;
- (b) $\Upsilon^Y(\tau' + \tau, x) = \Upsilon^Y(\tau', \Upsilon^Y(\tau, x))$, whenever both sides are defined.

The induced action β_τ^Y on (scalar) functions f is individuated by the requirement

$$(\beta_\tau^Y f)(\Upsilon^Y(\tau, x)) = J_\tau^Y(x)^{-1/4} f(x), \quad (3)$$

where J_τ^Y is the Jacobian *in Minkowskian coordinates* of the the map $x \mapsto \Upsilon^Y(\tau, x)$ and it is assumed that it does not vanish. The exponent $-1/4$ may be fixed differently [Sc08], however our choice will turn out to be useful later when discussing properties of the massless Klein-Gordon equation. In this paper we shall be concerned with the *infinitesimal generator* γ^Y of β_τ^Y :

$$(\gamma^Y f)(x) \doteq \left. \frac{d}{d\tau} \right|_{\tau=0} (\beta_\tau^Y f)(x) = -Y(f)(x) - \frac{1}{4} (\nabla_a Y^a) f(x), \quad (4)$$

There is an interesting interplay between the squared distance σ in \mathbb{M} and the transformations induced by the conformal Killing vectors:

$$\sigma(\Upsilon^Y(\tau, x), \Upsilon^Y(\tau, x')) = J_\tau^Y(x)^{1/4} J_\tau^Y(x')^{1/4} \sigma(x, x') \quad (5)$$

(see for instance [Ha92, DMS97], collecting all the results for every subgroup of conformal transformation of \mathbb{M}). Taking the derivative at $s = 0$, the useful relation arises:

$$Y_x(\sigma)(x, x') + Y_{x'}(\sigma)(x, x') = \frac{1}{4} (\nabla_a Y^a(x) + \nabla_a Y^a(x')) \sigma(x, x') \quad (6)$$

where the notation Y_x means that the action of Y is on the variable showed as a lower index. The previous relation implies that, if F is a differentiable function on the real line and γ^Y is defined in (4):

$$\gamma_x^Y F(\sigma(x, x')) + \gamma_{x'}^Y F(\sigma(x, x')) = \frac{1}{4} (\nabla_a Y^a(x) + \nabla_a Y^a(x')) (F(\sigma(x, x')) + \sigma F'(\sigma(x, x'))).$$

2.2 The standard double cones $D(p, q)$

Double cones in \mathbb{M} are open regions generated by the choice of two points, p and q , such that q lies in the chronological future of p , and are defined by the position

$$D(p, q) \doteq \text{Int}(J^+(p) \cap J^-(q)) .$$

Given $D(p, q)$, it is always possible to construct a Minkowskian reference frame, said to be *adapted to $D(p, q)$* , such that p and q stay on the t -axis, q in the future of p which, in turn, coincides to the origin of the coordinates. That coordinate frame is determined up to a spatial 3-rotation and, referring to those coordinates, $D(p, q)$ takes the canonical form, where $a \doteq \sqrt{-\sigma(p, q)}/2$ is the *radius* of double cone:

$$D(p, q) \equiv \{(t, \mathbf{x}) \in \mathbb{R}^4 \mid |t - a| + \|\mathbf{x}\| < a\} . \quad (7)$$

The boundary ∂D is decomposed into *three* disjoint sets: A couple of lightlike conical 3-surfaces, $V(p, q)$ and $W(p, q)$ with tips p and q respectively and a 2-sphere, $C(p, q)$, defining the common base of $V(p, q)$ and $W(p, q)$. Using light-coordinates as in (1):

$$V(p, q) \doteq \{(\omega, u, v) \in \mathbb{S}^2 \times \mathbb{R}^2 \mid v = 0, u \in [0, a]\} , \quad (8)$$

$$W(p, q) \doteq \{(\omega, u, v) \in \mathbb{S}^2 \times \mathbb{R}^2 \mid u = a, v \in (0, a]\} , \quad (9)$$

$$C(p, q) \doteq \{(\omega, u, v) \in \mathbb{S}^2 \times \mathbb{R}^2 \mid v = 0, u = a\} . \quad (10)$$

The embedding of the cylinder $\mathbb{S}^2 \times (0, a)$ into \mathbb{M} :

$$\kappa : \mathbb{S}^2 \times (0, a) \ni (\omega, u) \mapsto (u, u\omega) \in \mathbb{M} \quad (11)$$

realizes a *blow up* of the cone $V(p, q)$, removing its lower tip p .

To conclude, we notice that $(D(p, q), g \upharpoonright_{D(p, q)})$ is globally hyperbolic. Indeed, it is strongly causal, since \mathbb{M} is such and, for $r, s \in D(p, q)$, the intersection of causal sets $J^+(r; D(p, q)) \cap J^-(s; D(p, q))$ is compact since it coincide with the analogous set $J^+(r) \cap J^-(s)$ referred to the whole spacetime \mathbb{M} , which is compact in turn.

2.3 A relevant conformal Killing vector for D

For a fixed double cone $D(p, q)$, taking account of the general expression [Sc08] of smooth conformal Killing fields everywhere defined in \mathbb{M} , one easily sees that, up to a constant nonvanishing factor, there is a unique (nonvanishing) conformal Killing field $X^{(p, q)}$ of

\mathbb{M} such that its flow leaves $V(p, q) \cup C(p, q)$ fixed, $X^{(p, q)}$ is nonspacelike and invariant under 3-rotations on $V(p, q) \cup C(p, q)$. It is:

$$X^{(p, q)} \doteq (t - a)\mathbf{x} \cdot \nabla_{\mathbf{x}} + \left(\frac{t^2 + \mathbf{x}^2}{2} - at \right) \partial_t \quad \text{so } \nabla_b X^{(p, q)b} = 4(t - a), \quad (12)$$

where we have adopted a Minkowskian coordinate frame adapted to $D(p, q)$. It will be useful to have also the form of the field in light-cone coordinates, namely,

$$X^{(p, q)} = u(u - a)\partial_u + v(v - a)\partial_v .$$

Notice that $X^{(p, q)}$ is tangent to both $V(p, q)$ and $W(p, q)$ and it vanishes on p, q and $C(p, q)$, therefore its flow leaves $D(p, q), C(p, q), V(p, q), W(p, q), p, q$ separately fixed. As $\overline{D(p, q)} = D(p, q) \cup \partial D(p, q)$ is compact, the one-parameter group of diffeomorphisms generated by $X^{(p, q)}$ is global on $\overline{D(p, q)}$, i.e. its parameter ranges to the full real line \mathbb{R} when the orbit starts from a point in $\overline{D(p, q)}$ and the composition of flows in the right-hand side of (b) in section 2.1 is always defined if $x \in \overline{D(p, q)}$ and $\tau, \tau' \in \mathbb{R}$.

Remark 2.3.1. In the following, without loss of generality, we shall deal with a preferred double cone, indicated by D , assuming that $a = 1$. In that case, $p, q, C(p, q), V(p, q), W(p, q)$ and $X^{(p, q)}$ will be denoted by o, o^+, C, V, W and X respectively.

3 Symplectic and conformal structures for wave equations

In the following $\mathcal{S}_m(\mathbb{M})$ is the real vector space of the smooth solutions ϕ of the Klein-Gordon equation with mass $m \geq 0$,

$$g^{ab}\nabla_a\nabla_b\phi - m^2\phi = 0, \quad (13)$$

which are compactly supported on smooth spacelike Cauchy surfaces Σ of \mathbb{M} . (Notice that we use the Einstein summation convention and from now on italic indices a, b, c, \dots will run from 1 to 4.)

Furthermore,

$$\sigma_{\mathbb{M}}(\phi_1, \phi_2) \doteq \int_{\Sigma} (\phi_2 \nabla_{n_{\Sigma}} \phi_1 - \phi_1 \nabla_{n_{\Sigma}} \phi_2) d\mu_{\Sigma} \quad (14)$$

is the standard nondegenerate symplectic form on $\mathcal{S}_m(\mathbb{M})$, where n_{Σ} is the future-oriented normal-to- Σ unit vector and μ_{Σ} is the standard measure induced on Σ by g . $\sigma_{\mathbb{M}}$ is Cauchy-surface-independent in view of the Klein-Gordon equation and Stokes-Poincaré's theorem. $(\mathcal{S}_m(\mathbb{M}), \sigma_{\mathbb{M}})$ is a real symplectic space.

In the following $\Delta_m : C_0^{\infty}(\mathbb{M}) \rightarrow C^{\infty}(\mathbb{M})$ denotes the *causal propagator* of the equation (13), that is the difference of the *advanced* and *retarded fundamental solutions*, which exist because (\mathbb{M}, g) is globally hyperbolic.

3.1 Symplectic spaces associated to D and V and the Goursat problem.

All the mentioned structures can be defined replacing (\mathbb{M}, g) with a generic smooth globally hyperbolic spacetime [BGP07]. In particular, as the double cone $D \subset \mathbb{M}$ is open and $(D, g|_D)$ is globally-hyperbolic, one can define the symplectic space $(\mathcal{S}_m(D), \sigma_D)$ of Klein-Gordon solutions in D with compactly supported Cauchy data on Cauchy surfaces Σ_D of D , σ_D being defined as the right-hand side of (14) replacing Σ for Σ_D .

Remark 3.1.1. It is worth noticing that there is a unique linear map preserving the relevant symplectic forms (i.e. a *symplectic homomorphism*) $L_m^{\mathbb{M}D} : \mathcal{S}_m(D) \rightarrow \mathcal{S}_m(\mathbb{M})$, such that

$$(L_m^{\mathbb{M}D} \phi)|_D = \phi \quad \text{and} \quad \text{supp}(L_m^{\mathbb{M}D} \phi) \subset J^+(D) \cup J^-(D) \quad \text{for every } \phi \in \mathcal{S}_m(D). \quad (15)$$

Indeed, every smooth spacelike Cauchy surface Σ_D of D extends to a smooth spacelike Cauchy surface Σ of \mathbb{M} . As the Cauchy problem is well posed in every globally hyperbolic spacetime [Wa84, BGP07] and $\sigma_{\mathbb{M}}|_{\mathcal{S}_m(D) \times \mathcal{S}_m(D)} = \sigma_D$, there is a linear map $L_m^{\mathbb{M}D} : \mathcal{S}_m(D) \rightarrow \mathcal{S}_m(\mathbb{M})$ mapping $\phi \in \mathcal{S}_m(D)$ to $L_m^{\mathbb{M}D} \phi \in \mathcal{S}_m(\mathbb{M})$ with the same Cauchy data on Σ as ϕ on Σ_D . By the uniqueness property of the Cauchy problem, (15) turns out to be satisfied. The same argument proves that $L_m^{\mathbb{M}D}$ is uniquely determined. Finally, we remark that $L_m^{\mathbb{M}D}$ is injective because a linear map between two symplectic spaces that preserves the symplectic forms is injective provided the symplectic form of the domain is non degenerate, the proof being elementary.

Let us consider the symplectic structures associated with the lower boundary V of D . Equipping V with the topology induced by \mathbb{M} , we define the real vector space

$$\mathcal{S}(V) \doteq \{ \Phi : V \rightarrow \mathbb{R} \mid \Phi = uf|_V, \text{ for some } f \in C^\infty(\mathbb{M}), \text{supp}(\Phi) \text{ is compact} \}, \quad (16)$$

where u is the light-coordinate appearing in (8). By direct inspection one easily sees that every $\Phi \in \mathcal{S}(V)$ results to be smooth away from o , it vanishes in a neighborhood of $u = 1$, $|\Phi(\omega, u)| \leq C_\Phi |u|$, uniformly in ω , for some $C_\Phi \geq 0$ and all u -derivatives (of every order) of Φ are bounded functions.

$\mathcal{S}(V)$ becomes a symplectic space when equipped with the non degenerate symplectic form:

$$\sigma_V(\Phi_1, \Phi_2) \doteq \int_V \left(\Phi_2(\omega, u) \frac{\partial \Phi_1(\omega, u)}{\partial u} - \Phi_1(\omega, u) \frac{\partial \Phi_2(\omega, u)}{\partial u} \right) d\omega du \quad \Phi_1, \Phi_2 \in \mathcal{S}(V) \quad (17)$$

where $d\omega$ is the standard measure on the the 2-sphere surface \mathbb{S}^2 with unitary radius.

The relation between $(\mathcal{S}_m(D), \sigma_D)$ and $(\mathcal{S}(V), \sigma_V)$ is stated within the following proposition, where henceforth $\lim_{\rightarrow V}$ denotes the pointwise limit toward V of functions.

Proposition 3.1.2. *For every fixed $m \geq 0$, the map*

$$L_m^{VD} : \mathcal{S}_m(D) \ni \phi \mapsto u \lim_{\rightarrow V} \phi, \quad (18)$$

is an injective symplectic homomorphism from $(\mathcal{S}_m(D), \sigma_D)$ to $(\mathcal{S}(V), \sigma_V)$.

Proof. $\phi \in \mathcal{S}_m(D)$ is the restriction to D of an element $L_m^{\mathbb{M}D}\phi \in \mathcal{S}_m(\mathbb{M})$. Therefore $L_m^{VD}\phi$ is the restriction to V of a smooth function $f \doteq L_m^{\mathbb{M}D}\phi \in C^\infty(\mathbb{M})$. Furthermore $\text{supp}(L_m^{\mathbb{M}D}\phi)$ does not intersect C so that the support of $u(L_m^{\mathbb{M}D}\phi)|_V$ is compact in V as requested in the definition of $\mathcal{S}(V)$. (Let Σ be the Cauchy surface of \mathbb{M} at $t = 1$ and Σ_D its restriction to D . The Cauchy data of $L_m^{\mathbb{M}D}\phi$ are included in a compact $K \subset \Sigma_D$ so that $\text{supp}(L_m^{\mathbb{M}D}\phi) \subset J^+(K) \cup J^-(K)$; however $(J^+(K) \cup J^-(K)) \cap C = \emptyset$ by construction and thus $\text{supp}(L_m^{\mathbb{M}D}\phi) \cap C = \emptyset$.) Hence $L_m^{VD}(\phi) \in \mathcal{S}(V)$, by definition of $\mathcal{S}(V)$. To conclude it is enough proving that the linear map $L_m^{VD} : \mathcal{S}_m(M) \rightarrow \mathcal{S}(V)$ preserves the symplectic forms. This fact entails that L_m^{VD} is injective as σ_D is nondegenerate. For $\phi_1, \phi_2 \in \mathcal{S}_m(M)$, the 3-form $\eta \doteq \sqrt{|g|}g^{ea}(\phi_1\partial_e\phi_2 - \phi_2\partial_e\phi_1)\epsilon_{abcd}dx^b \wedge dx^c \wedge dx^d$ where ϵ is completely antisymmetric and $\epsilon_{1234} = 1$, satisfies $d\eta = 0$ in view of Klein-Gordon equation. Applying Stokes-Poincaré's theorem to η for the set $B \subset D$ whose boundary is the union of V and the intersection of \bar{D} and the surface Σ at $t = 1$, one easily gets $\sigma_D(\phi_1, \phi_2) = \sigma_V(L_m^{VD}(\phi_1), L_m^{VD}(\phi_2))$. \square

We aim now to investigate the possibility of inverting, on the whole $\mathcal{S}(V)$, the injective symplectic homomorphism $L_m^{VD} : \mathcal{S}(D) \rightarrow \mathcal{S}(V)$. In other words we are concerned with the surjectivity of L_m^{VD} . Hence we have to study the problem of determining a solution of Klein-Gordon equation in $D \cup V$ when its restriction to V is assigned. That is the *Goursat problem* for the Klein-Gordon equation in $D \cup V$, using V as the *characteristic surface*. We have the following crucial result.

Theorem 3.1.3. *Fix $m^2 \in \mathbb{R}$. If $\Phi \in \mathcal{S}(V)$, there is a unique function $\phi : D \rightarrow \mathbb{R}$ such that:*

- (i) ϕ is the restriction to D of a function $C_0^\infty(\mathbb{M})$,
- (ii) $\nabla^a \nabla_a \phi - m^2 \phi = 0$,
- (iii) $u \varinjlim_V \phi = \Phi$.

Proof. By definition, if $\Phi \in \mathcal{S}(V)$, Φ/u is the restriction to V of a smooth function f_Φ defined in \mathbb{M} whose support does not intersect C so that we can always assume $\text{supp}(f_\Phi)$ under the $t = 1$ Cauchy surface. With this choice and extending Φ/u to the zero function in $\partial J^+(o) \setminus V$, working in $J^+(o)$ and then restricting to D , Theorems 5.4.1 and 5.4.2 in [Fr75] straightforwardly imply the thesis. \square

Notice that the extension of ϕ to \bar{D} is uniquely determined by continuity. Coming back to our problem to invert L_m^{VD} , we notice that the only possible candidate ϕ for satisfying $\phi \in \mathcal{S}_m(D)$ and $L_m^{VD}(\phi) = \Phi$ with a given $\Phi \in \mathcal{S}(V)$ is just the ϕ of Proposition 3.1.3. If Σ_D is a Cauchy surface of D that extends to a Cauchy surface Σ of \mathbb{M} , there is however no guarantee for having $\phi \in \mathcal{S}_m(D)$. Because the Cauchy data of ϕ on Σ may not have compact support in Σ_D , as requested in the definition of $\mathcal{S}_m(D)$. $\phi|_\Sigma$ may vanish exactly on C but not on Σ_D and $\partial_t \phi|_\Sigma$ may be strictly different from 0 on the whole $\Sigma_D \cup C$. In two dimensional models, for $m = 0$ similar counterexamples can be constructed very easily. We conclude that the injective symplectic map $L_m^{VD} : \mathcal{S}_m(D) \rightarrow \mathcal{S}(V)$ is not

surjective. It could be possible to have a surjective restriction map by working within a *larger* space of wavefunctions, where $m \geq 0$ is fixed:

$$\tilde{\mathcal{S}}_m(D) \doteq \{\phi : D \rightarrow \mathbb{R} \mid (\nabla^a \nabla_a - m^2)\phi = 0, \phi = f|_D, f \in C_0^\infty(\mathbb{M}), uf|_V \in \mathcal{S}(V)\}. \quad (19)$$

Now, the map that naturally extends L_m^{VD} :

$$L_m^{V\tilde{D}} : \tilde{\mathcal{S}}_m(D) \ni \phi \mapsto u \lim_{\rightarrow V} \phi \in \mathcal{S}(V) \quad (20)$$

is surjective in view of Proposition 3.1.3. Following this way, and using again σ_D as symplectic form on $\tilde{\mathcal{S}}_m(D)$, $(\tilde{\mathcal{S}}_m(D), \sigma_D)$ turns out to be a well defined real symplectic space with a nondegenerate symplectic form and including $(\mathcal{S}_m(D), \sigma_D)$ as a symplectic subspace through the natural inclusion map that will be indicated as:

$$L_m^{\tilde{D}D} : \mathcal{S}_m(D) \rightarrow \tilde{\mathcal{S}}_m(D). \quad (21)$$

Notice that $L_m^{VD} = L_m^{V\tilde{D}} L_m^{\tilde{D}D}$ by construction.

Remark 3.1.4. (1) The extension of $(\mathcal{S}_m(D), \sigma_D)$ to $(\tilde{\mathcal{S}}_m(D), \sigma_D)$ is canonically determined by the requirement that $\tilde{\mathcal{S}}_m(D)$ contains the solutions of the Goursat problem with data in $\mathcal{S}(V)$.

(2) What we loose in the extension of $(\mathcal{S}_m(D), \sigma_D)$ to $(\tilde{\mathcal{S}}_m(D), \sigma_D)$ is the fact that $(\tilde{\mathcal{S}}_m(D), \sigma_D)$ is not identifiable as a symplectic subspace of $(\mathcal{S}_m(\mathbb{M}), \sigma_{\mathbb{M}})$. It is because the wavefunctions of $\tilde{\mathcal{S}}_m(D)$ do not extend into a unique canonical way to smooth wavefunctions in the whole \mathbb{M} through a suitable extension of the embedding L_m^{MD} in remark 3.1.1. That is extending the Cauchy initial data on Σ_D as the zero data on $\Sigma \setminus \Sigma_D$, as we did defining L_m^{MD} . In the case of a generic element of $\tilde{\mathcal{S}}_m(D)$ this procedure gives rise to discontinuous Cauchy datum $\partial_t \phi|_{\Sigma}$ crossing $C \subset \Sigma$, generating singularities of the v -derivative of the wavefunctions, propagating along null geodesics emanating from C , when interpreting the KG equation in a suitable weaker sense. Other smooth extensions to \mathbb{M} are however possible, but they depend on the initial element $\phi \in \tilde{\mathcal{S}}_m(D)$ and on on arbitrary choices (as smoothing functions).

To go on, proving that $L_m^{V\tilde{D}} : \tilde{\mathcal{S}}_m(D) \rightarrow \mathcal{S}(V)$ is injective and surjective, we state a technical lemma, that we shall use in several occasions later. Its proof is given in the appendix and it is based on uniqueness theorems for weak solutions of Klein-Gordon equation in suitable Sobolev spaces (to prove that (22) and (23) hold if $t \geq 0$) and exploits the Hörmander's theorem about propagation of singularities (to prove that they hold for $t \leq 0$).

Lemma 3.1.5. *If $\phi \in \tilde{\mathcal{S}}_m(D)$ with $m \geq 0$, for $(t, \mathbf{x}) \in D$ it holds*

$$\phi(t, \mathbf{x}) = \frac{1}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} \frac{d\mathbf{k}}{\sqrt{2E(\mathbf{k})}} e^{i\mathbf{k}\cdot\mathbf{x} - itE(\mathbf{k})} \left(\widehat{\phi}(\mathbf{k}) + \overline{\widehat{\phi}(-\mathbf{k})} \right) \quad (22)$$

$$\partial_t \phi(t, \mathbf{x}) = \frac{-i}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} d\mathbf{k} \sqrt{2E(\mathbf{k})} e^{i\mathbf{k}\cdot\mathbf{x} - itE(\mathbf{k})} \left(\widehat{\phi}(\mathbf{k}) - \overline{\widehat{\phi}(-\mathbf{k})} \right), \quad (23)$$

where, for every fixed t , the integrals above have to be understood in the sense of the Fourier-Plancherel transform, or as standard integrals when $\phi \in \mathcal{S}_m(D)$ and where:

$$\widehat{\phi}(\mathbf{k}) \doteq \frac{e^{iE(\mathbf{k})}}{(2\pi)^{3/2}} \int_{\Sigma} \left(\sqrt{2E(\mathbf{k})} \phi(1, \mathbf{x}) + i \sqrt{\frac{2}{E(\mathbf{k})}} \partial_t \phi(1, \mathbf{x}) \right) e^{-i\mathbf{k} \cdot \mathbf{x}} d\mathbf{x}. \quad (24)$$

Above Σ is the $t = 1$ Cauchy surface of \mathbb{M} and $E(\mathbf{k}) \doteq \sqrt{m^2 + \mathbf{k}^2}$.

Proposition 3.1.6. For every fixed $m \geq 0$, the map that extends $L_m^{VD} : \mathcal{S}_m(D) \rightarrow \mathcal{S}(V)$,

$$L_m^{V\widetilde{D}} : \widetilde{\mathcal{S}}_m(D) \ni \phi \mapsto u \lim_{\rightarrow V} \phi, \quad (25)$$

is a symplectic isomorphism from $(\widetilde{\mathcal{S}}_m(D), \sigma_D)$ onto $(\mathcal{S}(V), \sigma_V)$. Moreover, if $(t, \mathbf{x}) \in D$, it holds, for every t and almost everywhere in \mathbf{x} :

$$\left((L_m^{V\widetilde{D}})^{-1} \Phi \right) (t, \mathbf{x}) = \lim_{n \rightarrow +\infty} -2 \int_{\mathbb{S}^2} d\omega \int_0^1 du \Delta_m((t - i\epsilon_n), \mathbf{x}, (u, u\omega)) u \partial_u \Phi(\omega, u), \quad (26)$$

where $\Delta_m(\mathbf{x}, \mathbf{x}')$ is the Schwartz kernel of the causal propagator of mass m Klein-Gordon equation in \mathbb{M} and $\{\epsilon_n\}_{n \in \mathbb{N}} \subset \mathbb{R}_+$ is a sequence with $\epsilon_n \rightarrow 0$ as $n \rightarrow +\infty$.

Proof. Barring surjectivity, which is assured by Theorem 3.1.3, the proof of the rest of the first statement in the thesis is exactly the same as for Proposition 3.1.2. Let us pass to the second statement for the case $m > 0$. In view of the Lemma 3.1.5, (22) makes sense, at every fixed $t \in \mathbb{R}$, interpreted as Fourier-Plancherel transforms. By standard results of L^2 -convergence, there must be a sequence of reals $\epsilon_n > 0$ with $\epsilon_n \rightarrow 0^+$ as $n \rightarrow +\infty$, such that, almost everywhere in \mathbf{x} :

$$\phi(t, \mathbf{x}) = \lim_{n \rightarrow +\infty} \frac{1}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} \frac{d\mathbf{k}}{\sqrt{2E(\mathbf{k})}} e^{i\mathbf{k} \cdot \mathbf{x} - i(t - i\epsilon_n)E(\mathbf{k})} \left(\widehat{\phi}(\mathbf{k}) + \overline{\widehat{\phi}(-\mathbf{k})} \right), \quad (27)$$

where the integral in the right-hand side is interpreted as a standard Fourier integral (this is because, if $m > 0$, $E(\pm\mathbf{k})^{-1/2} |\widehat{\phi}(\mathbf{k})| \leq \text{constant}$, as follows from (24)). The right-hand side of (24) is the integral of $-i \frac{1}{3!} \sqrt{|g|} g^{ea} \left(\frac{e^{-ik_m x^m}}{(2\pi)^{3/2} \sqrt{2E}} \partial_e \phi - \phi \partial_e \frac{e^{-ik_m x^m}}{(2\pi)^{3/2} \sqrt{2E}} \right) \epsilon_{abcd} dx^b \wedge dx^c \wedge dx^d$ over Σ_D . Dealing with as in the proof of Proposition 3.1.2, employing Stokes-Poincaré's theorem and integrating by parts, we can rearrange the right-hand side of (24), obtaining:

$$\widehat{\phi}(\mathbf{k}) = \frac{-2i}{(2\pi)^{3/2} \sqrt{2E(\mathbf{k})}} \int_{\mathbb{S}^2 \times [0,1]} d\omega du u e^{-i(\mathbf{k} \cdot u\omega - Eu)} \partial_u L_m^{V\widetilde{D}}(\phi)(\omega, u). \quad (28)$$

Inserting it in (27), after some manipulations one gets

$$\phi(t, \mathbf{x}) = \lim_{n \rightarrow +\infty} -2 \int_{\mathbb{S}^2} d\omega \int_0^1 du \Delta_m((t - i\epsilon_n), \mathbf{x}, (u, u\omega)) u \partial_u L_m^{V\widetilde{D}} \phi,$$

for every t , a.e. in \mathbf{x} where $\Delta_m(\mathbf{x}, \mathbf{x}')$ is the integral kernel of the causal propagator (the advanced-minus-retarded fundamental solution of the Klein-Gordon equation with mass

m [RS80II]), which is a smooth function when analytically extend to complex values of t . Defining $\Phi \doteq L_m^{V\tilde{D}}\phi$ (noticing that, due to the definition of $\tilde{\mathcal{S}}_m(D)$, Φ varies everywhere in $\mathcal{S}(V)$ if ϕ ranges in $\tilde{\mathcal{S}}_m(D)$), we have $\phi = (L_m^{V\tilde{D}})^{-1}\Phi$, so that the found identity is (26). In the case $m = 0$ (26) is still true. Indeed, using the explicit expression

$$\Delta_0((t - i\epsilon_n, \mathbf{x}), (t', \mathbf{x}')) = \frac{i}{4\pi^2} \left[\frac{1}{\sigma((t + i\epsilon_n, \mathbf{x}), (t', \mathbf{x}'))} - \frac{1}{\sigma((t - i\epsilon_n, \mathbf{x}), (t', \mathbf{x}'))} \right], \quad (29)$$

after some straightforward manipulations, the right-hand side of (26) turns out to coincide with the explicit solution of the Goursat problem for the massless Klein-Gordon equation with characteristic datum Φ/u as represented in (5.4.17) in [Fr75], independently from the choice of the sequence $\{\epsilon_n\}_{n \in \mathbb{N}}$. \square

3.2 Action of conformal Killing fields on massless wavefunctions

As is well-known (e.g. see [DMS97]), if ϕ satisfies Klein-Gordon equation with $m = 0$, then $\beta_\tau^Y \phi$ in (3) satisfies the same equation, provided that Y is a conformal Killing vector and $\beta_\tau^Y \phi$ is well defined (the one-parameter group generated by Y is only local in the general case). Let us specialize to the case of $Y = X$ given in (12). In that case, the flow of X leaves separately D and ∂D and V invariant, in particular, giving rise to a well-defined action $\beta_\tau^X : C^\infty(D) \rightarrow C^\infty(D)$ for every $\tau \in \mathbb{R}$. There is a similar action of β_τ^X on the functions defined on V . We want to study the interplay of these actions restricting to $\tilde{\mathcal{S}}_0(D)$ and $\mathcal{S}(V)$ respectively. To this end we notice that from the very definition (12) of X , if $(\omega_\tau^X(\omega, u), u_\tau^X(\omega, u)) \doteq \Upsilon^X(\tau, (\omega, u))$, one has

$$X|_V(\omega, u) = u(u-1)\partial_u \quad \text{so that} \quad u_\tau^X(\omega, u) = \frac{u}{u + e^\tau(1-u)}, \quad \omega_\tau^X(\omega, u) = \omega. \quad (30)$$

Referring to (4), one easily proves from the first identity in (30) that, whenever a map $f : \mathbb{S}^2 \times (0, 1) \rightarrow \mathbb{R}$ is differentiable, then $u\gamma^X f = -uX(f) - u(u-1)f$. If we define $\Phi \doteq uf$, the found identity implies that the infinitesimal generator $\gamma^X(f)$ satisfies $u\gamma^X(f) = -X(\Phi)$. Passing from the infinitesimal action to the finite one finds that transforming $f \rightarrow \beta_\tau^X(f)$ (interpreted as in (3) with $Y = X$) is equivalent to transform the associated $\Phi \doteq uf$ as $\Phi \rightarrow \beta_\tau^X(\Phi)$ where, with a harmless misuse of notation, we have defined:

$$(\beta_\tau^X(\Phi))(\omega, u) \doteq \Phi \left(\omega, \frac{u}{u + e^{-\tau}(1-u)} \right). \quad (31)$$

From now on, $\beta_\tau^X(\Phi)$ with Φ defined on V will always mean (31) in spite of the definition (3) with $Y = X$, which we will only adopt for smooth functions defined in the bulk D or in the whole \mathbb{M} . We have the following technical result.

Proposition 3.2.1. $\beta_\tau^X : \mathcal{S}(V) \rightarrow \mathcal{S}(V)$ is well-defined for every $\tau \in \mathbb{R}$, giving rise to a one-parameter group of symplectic isomorphisms of $(\mathcal{S}(V), \sigma_V)$.

Proof. Define the vector field $X' \doteq \chi X$ on \mathbb{M} , where $\chi \in C_0^\infty(\mathbb{M})$ is such that $\chi = 1$ on \overline{D} . As X' has compact support, its integral lines are complete and the associated one-parameter group of diffeomorphisms is global. Consequently the Jacobian $J_\tau^{X'}$ does not vanish anywhere and thus it is strictly positive (since it is such for $\tau = 0$). Defining $\beta_\tau^{X'}(f)$ as in (3), if $f \in C_0^\infty(\mathbb{M})$ we get a well-behaved action for every value of $\tau \in \mathbb{R}$ and $\beta_\tau^{X'}(f) \in C_0^\infty(\mathbb{M})$ for every $\tau \in \mathbb{R}$. Furthermore, in \overline{D} , the action of X' is the same as that of X . If $\Phi \in \mathcal{S}(V)$, $\Phi = uf \upharpoonright_V$ for some $f \in C_0^\infty(\mathbb{M})$. Therefore, from the discussion before Proposition 3.2.1, $\beta_\tau^X \Phi = u(\beta_\tau^{X'} f) \upharpoonright_V$ so that $\beta_\tau^X \Phi = ug \upharpoonright_V$ for $g = \beta_\tau^{X'} f \in C_0^\infty(\mathbb{M})$. Finally, notice that computing the support in V , $\text{supp}(\beta_\tau^X \Phi)$ cannot reach the points at $u = 1$ since the flow of $-X \upharpoonright_V$ transforms the support of $\beta_\tau^{X'} \Phi$ to the support of Φ , but the points at $u = 1$ are fixed for the flow of $-X$ and $\text{supp}(\Phi)$ does not include those points by definition and thus they cannot belong to $\text{supp}(\beta_\tau^X \Phi)$. To prove that β^X preserves the symplectic form σ_V , thinking of V as the product $\mathbb{S}^2 \times [0, 1)$, we notice that the symplectic form σ_V (17) is trivially invariant under the (pull back) action of diffeomorphisms of $[0, 1]$ onto $[0, 1]$ preserving the endpoints. For each $\tau \in \mathbb{R}$, the map $[0, 1] \ni u \mapsto u(u + e^\tau(1 - u))^{-1}$ is just such a diffeomorphism, so that the bijective (as $\beta_\tau^X \beta_{-\tau}^X = \beta_{-\tau}^X \beta_\tau^X = id$) linear map $\beta_\tau^X : \mathcal{S}(V) \rightarrow \mathcal{S}(V)$ in (31) preserves the symplectic form. \square

We are now in place to state the theorem concerning the interplay of the action of β^X on D and that on V for massless wavefunctions. Below, the index $_0$ means $m=0$.

Theorem 3.2.2. $\beta_\tau^X : \tilde{\mathcal{S}}_0(D) \rightarrow \tilde{\mathcal{S}}_0(D)$ is well-defined for every $\tau \in \mathbb{R}$ giving rise to a one-parameter group of symplectic isomorphisms of $(\tilde{\mathcal{S}}_0(V), \sigma_V)$. Moreover it holds

$$\beta_\tau^X \circ L_0^{V\tilde{D}} = L_0^{V\tilde{D}} \circ \beta_\tau^X \quad \text{for all } \tau \in \mathbb{R}. \quad (32)$$

Proof. Since the flow of X preserves D and it transforms massless solutions of Klein-Gordon equation to massless solutions of Klein-Gordon equation, β_τ^X transforms smooth solutions of massless Klein-Gordon equation defined in D to smooth solutions of massless Klein-Gordon equation defined in D . With the same procedure as that exploited in the proof of Proposition 3.2.1, using the vector field X' instead of X , one sees that if $f' \in C_0^\infty(\mathbb{M})$, $\beta_\tau^{X'}(f') \in C_0^\infty(\mathbb{M})$ for every fixed τ and $(\beta_\tau^{X'}(f')) \upharpoonright_D = \beta_\tau^X(f' \upharpoonright_D)$. Those results imply that, if $\phi \in \tilde{\mathcal{S}}_0(D)$, then $\beta_\tau^X \phi$ is a solution of Klein-Gordon equation in D and it is the restriction to D of some smooth compactly supported function defined on \mathbb{M} . To prove that $\beta_\tau^X \phi \in \tilde{\mathcal{S}}_0(D)$ it is enough establishing that the u -rescaled restriction of $\beta_\tau^X \phi$ to V belongs to $\mathcal{S}(V)$. By construction (since X is everywhere defined on \overline{D} and all the functions are jointly smooth) $u \lim_{\rightarrow V} \beta_\tau^X \phi = \beta_\tau^X L_0^{V\tilde{D}} \phi$, so that $\beta_\tau^X \phi \in \tilde{\mathcal{S}}_0(D)$, because $\beta_\tau^X L_0^{V\tilde{D}} \phi \in \mathcal{S}(V)$ in view of Proposition 3.2.1. The found identity can alternately be written $L_0^{V\tilde{D}} \beta_\tau^X \phi = \beta_\tau^X L_0^{V\tilde{D}} \phi$ proving (32). The fact that β_τ^X preserves the symplectic form σ_D follows from (32), Proposition 3.2.1 and the fact that $L_0^{V\tilde{D}} : \tilde{\mathcal{S}}_0(D) \rightarrow \mathcal{S}(V)$ preserves the corresponding symplectic forms as established in Theorem 3.1.6. \square

4 Algebras and states for the bulk and the boundary

This section is devoted to study the interplay of the various algebras of observables associated with the quantum Klein-Gordon field in D and restricted to V , and some relevant states. We shall prove that, remarkably, the restriction ω_m of the standard Minkowski vacuum to the bulk algebra $\mathscr{W}_m(D)$ corresponds to a unique state λ on the boundary algebra $\mathscr{W}(V)$, independently from the mass $m > 0$ of the field.

4.1 Weyl algebras on D and V and their relations

As $(S_m(\mathbb{M}), \sigma_{\mathbb{M}})$ is a real symplectic space with non-degenerate symplectic form, there is a unique (up to $*$ -algebra isomorphisms) C^* -algebra, $\mathscr{W}_m(\mathbb{M})$, with (nonvanishing) generators $W_{m,\mathbb{M}}(\phi) \neq 0$ satisfying Weyl relations [BGP07, BR96I], for every choice of $\phi_i, \phi \in S_m(\mathbb{M})$:

$$W_{m,\mathbb{M}}(\phi_1)W_{m,\mathbb{M}}(\phi_2) = e^{\frac{i\sigma_{\mathbb{M}}(\phi_1,\phi_2)}{2}}W_{m,\mathbb{M}}(\phi_1 + \phi_2), \quad W_{m,\mathbb{M}}(\phi)^* = W_{m,\mathbb{M}}(-\phi).$$

That is the the *Weyl algebra* of Klein-Gordon field on \mathbb{M} . As is well known, it can equivalently be realized in terms of generators smeared with real functions $f \in C_0^\infty(\mathbb{M})$ when defining:

$$W_{m,\mathbb{M}}[f] \doteq W_{m,\mathbb{M}}(\Delta_{m,\mathbb{M}}f) \quad \text{for all } f \in C_0^\infty(\mathbb{M}). \quad (33)$$

$\Delta_{m,\mathbb{M}} : C_0^\infty(M) \rightarrow S_m(\mathbb{M})$ is the causal propagator introduced previously. Known properties of $\Delta_{m,\mathbb{M}}$ immediately imply that locality holds

$$[W_{m,\mathbb{M}}[f], W_{m,\mathbb{M}}[h]] = 0 \quad \text{whenever } \text{supp}(f) \text{ and } \text{supp}(h) \text{ are causally separated.} \quad (34)$$

Similar Weyl algebras can be constructed for the symplectic spaces introduced above $(S_m(D), \sigma_D)$, $(\tilde{S}_m(D), \sigma_D)$ and $(S(V), \sigma_V)$. With the same order, the corresponding Weyl algebras will be denoted by $\mathscr{W}_m(D)$, $\tilde{\mathscr{W}}_m(D)$, $\mathscr{W}(V)$ with generators, respectively, $W_m(\phi)$, $\tilde{W}_m(\phi)$ and $W(\Phi)$. Since symplectic homomorphisms can be lifted to isometric $*$ -isomorphisms of the corresponding Weyl algebras [BR96II] and these $*$ -homomorphisms are $*$ -isomorphisms whenever the symplectic homomorphisms are symplectic isomorphisms, we have the following immediate result arising by the analogous properties already established for the associated symplectic homomorphisms.

Proposition 4.1.1. *For $m \geq 0$ the following holds. There is an isometric $*$ -homomorphism $\ell_m^{\mathbb{M}D} : \mathscr{W}_m(D) \rightarrow \mathscr{W}_m(\mathbb{M})$ uniquely individuated by the requirements*

$$\ell_m^{\mathbb{M}D}(W_m(\phi)) = W_{m,\mathbb{M}}(L_m^{\mathbb{M}D}\phi) \quad \text{for all } \phi \in S_m(D). \quad (35)$$

There is an isometric $$ -homomorphism $\ell_m^{\tilde{D}D} : \mathscr{W}_m(D) \rightarrow \tilde{\mathscr{W}}_m(D)$ uniquely individuated by*

$$\ell_m^{\tilde{D}D}(W_m(\phi)) = \tilde{W}_m(L_m^{\tilde{D}D}\phi) \quad \text{for all } \phi \in S_m(D). \quad (36)$$

There is an isometric $*$ -homomorphism $\ell_m^{VD} : \mathscr{W}_m(D) \rightarrow \mathscr{W}(V)$ uniquely individuated by

$$\ell_m^{VD}(W_m(\phi)) = W(L_m^{VD}\phi) \quad \text{for all } \phi \in \mathbb{S}_m(D). \quad (37)$$

There is a $*$ -isomorphism $\ell_m^{V\tilde{D}} : \widetilde{\mathscr{W}}_m(D) \rightarrow \mathscr{W}(V)$ uniquely individuated by

$$\ell_m^{V\tilde{D}}(\widetilde{W}_m(\phi)) = W(L_m^{V\tilde{D}}\phi) \quad \text{for all } \phi \in \widetilde{\mathbb{S}}_m(D). \quad (38)$$

Finally

$$\ell_m^{VD} = \ell_m^{V\tilde{D}} \ell_m^{\tilde{D}D}. \quad (39)$$

Remark 4.1.2. Similarly to the case of symplectic spaces, the Weyl algebra $\widetilde{\mathscr{W}}_m(D)$ cannot be naturally identified to a sub C^* -algebra of $\mathscr{W}_m(\mathbb{M})$. However this drawback will be remedied at the level of von Neumann algebras as we are going to prove.

4.2 States on $\mathscr{W}_m(D)$, $\mathscr{W}(S)$, $\widetilde{\mathscr{W}}_m(D)$ and their relations

Henceforth we adopt the definition of *quasifree state* as stated in [Wa94, KW91] (with the conventions of the latter concerning the signs of the symplectic forms) and assume the reader is acquainted with the basic theory of those states and their GNS Fock representations. In the following $\omega_m : \mathscr{W}_m(D) \rightarrow \mathbb{C}$ denotes the restriction to $\mathscr{W}_m(D)$ of the standard Poincaré-invariant quasifree vacuum state on $\mathscr{W}_m(\mathbb{M})$ [Ha92], obtained by standard decomposition of the $\phi \in \mathbb{S}_m(\mathbb{M})$ into positive and negative frequency parts. Taking advantage from the embedding $L_m^{VD} : \mathscr{W}_m(D) \rightarrow \mathscr{W}(V)$ (18) and (37) and viewing $\mathscr{W}_m(D)$ as a subalgebra of $\mathscr{W}(V)$, we intend to prove that $\omega_m = \lambda \circ L_m^{VD}$, for a quasifree state, independent from the value of m , $\lambda : \mathscr{W}(V) \rightarrow \mathbb{C}$ completely individuated by the requirements, if $\Phi, \Phi' \in \mathbb{S}(V)$,

$$\lambda(W(\Phi)) \doteq e^{-\frac{\mu_\lambda(\Phi, \Phi)}{2}}, \quad (40)$$

where

$$\mu_\lambda(\Phi, \Phi') \doteq \text{Re} \left(\lim_{\epsilon \rightarrow 0_+} - \int_{\mathbb{S}^2} d\omega \int_0^1 du \int_0^1 du' \frac{\Phi(\omega, u)\Phi'(\omega, u')}{\pi(u - u' - i\epsilon)^2} \right). \quad (41)$$

Proposition 4.2.1. *There exists a unique quasifree state λ on $\mathscr{W}(\mathbb{S}(V))$ satisfying (40)-(41). Its one-particle space structure $(\mathbb{K}_\lambda, \mathbb{H}_\lambda)$, which uniquely individuates the state (see Proposition 3.1 and subsequent comments in [KW91]) is the following.*

(1) $\mathbb{K}_\lambda : \mathbb{S}(V) \rightarrow \mathbb{H}_\lambda$ associates every $\Phi \in \mathbb{S}(V)$ with the corresponding $\widehat{\Phi}$:

$$\widehat{\Phi}(\omega, k) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} du e^{iku} \Phi(\omega, u) \quad \text{for every } (\omega, k) \in \mathbb{S}^2 \times \mathbb{R}_+, \quad (42)$$

where, in the integrand, we have extended Φ to the null function for $u \notin [0, 1)$;

(2) *The one-one particle space H_λ is the closure of $K_\lambda(S) + iK_\lambda(S)$ in $L^2(\mathbb{S}^2 \times \mathbb{R}_+, 2kd\omega dk)$.*

Proof. Following Sec. 3.2 in [KW91], a quasifree state is defined by the requirements (40)-(41) as stated in the thesis, if (i) μ_λ is a (real) positive-defined scalar product on $S(V)$ and (ii) $|\sigma_V(\Phi, \Phi')|^2 \leq 4\mu_\lambda(\Phi, \Phi)\mu_\lambda(\Phi', \Phi')$ for all $\Phi, \Phi' \in S(V)$. Then the one-particle structure space (K_λ, H_λ) is individuated, up to Hilbert-space isomorphisms, by $\langle K_\lambda\Phi, K_\lambda\Phi' \rangle_{H_\lambda} = \mu_\lambda(\Phi, \Phi') - i\sigma_V(\Phi, \Phi')/2$ and H_λ is the completion of $K_\lambda(S(V)) + iK_\lambda(S(V))$ with respect to the so-defined Hermitean scalar product. We want to check the validity of (i) and (ii). To this end, taking in general $\Phi, \Phi' \in S(V) + iS(V)$ we study the well-posedness of

$$\lambda(\overline{\Phi}, \Phi') \doteq \lim_{\epsilon \rightarrow 0_+} -\frac{1}{\pi} \int_{\mathbb{S}^2 \times [0,1] \times [0,1]} d\omega du du' \frac{\overline{\Phi(\omega, u)}\Phi'(\omega, u')}{(u - u' - i\epsilon)^2}, \quad \text{for } \Phi, \Phi' \in S(V) + iS(V). \quad (43)$$

Φ and Φ' are $L^2(\mathbb{S}^2 \times \mathbb{R}, d\omega du)$ functions such that, by construction, vanish uniformly in the angle ω as $u \rightarrow 0_+$, moreover the u -derivatives of Φ and Φ' belong to $L^2(\mathbb{S}^2 \times \mathbb{R}, d\omega du)$ (in particular, as one sees by direct inspection, the u -derivatives turn out to be bounded about $u = 0$ because $u^{-1}\Phi$ and $u^{-1}\Phi'$ are restrictions to V of smooth functions defined about V). Therefore the following lemma (proved in the appendix), leads to the wanted well-posedness.

Lemma 4.2.2. *Consider $\Phi, \Phi' \in L^2(\mathbb{S}^2 \times \mathbb{R}, d\omega du) \cap C^1(\mathbb{S}^2 \times (0, +\infty); \mathbb{C})$ such that they vanish for $u \in \mathbb{R} \setminus [0, 1)$ and $\omega \in \mathbb{S}^2$ and furthermore one (or both) of the following facts holds:*

- (a) $\partial_u \Phi \in L^2(\mathbb{S}^2 \times \mathbb{R}, d\omega du)$ and $\Phi(\omega, u) \rightarrow 0$ as $u \rightarrow 0_+$ uniformly in $\omega \in \mathbb{S}^2$;
- (b) $\partial_u \Phi' \in L^2(\mathbb{S}^2 \times \mathbb{R}, d\omega du)$ and $\Phi'(\omega, u) \rightarrow 0$ as $u \rightarrow 0_+$ uniformly $\omega \in \mathbb{S}^2$;

then:

$$\lim_{\epsilon \rightarrow 0_+} -\frac{1}{\pi} \int_{\mathbb{S}^2 \times [0,1] \times [0,1]} d\omega du du' \frac{\overline{\Phi(\omega, u)}\Phi''(\omega, u')}{(u - u' - i\epsilon)^2} = \int_{\mathbb{S}^2 \times \mathbb{R}_+} d\omega dk 2k \overline{\widehat{\Phi}(\omega, k)} \widehat{\Phi}'(\omega, k). \quad (44)$$

Coming back to the validity of requirements (i) and (ii) we notice that, in view of the lemma proved above, $\lambda(\overline{\Phi}, \Phi')$ is well-defined and $\lambda(\overline{\Phi}, \Phi) \geq 0$. Let us now restrict to the real case: $\Phi \in S(V)$ so that, in particular, $\mu_\lambda(\Phi, \Phi) = \text{Re}\lambda(\Phi, \Phi) = \text{Re}\lambda(\overline{\Phi}, \Phi) = \lambda(\Phi, \Phi)$. To prove (i) we only have to prove the strict positivity of μ_λ . If $\mu_\lambda(\Phi, \Phi) (= \lambda(\Phi, \Phi)) = 0$, (44) entails that $\widehat{\Phi}(\omega, k) = 0$, a.e, for $k > 0$, $\omega \in \mathbb{S}^2$. On the other hand, as $\widehat{\Phi}(\omega, -k) = \overline{\widehat{\Phi}(\omega, k)}$ because Φ is real, it must hold $\widehat{\Phi}(\omega, k) = 0$, a.e, for $k \in \mathbb{R}$, $\omega \in \mathbb{S}^2$. Since the Fourier-Plancherel transform is injective, it implies $\Phi = 0$ a.e. and thus $\Phi = 0$ everywhere as $\Phi \in S(V)$ is continuous by definition. We pass to (ii) noticing that, from (41) by trivial manipulations of the right-hand side employing integration by parts and Sokhotsky's formula:

$$-2\text{Im}\lambda(\Phi, \Phi') = \sigma_V(\Phi, \Phi') \quad \text{for all } \Phi, \Phi' \in S(V). \quad (45)$$

If (i) holds, as the Hermitean quadratic form $\lambda(\overline{\Phi}, \Phi')$ is semi-positive defined, the Cauchy-Schwartz inequality holds and so $|Im\lambda(\overline{\Phi}, \Phi')|^2 \leq |\lambda(\overline{\Phi}, \Phi')|^2 \leq \lambda(\overline{\Phi}, \Phi)\lambda(\overline{\Phi'}, \Phi')$. Restricting again to the real case $\Phi, \Phi' \in \mathcal{S}(V)$ and exploiting $\mu(\Phi, \Phi) = \lambda(\overline{\Phi}, \Phi)$ and (45), (ii) arises immediately. To conclude, we observe that in view of the definition of H_λ, μ_λ , (43), (44) and (45), one has: $\overline{K_\lambda \mathcal{S}(V)} + iK_\lambda \mathcal{S}(V) = H_\lambda$ and $\langle K_\lambda \Phi, K_\lambda \Phi' \rangle_{H_\lambda} = \mu_\lambda(\Phi, \Phi') - \frac{i}{2}\sigma_V(\Phi, \Phi')$ if $\Phi, \Phi' \in \mathcal{S}(V)$. These are the hypotheses of Proposition 3.1 in [KW91] such that (K_λ, H_λ) defines (up to unitary equivalence) the one-particle space structure of λ . \square

We are now in place to establish the main theorem of this section. The quasifree state ω_m on $\mathcal{W}_m(D)$, obtained restricting the standard Minkowski vacuum to $\mathcal{W}_m(D)$, is completely individuated by:

$$\omega_m(W_m(\phi)) \doteq e^{-\frac{\mu_m(\phi, \phi)}{2}}, \quad (46)$$

with

$$\mu_m(\phi_1, \phi_2) = Re \left(\int_{\mathbb{R}^3} d\mathbf{k} \overline{\widehat{\phi}_1(\mathbf{k})} \widehat{\phi}_2(\mathbf{k}) \right) \quad \forall \phi_1, \phi_2 \in \mathcal{S}_m(D), \quad (47)$$

and this is equivalent to say that one-particle space structure of ω_m is (K_m, H_m) where, for $\widehat{\phi}$ is defined in (24):

$$K_m \phi \doteq \widehat{\phi} \quad \text{for } \phi \in \mathcal{S}_m(D), \quad H_m \doteq \overline{K_m(\mathcal{S}_m(D))} + iK_m(\mathcal{S}_m(D)) \subset L^2(\mathbb{R}^3, d\mathbf{k}). \quad (48)$$

Theorem 4.2.3. *Let $\omega_m : \mathcal{W}_m(D) \rightarrow \mathbb{C}$ be the restriction to $\mathcal{W}_m(D)$ of the standard Poincaré invariant vacuum of the real Klein-Gordon field with mass $m \geq 0$ in \mathbb{M} . It holds:*

$$\omega_m = \lambda \circ \ell_m^{VD}. \quad (49)$$

Proof. As all the involved states are completely individuated by the real scalar products μ_λ and μ_m and ℓ_m^{VD} satisfies (37), the thesis is a consequence of the following lemma whose very technical proof is given in the appendix. \square

Lemma 4.2.4. *If $m \geq 0$, for all $\phi_1, \phi_2 \in \mathcal{S}_m(D)$ one has*

$$\int_{\mathbb{R}^3} d\mathbf{k} \overline{\widehat{\phi}_1(\mathbf{k})} \widehat{\phi}_2(\mathbf{k}) = \lim_{\epsilon \rightarrow 0^+} -\frac{1}{\pi} \int_{\mathbb{S}^2 \times [0,1] \times [0,1]} d\omega du du' \frac{L_m^{VD}(\phi_1)(\omega, u) L_m^{VD}(\phi_2)(\omega, u')}{(u - u' - i\epsilon)^2}, \quad (50)$$

where $\widehat{\phi}$ is given in (24) for every $\phi \in \mathcal{S}_m(D)$.

Remark 4.2.5. The statement of Theorem 4.2.3 resembles the uniqueness theorem proved in [KW91] for the restriction of the two-point function of a quasifree Hadamard states to bifurcate Killing horizons. However there is here an important difference that, in fact, required a different proof: The conical light surface V is not the union of the orbits of a causal Killing field.

The structure (40)-(41) of the state λ arises in several different contexts, even in QFT in curved spacetime and referring to conformal field theory and lightfront holography, where one deals with the “restrictions” of states on null surfaces or with the procedure to induce states on the bulk algebra from states defined on the boundary [KW91, Ho00, Sch03, Sch08, DMP06, DMP09, DMP09b, BMRW09, DPP10]. In the second case, very interestingly, the states induced in the bulk by the ones individuated on the boundary with the shape similar to (40)-(41) turn out to be of Hadamard class [Ho00, Mo08, DMP09b, DPP10] and are invariant under the action of natural symmetries induced by the spacetime background [DMP06, Mo06, DMP09, DPP10]. In the presently discussed case these features are automatically satisfied as the state induced on the algebra of the bulk is the very Poincaré invariant vacuum.

To conclude we pass from $\mathscr{W}_m(D)$ to the larger algebra $\widetilde{\mathscr{W}}_m(D)$ proving that it admits natural extension of ω_m that coincides to λ through the *-isomorphism $\ell_m^{\widetilde{D}M}$.

Theorem 4.2.6. *For $m > 0$, there exists a quasifree state $\widetilde{\omega}_m$ on $\widetilde{\mathscr{W}}_m(D)$ individuated by, for all $\phi_1, \phi_2 \in \widetilde{\mathcal{S}}_m(D)$:*

$$\widetilde{\omega}_m(\widetilde{W}_m(\phi)) \doteq e^{-\frac{\widetilde{\mu}_m(\phi, \phi)}{2}} \quad \text{with} \quad \widetilde{\mu}_m(\phi_1, \phi_2) = \text{Re} \left(\int_{\mathbb{R}^3} d\mathbf{k} \overline{\widehat{\phi}_1(\mathbf{k})} \widehat{\phi}_2(\mathbf{k}) \right), \quad (51)$$

where $\widehat{\phi}$ is given in (24). Moreover $\widetilde{\omega}_m$ coincides to λ through the previously-introduced *-isomorphism $\ell_m^{\widetilde{D}M} : \widetilde{\mathscr{W}}_m(D) \rightarrow \mathscr{W}(V)$ (see (38)):

$$\widetilde{\omega}_m \doteq \lambda \circ \ell_m^{\widetilde{D}M}. \quad (52)$$

Proof. Following Sec. 3.2 in [KW91], a quasifree state is defined by the requirement (51), if (i) $\widetilde{\mu}_m$ is a (real) positive-defined scalar product on $\widetilde{\mathcal{S}}_m(D)$ and (ii) $|\sigma_D(\phi, \phi')|^2 \leq 4\widetilde{\mu}_m(\phi, \phi)\widetilde{\mu}_m(\phi', \phi')$ for all $\phi, \phi' \in \widetilde{\mathcal{S}}_m(D)$. By direct inspection one sees that these are fulfilled. Moreover if $\widetilde{\mathcal{K}}_m\phi \doteq \widehat{\phi}$ given in (24) for $\phi \in \widetilde{\mathcal{S}}_m(D)$ and we define the Hilbert space $\widetilde{\mathcal{H}}_m \doteq \widetilde{\mathcal{K}}_m(\widetilde{\mathcal{S}}_m(D)) + i\widetilde{\mathcal{K}}_m(\widetilde{\mathcal{S}}_m(D))$ in $L^2(\mathbb{R}^3, d\mathbf{k})$, it holds $\langle \widetilde{\mathcal{K}}_m\phi, \widetilde{\mathcal{K}}_m\phi' \rangle_{\widetilde{\mathcal{H}}_m} = \widetilde{\mu}_m(\phi, \phi') - \frac{i}{2}\sigma_D(\phi, \phi')$ if $\phi, \phi' \in \widetilde{\mathcal{S}}_m(D)$. These are the hypotheses of Proposition 3.1 in [KW91] such that $(\widetilde{\mathcal{K}}_m, \widetilde{\mathcal{H}}_m)$ defines (up to unitary equivalence) the one-particle space structure of $\widetilde{\omega}_m$. Now the proof of (52) is the same as that of (49) since Lemma 4.2.4 is valid, with the same proof, also referring to $\phi_1, \phi_2 \in \widetilde{\mathcal{S}}_m(D)$ simply replacing L_m^{VD} for $L_m^{V\widetilde{D}}$. \square

5 The modular groups for a double cone and its boundary

In view of Reeh-Schlieder’s [Ha92] and locality properties (34) promoted to the von Neumann algebra level, the cyclic GNS vector Ψ_ω of the state ω_m is separating for the von Neumann algebra $\pi_\omega(\mathscr{W}_m(D))''$ generated by the GNS representation of $\mathscr{W}_m(D)$ in the GNS Hilbert space. In other words, $(\pi_m(\mathscr{W}_m(D))'', \Psi_m)$ is in *standard form* and thus it admits a unique modular group [Hu72, BR96II, Ha92]. We want to study it exploiting

the link between modular theory and KMS theory [Hu72, BR96II, Ha92]: If a state ω on the unital C^* -algebra \mathcal{A} is KMS with inverse temperature $\beta \neq 0$, then the pair $(\pi_\omega(\mathcal{A})'', \Psi_\omega)$ is in “standard form” and, if U_τ is the unitary implementing α_τ in the GNS representation of ω leaving Ψ_ω fixed, the modular group is $\sigma_\tau(A) = U_{-\beta\tau} A U_{-\beta\tau}^*$ for all $\tau \in \mathbb{R}$ and $A \in \pi_\omega(\mathcal{A})''$.

5.1 The modular group for the boundary V

We start our analysis building up the modular group associated with the von Neumann algebra of the GNS representation the state λ defined on the algebra of the boundary $\mathcal{W}(V)$.

Theorem 5.1.1. *Consider the quasifree state λ on $\mathcal{W}(V)$ defined in (40) and its GNS triple $(\mathcal{H}_\lambda, \pi_\lambda, \Psi_\lambda)$. Referring to (31), let $\{\alpha_\tau^X\}_{\tau \in \mathbb{R}}$ be the one-parameter group of $*$ -automorphisms of $\mathcal{W}(V)$ completely individuated by the requirement*

$$\alpha_\tau^X(W(\Phi)) = W(\beta_\tau^X(\Phi)) \quad \text{for all } \Phi \in \mathfrak{S}(V), \quad (53)$$

with β^X as in (31). The following hold.

- (a) λ is invariant and KMS with inverse temperature $\beta = 2\pi$ with respect to α^X .
- (b) The modular group of $(\pi_\lambda(\mathcal{W}(V))'', \Psi_\lambda)$ is $\mathbb{R} \ni \tau \mapsto U_{-2\pi\tau}^{(\lambda)} \cdot U_{-2\pi\tau}^{(\lambda)*}$ where $U_\tau^{(\lambda)} \Psi_\lambda = \Psi_\lambda$ if $\tau \in \mathbb{R}$ and $U^{(\lambda)}$ is the second-quantization of the one-parameter group of unitaries $V_\tau^{(\lambda)} : \mathbb{H}_\lambda \rightarrow \mathbb{H}_\lambda$,

$$V_\tau^{(\lambda)} \mathbb{K}_\lambda = \mathbb{K}_\lambda \beta_\tau^X \quad \text{for all } \tau \in \mathbb{R}. \quad (54)$$

Sketch of the proof (details in the Appendix). Our strategy is constructing a quasifree state λ' on $\mathcal{W}(V)$ which admits $(\mathbb{K}'_\lambda, \mathbb{H}'_\lambda)$ as one-particle space structure with

$$\mathbb{H}'_\lambda \doteq L^2(\mathbb{S}^2 \times \mathbb{R}; m(h) d\omega dh) \quad \text{where} \quad m(h) \doteq \frac{he^{2\pi h}}{e^{2\pi h} - e^{-2\pi h}}, \quad (55)$$

$$(\mathbb{K}'_\lambda(\Phi))(\omega, h) \doteq \int_{\mathbb{R}} dl \frac{e^{ilh}}{\sqrt{2\pi}} \Phi(\omega, \ell) \doteq \tilde{\Phi}(\omega, h), \quad (56)$$

where Φ has been represented in coordinates $(\omega, \ell) \in \mathbb{S}^2 \times \mathbb{R}$, with $u \doteq 1/(1 + e^{-\ell})$, over V . Then we prove that α^X is implemented in the bosonic Fock space \mathcal{H}'_λ by second-quantization of the strongly continuous one-parameter group of unitary operators in the one-particle space \mathbb{H}'_λ

$$V_\tau^{(\lambda')} \doteq e^{i\tau\hat{h}} \quad \text{with} \quad (\hat{h}\psi)(\omega, h) \doteq h\psi(\omega, h) \quad \text{for } (\omega, h) \in \mathbb{S}^2 \times \mathbb{R} \quad (57)$$

($\text{Dom}(\hat{h})$ is the usual one for multiplicative operators). We will establish that λ' is KMS with respect to α^X with $\beta = 2\pi$, so that the second-quantization of $V^{(\lambda')}$ (which

satisfies $V_\tau^{(\lambda')} \mathbf{K}'_\lambda = \mathbf{K}'_\lambda \beta_\tau^X$ defines the modular group of $(\pi_{\lambda'}(\mathscr{W}(V))'', \Psi_{\lambda'})$ in view of the link between KMS condition and modular group outlined at the beginning of section 5. Finally we will prove that $(\mathbf{K}'_\lambda, \mathbf{H}'_\lambda)$ is unitarily equivalent to that of λ , so that $\lambda = \lambda'$ in view of Proposition 3.1 in [KW91]. In this picture, $V_\tau^{(\lambda)}$ is the unitary corresponding to $V_\tau^{(\lambda')}$ in the structure $(\mathbf{K}_\lambda, \mathbf{H}_\lambda)$. \square

5.2 The modular group of the double cone and its relation with that of V

From Proposition 4.1.1, we know that, for every $m \geq 0$, $\mathscr{W}_m(D)$ identifies to a subalgebra of $\mathscr{W}(V)$ in view of the existence of the isometric $*$ -homomorphism ℓ_m^{VD} and that $\widetilde{\mathscr{W}}_m(D)$ identifies to the whole algebra of $\mathscr{W}(V)$ due to the $*$ -isomorphism $\ell_m^{V\widetilde{D}}$. However, it does not hold $\ell_m^{VD}(\mathscr{W}_m(D)) = \mathscr{W}(V)$ and $\ell_m^{\widetilde{D}D}(\mathscr{W}_m(D)) = \widetilde{\mathscr{W}}_m(D)$ in any cases. Conversely, the three algebras $\mathscr{W}_m(D), \widetilde{\mathscr{W}}_m(D), \mathscr{W}(V)$, in a precise sense, coincide if they are promoted to von Neumann algebras in the GNS representations of $\omega_m, \widetilde{\omega}_m$ and λ as we are going to discuss. We start with a fundamental technical lemma proved in the appendix.

Lemma 5.2.1. *For $m \geq 0$, if $\Phi \in \mathfrak{S}(V)$ and defining the corresponding $\phi \doteq (L_m^{V\widetilde{D}})^{-1}(\Phi)$, so that $\phi \in \widetilde{\mathfrak{S}}(D)$, there is a sequence $\{\phi_n\}_{n \in \mathbb{N}} \subset \mathfrak{S}_m(D)$ such that both the limits hold $\mathbf{K}_\lambda L_m^{VD}(\phi_n) \rightarrow \mathbf{K}_\lambda \Phi$ in \mathbf{H}_λ and $\widetilde{\mathbf{K}}_m L_m^{\widetilde{D}D}(\phi_n) \rightarrow \widetilde{\mathbf{K}}_m \phi$ in $\widetilde{\mathbf{H}}_m$ as $n \rightarrow +\infty$, so that, in particular:*

$$\widetilde{\mathbf{H}}_m = \mathbf{H}_m \quad \text{and} \quad \mathbf{K}_m = \widetilde{\mathbf{K}}_m L_m^{\widetilde{D}D}. \quad (58)$$

A consequence of the lemma concerns the existence of a unitary Z_m that implements the (nonsurjective) embedding $\ell_m^{VD} : \mathscr{W}_m(D) \rightarrow \mathscr{W}(V)$ at the level of the GNS representations.

Proposition 5.2.2. *For $m \geq 0$, consider the states λ on $\mathscr{W}(V)$, ω_m on $\mathscr{W}_m(D)$ and $\widetilde{\omega}_m$ on $\widetilde{\mathscr{W}}_m(D)$ with GNS representations $(\mathcal{H}_m, \pi_m, \Psi_m)$, $(\mathcal{H}_\lambda, \pi_\lambda, \Psi_\lambda)$ and $(\widetilde{\mathcal{H}}_\lambda, \widetilde{\pi}_\lambda, \widetilde{\Psi}_\lambda)$ respectively. The following holds.*

(a) *There is a unique unitary operator $Z_m : \mathcal{H}_m \rightarrow \mathcal{H}_\lambda$ such that*

$$Z_m \Psi_m = \Psi_\lambda, \quad (59)$$

$$Z_m \pi_m(a) Z_m^{-1} = \pi_\lambda(\ell_m^{VD}(a)) \quad \text{for all } a \in \mathscr{W}_m(D). \quad (60)$$

(b) *Up to unitary equivalence $(\mathcal{H}_m, \pi_m, \Psi_m) = (\widetilde{\mathcal{H}}_\lambda, \widetilde{\pi}_\lambda \circ \ell_m^{\widetilde{D}D}, \widetilde{\Psi}_\lambda)$.*

(c) *Taking (b) and (58) into account, the unitary Z_m also verifies:*

$$(i) \quad Z_m(\mathbf{H}_m) = Z_m(\widetilde{\mathbf{H}}_m) = \mathbf{H}_\lambda$$

- (ii) $Z_m \mathbf{K}_m = \mathbf{K}_\lambda L_m^{VD}$ and $Z_m \tilde{\mathbf{K}}_m = \mathbf{K}_\lambda L_m^{V\tilde{D}}$
- (iii) $Z_m \tilde{\pi}_m(a) Z_m^{-1} = \pi_\lambda(\ell_m^{V\tilde{D}}(a))$ for all $a \in \tilde{\mathcal{W}}_m(D)$.

Proof. (a) Define the \mathbb{R} -linear map $z_m : \mathbf{K}_m \phi \mapsto \mathbf{K}_\lambda L_m^{VD}(\phi)$ for every $\phi \in \mathbf{S}_m(D)$. First notice that z_m is isometric. This follows decomposing the scalar products into real and imaginary part and then using Lemma 4.2.2 for the real parts and the fact that the imaginary parts equal the corresponding symplectic forms which, in turn, are preserved by the map L_m^{VD} as discussed in the proof of Proposition 3.1.2. As a consequence of Lemma 5.2.1, the space spanned by $\mathbf{K}_\lambda L_m^{VD}(\phi) + i\mathbf{K}_\lambda L_m^{VD}(\phi)$ for $\phi \in \mathbf{S}_m(D)$ is dense in \mathbf{H}_λ . By Lemma A.2 in [KW91], we conclude that z_m extends by linearity and continuity to a unitary map from \mathbf{H}_m to \mathbf{H}_λ . In turn, by standard procedures of second quantization in bosonic Fock spaces, this map uniquely extends by linearity and continuity to a unitary map from the Fock space \mathcal{H}_m (with one-particle space structure $(\mathbf{K}_m, \mathbf{H}_m)$) onto the Fock space \mathcal{H}_λ (with one-particle space structure $(\mathbf{K}_\lambda, \mathbf{H}_\lambda)$) verifying (59) and (60). Since Ψ_m is cyclic for π_m and Z_m is unitary, we conclude from (59) and (60) that Ψ_λ has to be cyclic for $\pi_\lambda(\ell_m^{VD}(\mathcal{W}_m(D)))$. This result implies the uniqueness of any unitary $Z_m : \mathcal{H}_m \rightarrow \mathcal{H}_\lambda$ satisfying (59) and (60). The proof of the statement (b) immediately follows from (58) and the definition of $\ell_m^{V\tilde{D}}$ (and the known Fock space structure of the GNS representation for quasifree states). Concerning (c), we observe that the (i) and first identity in (ii) are true from the construction of Z_m above, also taking (58) into account. Afterwards, a unitary $Z'_m : \tilde{\mathcal{H}}_m (= \mathcal{H}_m) \rightarrow \mathcal{H}_\lambda$ verifying (iii) and $Z'_m \tilde{\Psi}_m (= Z'_m \Psi_m) = \Psi_\lambda$ can be constructed as before, thus verifying the second in (ii). In view of (39), Z'_m fulfills the same requirement satisfied (59) and (60) so it must coincide to Z_m by the uniqueness property of the latter. \square

The unitary Z_m gives a natural way to embed the representation π_m into π_λ in accordance with the (non surjective) embedding of algebras $\ell_m^{VD} : \mathcal{W}_m(D) \rightarrow \mathcal{W}(V)$. Passing to the corresponding von Neumann algebras this embedding becomes an identification that allows us to characterize the modular group of $(\pi_m(\mathcal{W}_m(D))'', \Psi_m)$ (and that of $(\tilde{\pi}_m(\tilde{\mathcal{W}}_m(D))'', \tilde{\Psi}_m)$) in terms of that of $(\pi_\lambda(\mathcal{W}(V))'', \Psi_\lambda)$.

Theorem 5.2.3. *Referring to the three von Neumann algebras $(\pi_m(\mathcal{W}_m(D))'', \Psi_m)$, $(\tilde{\pi}_m(\tilde{\mathcal{W}}_m(D))'', \tilde{\Psi}_m)$ and $(\pi_\lambda(\mathcal{W}(V))'', \Psi_\lambda)$ in standard form and to the unitary operator $Z_m : \mathcal{H}_m \rightarrow \mathcal{H}_\lambda$ introduced in Proposition 5.2.2, it holds:*

$$\pi_m(\mathcal{W}_m(D))'' = \pi_m(\tilde{\mathcal{W}}_m(D))'' = Z_m^{-1} \pi_\lambda(\mathcal{W}(V))'' Z_m. \quad (61)$$

The modular group $\{\sigma_\tau^{(m,D)}\}_{\tau \in \mathbb{R}}$ of $(\pi_m(\mathcal{W}_m(D))'', \Psi_m)$ and $(\tilde{\pi}_m(\tilde{\mathcal{W}}_m(D))'', \tilde{\Psi}_m)$ is

$$\sigma_\tau^{(m,D)}(A) = U_{-2\pi\tau}^{(m)} A U_{-2\pi\tau}^{(m)*} \quad \text{for every } A \in \pi_m(\mathcal{W}_m(D))'' \text{ and } \tau \in \mathbb{R}, \quad (62)$$

where $U_\tau^{(m)} \doteq Z_m^{-1} U_\tau^{(\lambda)} Z_m$ and $U_\tau^{(\lambda)}$ is that in (b) of Theorem 5.1.1.

Proof. The identity $\pi_m(\tilde{\mathcal{W}}_m(D))'' = Z_m^{-1} \pi_\lambda(\mathcal{W}(V))'' Z_m$ holds trivially from (iii) of (c) in Proposition 5.2.2 since Z_m is unitary and $\ell_m^{V\tilde{D}}$ is a *-isomorphism. We have to prove

the remaining identity in (61) only, since the last statement in the thesis is an immediate consequence of it, of the fact that $Z_m \Psi_m = \Psi_\lambda$ and of (b) in Theorem 5.1.1. Since, by basic properties of the commutant, $Z_m \pi_m(\mathscr{W}_m(D))'' Z_m^{-1} = (Z_m \pi_m(\mathscr{W}_m(D)) Z_m^{-1})''$, we have to prove that:

$$(Z_m \pi_m(\mathscr{W}_m(D)) Z_m^{-1})'' = \pi_\lambda(\mathscr{W}(V))'' . \quad (63)$$

We already know that $(Z_m \pi_m(\mathscr{W}_m(D)) Z_m^{-1})'' \subset \pi_\lambda(\mathscr{W}(V))''$, because $\ell_m^{VD}(\mathscr{W}_m(D)) \subset \mathscr{W}_\lambda(V)$. To pass to (63), using the fact that a von Neumann algebra is closed in the strong (operatorial) topology and that the space of finite linear combinations of elements $\pi_\lambda(W(\Phi))$, $\Phi \in \mathbf{S}(V)$, is dense in $\pi_\lambda(W(\Phi))''$ in that topology, it is enough to prove that, for every $\Phi \in \mathbf{S}(V)$, there is a sequence $\{\phi_n\}_{n \in \mathbb{N}} \subset \mathbf{S}_m(D)$ such that $\pi_\lambda(W(L_m^{VD}(\phi_n))) \rightarrow \pi_\lambda(W(\Phi))$ in the operatorial topology because $\pi_\lambda(W(L_m^{VD}(\phi_n))) = Z_m \pi_m(W_m(\phi_n)) Z_m^{-1}$ by (60). To conclude we observe that if, for a fixed $\Phi \in \mathbf{S}(V)$, the sequence $\{\phi_n\}_{n \in \mathbb{N}} \subset \mathbf{S}_m(D)$ is chosen as in Lemma 5.2.1 so that $\mathbf{K}_\lambda L_m^{VD}(\phi_n) \rightarrow \mathbf{K}_\lambda \Phi$ in \mathbf{H}_λ as $n \rightarrow +\infty$, (1) in Proposition 5.2.3 in [BR96II] implies that $\pi_\lambda(W(\ell_m^{VD}(\phi_n))) \rightarrow \pi_\lambda(W(\Phi))$, for $n \rightarrow +\infty$, in the strong operatorial topology as requested. \square

6 The spacetime action of the modular group for $m \geq 0$.

Referring to the last statement of Theorem 5.2.3, in the following $\{V_\tau^{(m)}\}_{\tau \in \mathbb{R}}$ denotes the one-parameter group of unitaries acting in the one-particle space \mathbf{H}_m whose second-quantization on the corresponding Fock space $\tilde{\mathcal{H}}_m$ (assuming the cyclic vector Ψ_m to be invariant) is the unitary $U^{(m)}$ implementing the modular group of $(\tilde{\pi}_m(\tilde{\mathscr{W}}_m(D))'', \Psi_m)$. We aim to study the action of the modular group *on the wavefunctions in the spacetime*, for $m \geq 0$. In other words, we are interested in the existence of a one-parameter group $\{s_\tau^{(m)}\}_{\tau \in \mathbb{R}}$ of symplectic isomorphisms of $\tilde{\mathbf{S}}_m(D)$ such that:

$$V_\tau^{(m)} \tilde{\mathbf{K}}_m \phi = \tilde{\mathbf{K}}_m s_\tau^{(m)}(\phi) \quad \text{for every } \phi \in \tilde{\mathbf{S}}_m(D) \text{ and } \tau \in \mathbb{R}. \quad (64)$$

If $\{s_\tau^{(m)}\}_{\tau \in \mathbb{R}}$ exists, it is unique because $\tilde{\mathbf{K}}_m$ in (64) is injective. From (54), the second identity in (ii) in Proposition 5.2.2 and the last statement of the thesis of Theorem 5.2.3, we have $V_\tau^{(m)} \tilde{\mathbf{K}}_m = \tilde{\mathbf{K}}_m (L_m^{V\tilde{D}})^{-1} \beta_\tau^X L_m^{V\tilde{D}}$. So that it must be:

$$s_\tau^{(m)} \doteq (L_m^{V\tilde{D}})^{-1} \beta_\tau^X L_m^{V\tilde{D}} \quad (65)$$

Notice that the right-hand side is well defined, it is a symplectic isomorphisms of $\tilde{\mathbf{S}}_m(D)$ for every τ , and (64) holds as a consequence of the given definition.

A first result is obtained taking Theorem 3.2.2 into account for the case $m = 0$: Since it holds $(L_0^{V\tilde{D}})^{-1} \beta_\tau^X = \beta_\tau^X (L_0^{V\tilde{D}})^{-1}$, inserting it in the right-hand side of (65) with $m = 0$, we obtain:

$$s_\tau^{(0)} = \beta_\tau^X . \quad (66)$$

This is nothing but the known result obtained in [HL95] written in our language.

Remark 6.0.4. It should be clear that we are working in the space $\tilde{\mathfrak{S}}_m(D)$ rather than $\mathfrak{S}_m(D)$ because $L_m^{VD} : \mathfrak{S}_m(D) \rightarrow \mathfrak{S}(V)$ is not surjective and so $(L_m^{VD})^{-1}$, that would appear in the corresponding of (65), would not exist. This is the ultimate reason to enlarge $\mathfrak{S}_m(D)$ to the space $\tilde{\mathfrak{S}}_m(D)$. This enlargement makes sense since it does not affect the relevant von Neumann algebra and the modular group as already established in Theorem 5.2.3.

6.1 The relative action of the modular group and its infinitesimal generator

We pass to study the *relative action* of the modular operator on wavefunctions, that is:

$$s_\tau^{(m)}\phi - s_\tau^{(0)}\phi \quad \text{for } \phi \in \tilde{\mathfrak{S}}_m(D), \quad (67)$$

giving its explicit expression. Three preliminary remarks are necessary. (1) $s_\tau^{(0)}\phi$ is well defined for every $\phi \in \tilde{\mathfrak{S}}_m(D)$ even if $m > 0$ since it coincides to $\beta_\tau^X\phi$ as remarked above. However, differently from $s_\tau^{(m)}\phi$, one may have $s_\tau^{(0)}\phi \notin \tilde{\mathfrak{S}}_m(D)$ if $m > 0$ and $\phi \in \tilde{\mathfrak{S}}_m(D)$. (2) Every map $L_m^{V\tilde{D}} : \tilde{\mathfrak{S}}_m(D) \rightarrow \mathfrak{S}(V)$, for all values of m , can be seen as the restriction to the corresponding space $\tilde{\mathfrak{S}}_m(D)$ of a common map $L : C_0^\infty(\mathbb{M}) \rightarrow C^\infty(\mathbb{S}^2 \times (0, 1))$ associating a smooth function f to the function $(Lf)(\omega, u) \doteq uf|_V(\omega, u)$. To do it, every element $\phi \in \tilde{\mathfrak{S}}_m(D)$ has to be smoothly extended to the whole \mathbb{M} (as is allowed by the very definition of $\tilde{\mathfrak{S}}_m(D)$) and the function $L\phi$ does not depend on the extension. (3) if $\phi \in \tilde{\mathfrak{S}}_{m'}(D)$ then $L\phi \in \mathfrak{S}(V)$ so that $(L_m^{V\tilde{D}})^{-1}(L\phi)$ is however well-defined regardless if $m \neq m'$. More generally, $(L_m^{V\tilde{D}})^{-1}(Lf)$ is well defined for $f \in C_0^\infty(\mathbb{M})$ provided that $Lf \in \mathfrak{S}(V)$.

Remark 6.1.1. From now on we simplify the notation as $L_m \doteq L_m^{V\tilde{D}}$.

Theorem 6.1.2. *If $m \geq 0$, $\phi \in \tilde{\mathfrak{S}}_m(D)$ and $(t, \mathbf{x}) \in D$ then:*

$$s_\tau^{(m)}\phi - s_\tau^{(0)}\phi = [L_m^{-1}L - L_0^{-1}L, \beta_\tau^X]\phi \quad \text{where, if } f \in C_0^\infty(\mathbb{M}) \text{ and } Lf \in \mathfrak{S}(V) \quad (68)$$

$$((L_m^{-1}L - L_0^{-1}L)f)(t, \mathbf{x}) = \int_{\mathbb{S}^2} \int_0^{u^*(t, \mathbf{x}, \omega)} F_m(\sigma((t, \mathbf{x}), (u, u\omega))) u \partial_u u f|_V(\omega, u) du d\omega \quad (69)$$

$$\text{with} \quad u^*(t, \mathbf{x}, \omega) \doteq \frac{t^2 - \mathbf{x}^2}{2(t - \omega \cdot \mathbf{x})}, \quad F_m(z) \doteq \frac{m^2}{8\pi} \sum_{k=0}^{+\infty} \binom{m}{2} \frac{z^k}{k!(k+1)!}, \quad z \in \mathbb{R} \quad (70)$$

Remark 6.1.3. For fixed $(t, \mathbf{x}) \in D$ and $\omega \in \mathbb{S}^2$, the real $u^* \doteq u^*(t, \mathbf{x}, \omega)$ is the unique solution of $\sigma((t, \mathbf{x}), (u^*, u^*\omega)) = 0$, so that $0 < u^* < 1$. The function F_m , due to the rapid convergence of the series, is analytic it being the restriction to \mathbb{R} of an entire analytic function. F_m vanishes, as expected, for $m = 0$.

Proof of Theorem 6.1.2. By definition of $s_\tau^{(m)}$ and it being $L_m^{-1}L_m = I_{\tilde{\mathcal{S}}_m(D)}$, we immediately obtain $s_\tau^{(m)}\phi - s_\tau^{(0)}\phi = L_m^{-1}\beta_\tau^X L_m\phi - \beta_\tau^X L_m^{-1}L_m\phi$. That is $s_\tau^{(m)}\phi - s_\tau^{(0)}\phi = L_m^{-1}\beta_\tau^X L\phi - \beta_\tau^X L_m^{-1}L\phi$. Now notice that $\beta_\tau^X Lf = L\beta_\tau^X f$, essentially because X is tangent to V and L is the restriction to V (however it can be checked by direct inspection). We have found that $s_\tau^{(m)}\phi - s_\tau^{(0)}\phi = L_m^{-1}L\beta_\tau^X\phi - \beta_\tau^X L_m^{-1}L\phi$, where now, in general, $L_m^{-1}L\beta_\tau^X\phi \neq L_m^{-1}L_m\beta_\tau^X\phi = \beta_\tau^X\phi$ since $\beta_\tau^X\phi \notin \tilde{\mathcal{S}}_m(D)$ if $m > 0$. We can re-arrange the found result as $s_\tau^{(m)}\phi - s_\tau^{(0)}\phi = (L_m^{-1}L - L_0^{-1}L)\beta_\tau^X\phi - (\beta_\tau^X L_m^{-1}L - L_0^{-1}L\beta_\tau^X)\phi$. The found identity gives rise to (68) if observing that $L_0^{-1}L\beta_\tau^X = L_0^{-1}\beta_\tau^X L = \beta_\tau^X L_0^{-1}L$ where, in the last passage, we have used Theorem 3.2.2. Let us pass to (69) exploiting (26). If $f \in C_0^\infty(\mathbb{M})$ and $Lf \in \mathcal{S}(V)$:

$$\begin{aligned} ((L_m^{-1}L - L_0^{-1}L)f)(t, \mathbf{x}) = \\ \lim_{n \rightarrow +\infty} -2 \int_{\mathbb{S}^2} \int_0^1 (\Delta_m - \Delta_0)((t - i\epsilon_n, \mathbf{x}), (u, u\omega)) u \partial_u (Lf)(\omega, u) du d\omega . \end{aligned}$$

The integral defining the advanced and retarded fundamental solutions of Klein-Gordon equation [RS80II] produces, for $m > 0$:

$$\begin{aligned} \Delta_m((t - i\epsilon_n, \mathbf{x}), (t', \mathbf{x}')) = \\ \frac{im}{4\pi^2} \left[\frac{K_1(m\sigma((t + i\epsilon_n, \mathbf{x}), (t', \mathbf{x}'))^{1/2})}{\sigma((t + i\epsilon_n, \mathbf{x}), (t', \mathbf{x}'))^{1/2}} - \frac{K_1(m\sigma((t - i\epsilon_n, \mathbf{x}), (t', \mathbf{x}'))^{1/2})}{\sigma((t - i\epsilon_n, \mathbf{x}), (t', \mathbf{x}'))^{1/2}} \right] , \end{aligned}$$

the analog for $m = 0$ appears in (29). Using the standard expansion of the Bessel functions K_1 and I_1 (see 8.446 and 2 of 8.447 in [GR95]), one sees that, at $(t, \mathbf{x}) \in D$ fixed, $-2(\Delta_m - \Delta_0)((t - i\epsilon_n, \mathbf{x}), (u, u\omega))$ is the smooth function $F_m(\sigma((t, \mathbf{x}), (u\omega, u)))$ multiplied with a function weakly tending to $\theta(-\sigma((t, \mathbf{x}), (u, u\omega)))$ as $n \rightarrow +\infty$, where $\theta(x) = 0$ for $x < 0$ and $\theta(x) = 1$ otherwise. This restricts the final integration to $\omega \in \mathbb{S}^2$ and $u \in [0, u^*(t, \mathbf{x}, \omega)]$ giving rise to (69) almost everywhere in \mathbf{x} . Since both members of (69) are continuous in (t, \mathbf{x}) , the identity holds everywhere. \square

If $\phi \in \tilde{\mathcal{S}}_m(D)$, the function $(\tau, t, \mathbf{x}) \mapsto (s_\tau^{(m)}\phi)(t, \mathbf{x})$ is jointly smooth for $\tau \in \mathbb{R}$ and $(t, \mathbf{x}) \in D$: The first and, for Theorem 6.1.2, the third term of the sum in the right-hand side of

$$s_\tau^{(m)}\phi = s_\tau^{(0)}\phi + (L_m^{-1}L - L_0^{-1}L)\beta_\tau^X\phi - \beta_\tau^X(L_m^{-1}L - L_0^{-1}L)\phi$$

are smooth by construction, the second one is smooth as it can be trivially proved by recursively deriving the integral in (69) using Leibniz rule and the Theorem of derivation under the symbol of integration. So, it makes sense to compute the *infinitesimal generator* $\delta^{(m)}$ of $s_\tau^{(m)}$, taking the derivative of $s_\tau^{(m)}\phi$ at $\tau = 0$. If γ^X is defined as in (4) with Y replaced for X , (68) leads to:

$$\delta^{(m)}\phi = \gamma^X\phi + [L_m^{-1}L - L_0^{-1}L, \gamma^X]\phi \quad \text{for every } \phi \in \tilde{\mathcal{S}}_m(D), m \geq 0. \quad (71)$$

The right-hand side of (71) can explicitly be computed taking (69) into account. Notice that, for $m = 0$, (71) produces, as it has to do,

$$\delta^{(0)}\phi = \gamma^X \phi. \quad (72)$$

Theorem 6.1.4. *For every $\phi \in \tilde{\mathcal{S}}_m(D)$, $m \geq 0$, it holds:*

$$(\delta^{(m)}\phi)(t, \mathbf{x}) = (\gamma^X \phi)(t, \mathbf{x}) + \int_{\mathbb{S}^2} \int_0^{u^*(t, \mathbf{x}, \omega)} I_\phi(t, \mathbf{x}, u, \omega, \sigma) \, du d\omega \quad (73)$$

where

$$I_\phi(t, \mathbf{x}, u, \omega, \sigma) \doteq u[2 - (t + u)] [F_m(\sigma) + \sigma F'_m(\sigma)] \partial_u u \phi|_V(\omega, u), \quad (74)$$

with $\sigma \doteq \sigma((t, \mathbf{x}), (u\omega, u))$ and $(t, \mathbf{x}) \in D$.

Proof. The commutator in the right-hand side of (71) can be explicitly computed using (69) and the identity (2.1) for $Y = X$, that permits to pass the symbol γ^X from the variables (t, \mathbf{x}) to the variables $(u, u\omega)$ in V cancelling some terms. Though the computation is long, it is straightforward and produces just the result (73) up to a further added term to the right-hand side:

$$- \int_{\mathbb{S}^2} d\omega F_m(0) \{u[u(1-u) + X^a(t, \mathbf{x})\partial_a u^*(t, \mathbf{x}, \omega)] \partial_u u \phi|_V(\omega, u)\}_{u=u^*(t, \mathbf{x}, \omega)}. \quad (75)$$

This term vanishes for the following reason. By the definition of $u^*(x, \omega)$ (see remark 6.1.3), $\sigma(x, (u^*(x, \omega), u^*(x, \omega)\omega)) = 0$ for all $x \equiv (t, \mathbf{x}) \in D$. Therefore we also have that, applying X to the complete dependence on x , $X^a(x) \frac{\partial}{\partial x^a} \sigma|_{u=u^*} + X^a(x) \frac{\partial u^*}{\partial x^a} \frac{\partial}{\partial u} \sigma|_{u=u^*} = 0$. Making use of (6), it implies $(u(1-u)\partial_u \sigma) + ((t-1) + (u-1))\sigma|_{u=u^*} + X^a(x) \frac{\partial u^*}{\partial x^a} \frac{\partial}{\partial u} \sigma|_{u=u^*} = 0$. Since $\sigma|_{u=u^*} = 0$, and $\sigma = -t^2 + \mathbf{x}^2 + 2u(t - \mathbf{x} \cdot \omega)$, the found identity is nothing but: $[u(1-u) + X^a(t, \mathbf{x})\partial_a u^*(t, \mathbf{x}, \omega)]|_{u=u^*}(t - \mathbf{x} \cdot \omega) = 0$. However $t - \mathbf{x} \cdot \omega > 0$ because $t > \|\mathbf{x}\|$ in D and $\|\omega\| = 1$. Therefore $[u(1-u) + X^a(t, \mathbf{x})\partial_a u^*(t, \mathbf{x}, \omega)]|_{u=u^*} = 0$ and the integral in (75) vanishes in any case. \square

6.2 $\delta^{(m)} - \delta^{(0)}$ is a pseudo-differential operator of class $L_{1,1}^0$

We consider the difference of generators $\delta^{(m)} - \delta^{(0)}$ (that is $\delta^{(m)} - \gamma^X$) as an operator $C_0^\infty(\mathbb{M}) \rightarrow C^\infty(D)$ in view of the fact that the integral in the right-hand side of (73) is well-defined also for ϕ replaced with a generic $f \in C_0^\infty(\mathbb{M})$, though its physical meaning is guaranteed only when it acts on $\phi \in \tilde{\mathcal{S}}_m(D)$. Hence we consider, for $f \in C_0^\infty(\mathbb{M})$ and $(t, \mathbf{x}) \in D$:

$$((\delta^{(m)} - \delta^{(0)})f)(t, \mathbf{x}) \doteq \int_{\mathbb{S}^2} d\omega \int_0^{u^*(t, \mathbf{x}, \omega)} du I_f(t, \mathbf{x}, u, \omega, \sigma), \quad (76)$$

with $I_f(t, \mathbf{x}, u, \omega, \sigma)$ as in (74).

We pass to prove that $\delta^{(m)} - \delta^{(0)}$ is a *pseudodifferential operator* of class $L_{1,1}^0(D)$, proving and making more precise Fredenhagen's conjecture mentioned in the introduction.

Remark 6.2.1. We address the reader to specialized texts (e.g. [Ta96],[GS94],[Du73]) for more information on this subject while recalling just two relevant facts here.

(1) A *pseudodifferential operator* $A : C_0^\infty(\mathbb{R}^n) \rightarrow C^\infty(X)$ – where $X \subset \mathbb{R}^n$ is open – is individuated by its *amplitude*, a smooth function on $a : X \times \mathbb{R}^n \rightarrow \mathbb{C}$, which satisfies

$$\left| \partial_x^\alpha \partial_\theta^\beta a(x, k) \right| \leq C_{\alpha, \beta, K}(a) (1 + \langle k, k \rangle^{1/2})^{m - \rho|\beta| + \delta|\alpha|} \quad (77)$$

for every compact set $K \subset X$, every pair of multindices $\alpha, \beta \in \mathbb{N}^n$ and corresponding constants $C_{\alpha, \beta, K}(a) \geq 0$. Above $|\alpha| \doteq \alpha_1 + \dots + \alpha_n$ when $\alpha = (\alpha_1, \dots, \alpha_n)$ and $\langle \cdot, \cdot \rangle$ denotes the standard scalar product in \mathbb{R}^n . Then, by definition, the *pseudodifferential operator* associated to a is

$$(Au)(x) \doteq \int_{\mathbb{R}^n} a(x, k) e^{i\langle \theta, x \rangle} \hat{u}(k) dk \quad \text{if } u \in C_0^\infty(\mathbb{R}^n), \quad (78)$$

where \hat{u} indicates the Fourier transform of u . Holding (77), we say that *the amplitude* a belongs to the functional class $S_{\rho, \delta}^m(X \times \mathbb{R}^n)$ and, equivalently, that *the pseudodifferential operator* A belongs to the class $L_{\rho, \delta}^m(X)$.

(2) Compositions of pseudodifferential operators can be made in various fashions. However we only stick to the following simple result (see, for instance, Proposition 3.3 in [Ta96]). Given $A_i \in L_{\rho_i, \delta_i}^{m_i}(X)$ $i = 1, 2$, suppose that $0 \leq \delta_2 < \rho \leq 1$, where $\rho = \min\{\rho_1, \rho_2\}$, then $A_1 A_2$ is well defined and belongs to $L_{\rho, \delta}^{m_1 + m_2}(X)$ where $\delta = \max\{\delta_1, \delta_2\}$.

To prove that $\delta^{(m)} - \delta^{(0)}$ is a *pseudodifferential operator* of class $L_{1,1}^0(D)$, it would be sufficient to recast (76) into the form (78) where, of course $a \in S_{\rho, \delta}^m(X \times \mathbb{R}^n)$ with $m = 0$ and $\rho = \delta = 1$. We will indirectly obtain that result exploiting the composition rule mentioned in (2) in Remark 6.2.1. We preventively need two lemmata concerning the restriction of a smooth function on V and the form of $\delta^{(m)} - \delta^{(0)}$ as integral operator. Henceforth we equip V with the topology induced by \mathbb{M} that, in turn, coincides with that of the identification $V \equiv \mathbb{S}^2 \times [0, 1)$, and the associated Borel σ -algebra.

Lemma 6.2.2. *If $g : V \rightarrow \mathbb{R}$ is measurable and bounded and $f \in C_0^\infty(\mathbb{M})$ then:*

$$\int_V gf \upharpoonright_V dud\omega = \frac{1}{(2\pi)^4} \int_{\mathbb{R}^4} dk \int_V e^{i(k_0 u + \mathbf{k} \cdot u\omega)} g(u, \omega) dud\omega \int_{\mathbb{R}^4} e^{-i\langle k, y \rangle} f(y) dy \quad (79)$$

where $k = (k_0, \mathbf{k})$, $y = (t', \mathbf{x}')$ and $\langle k, y \rangle \doteq k_0 t' + \mathbf{k} \cdot \mathbf{x}'$.

Remark 6.2.3. The map $\kappa : \mathbb{S}^2 \times (0, 1) \ni (\omega, u) \mapsto (u, u\omega) \in \mathbb{M}$ is a smooth embedding into \mathbb{M} of a cylinder with image given by the embedded submanifold $V^* \doteq V \setminus \{o\}$ of \mathbb{M} . This embedding, in fact, gives rise to a *blow up* of the singularity at the tip o of V . Since o has zero measure with respect to $dud\omega$, removing the tip of the cone does not affect the integration on V and it permits to exploit standard results about restrictions of smooth functions to embedded manifolds V^* .

Proof of Lemma 6.2.2. The pull back action of κ on $f \in C_0^\infty(\mathbb{M})$, interpreted as a restriction $f \mapsto f|_{V^*}$ can be represented as a *Fourier integral operator* [Du73]. In practice, if $g \in C_0^\infty(\mathbb{S}^2 \times (0, 1))$ we have

$$\int_{V^*} gf|_V dud\omega = \frac{1}{(2\pi)^4} \int dk d\omega dudx e^{-i\langle k, x - (u, u\omega) \rangle} g(u, \omega) f(x)$$

where the right-hand side has to be interpreted as an *oscillatory integral* [Du73]: Inserting a further smooth compactly supported function $\chi(\epsilon k)$ in the integrand, with $\chi = 1$ in a neighborhood of $k = 0$, and taking the limit $\epsilon \rightarrow 0^+$. Rearranging the integrals, the limit produces the right-hand side of (79), using the fact that $\int_{\mathbb{R}^4} e^{-i\langle k, y \rangle} f(y) dy$ belongs to Schwartz space in the variable k , that $\int_{V^*} e^{i(k_0 u + \mathbf{k} \cdot u\omega)} g(u, \omega) dud\omega$ is bounded in the variable k , and exploiting Lebesgue's dominated convergence theorem. So (79) holds true if $g \in C_0^\infty(\mathbb{S}^2 \times (0, 1))$ and $f \in C_0^\infty(\mathbb{M})$. By a standard corollary of Luzin's theorem, if the function $g : \mathbb{S}^2 \times (0, 1) \rightarrow \mathbb{C}$ is bounded by $M < +\infty$ and is Borel-measurable, there is a sequence of compactly-supported continuous functions g_n , all bounded by M , with $g_n \rightarrow g$ almost everywhere, for $n \rightarrow +\infty$, with respect the natural Borel measure on $\mathbb{S}^2 \times (0, 1)$. Stone-Weierstrass theorem implies that the g_n can be chosen in $C_0^\infty(\mathbb{S}^2 \times (0, 1))$. Using the sequence of these g_n , (79) extends to the case of g bounded and measurable, exploiting Lebesgue's dominated convergence theorem twice, the fact that $\int_{\mathbb{R}^4} e^{-i\langle k, y \rangle} f(y) dy$ belongs to Schwartz space in the variable k , the uniform bound M for g and the g_n and the further bound of $\int_{V^*} e^{i(k_0 u + \mathbf{k} \cdot u\omega)} g(u, \omega) dud\omega$ and all the $\int_{V^*} e^{i(k_0 u + \mathbf{k} \cdot u\omega)} g_n(u, \omega) dud\omega$ by $4\pi M$. We have proved that (79) holds with V replaced for V^* in both sides. However, since o has zero measure with respect to $dud\omega$, the integration over V^* can be extended to V without affecting the final result. \square

Lemma 6.2.4. *The operator $\delta^{(m)} - \delta^{(0)} : C_0^\infty(\mathbb{M}) \rightarrow C^\infty(D)$ in (76) can be written as*

$$((\delta^{(m)} - \delta^{(0)})f)(x) = \int_{\mathbb{R}^4} dk e^{i\langle k, x \rangle} b(x, k) \widehat{Z(f)}(k) \quad \text{for } f \in C_0^\infty(\mathbb{M}) \quad (80)$$

with $x = (t, \mathbf{x})$, $k = (k_0, \mathbf{k})$ and (with notation as in Theorem 6.1.2):

$$b(x, k) \doteq e^{-i\langle k, x \rangle} \int_{\mathbb{S}^2 \times [0, 1]} d\omega du e^{i(k_0 u + \mathbf{k} \cdot u\omega)} g(x, \omega, u), \quad (81)$$

$$g(x, \omega, u) \doteq (2\pi)^{-4} \theta(u^*(t, \mathbf{x}, \omega) - u) u [2 - (t + u)] [F_m(\sigma) + \sigma F'_m(\sigma)], \quad (82)$$

$$Z \doteq 1 + t\partial_t + \mathbf{x} \cdot \nabla_{\mathbf{x}}. \quad (83)$$

Furthermore $b \in S_{1,1}^{-1}(D \times \mathbb{R}^4)$.

Proof. By direct inspection, one sees that $Z(f)|_{V^*} = \partial_u u(f|_{V^*})$. Next, for every $x \in D$, one has

$$((\delta^{(m)} - \delta^{(0)})f)(x) = \int_{\mathbb{R}^4} dk \int_{\mathbb{R}^4} dy e^{i\langle k, x - y \rangle} b(x, k) Z(f)(y) \quad \text{for } f \in C_0^\infty(\mathbb{M}) \quad (84)$$

by (79) in the case of $g = g(x, \omega, u)$ as in (82), if taking (76) into account and where $b(x, k)$ is defined in (81). The right-hand side of (84) can equivalently be written as (80).

Let us prove that $b \in S_{11}^{-1}(D)$. By definition $b \in C^\infty(D \times \mathbb{R}^4)$, as one can easily check. Furthermore

$$|b(x, k)| = \left| \int_{\mathbb{S}^2 \times [0,1]} d\omega du e^{i(k_0 u + u \mathbf{k} \cdot \omega)} u g_0(x, \omega, u) \right| ,$$

where $|g_0(x, \omega, u)| \leq M_K < +\infty$ for $(x, \omega, u) \in K \times \mathbb{S}^2 \times [0, 1]$ for every compact $K \subset D$, and the derivatives of g_0 , which exists almost everywhere, are similarly bounded. Using polar coordinates with polar axis z along \mathbf{k} , to evaluate the integral above, integrating by parts one sees that $\|\mathbf{k}\| |b(x, k)| \leq C'_K < +\infty$ if $k \in \mathbb{R}^4$ and $x \in K$. Similarly, reducing to the previous case by integrating by parts again, one finds $|k_0| |b(x, k)| \leq C''_K < +\infty$ if $k \in \mathbb{R}^4$ and $x \in K$. Thus

$$|b(x, k)| \leq C_K (1 + \langle k, k \rangle^{1/2})^{-1} .$$

Let us improve this estimate passing to consider the derivatives. By Leibniz and the differentiation rules under the integral sign, we get that $\partial_{x_i}^{\alpha_i} b(x, k)$ with $\alpha_i \geq 1$ can be decomposed as a sum of terms as follows

$$\partial_{x_i}^{\alpha_i} b(x, k) = \int_{\mathbb{S}^2} d\omega \{ I_1(x, k, u^*, \omega) + I_2(x, k, u^*, \omega) + I_3(x, k, u^*, \omega) \} \quad (85)$$

where, respectively,

$$\begin{aligned} I_1(x, k, u^*, \omega) &= \sum_{\gamma_i=0}^{\alpha_i-1} c_{\alpha_i, \gamma_i} \partial_{x_i}^{\gamma_i} I(x, k, u^*, \omega) \partial_{x_i}^{\alpha_i - \gamma_i} u^*(x, \omega) , \\ I_2(x, k, u^*, \omega) &= \sum_{\lambda_i=1}^{\alpha_i-1} \sum_{\mu_i=0}^{\alpha_i - \lambda_i - 1} d_{\lambda_i, \mu_i, \alpha_i} \partial_{x_i}^{\mu_i} \left[\left(\partial_{x_i}^{\lambda_i} I(x, u, k, \omega) \right) \Big|_{u=u^*} \right] \partial_{x_i}^{\alpha_i - \lambda_i - \mu_i} u^*(x, \omega) , \\ I_3(x, k, u^*, \omega) &= \int_0^{u^*(x, \omega)} du \partial_{x_i}^{\alpha_i} I(x, k, u, \omega) , \end{aligned}$$

where, $c_{\alpha_i, \gamma_i}, d_{\lambda_i, \mu_i, \alpha_i}$ are some numerical coefficients, and finally, $I(x, k, u, \omega)$ represents the integrand function in (81).

By a direct inspection of the formulas it is clear that the dominant term w.r.t. the variable k , for $\langle k, k \rangle^{1/2} \geq 1$, comes from the term in I_1 with the highest power of derivatives of the function $I(x, u^*, k, \omega)$. Indeed, the derivatives act on the exponential term because of the u^* -dependence on the space variables, dropping down terms depending solely on the variable k . By differentiating w.r.t. all variables $x_i, i = 1, \dots, 4$, one gets a polynomial in k of order $|\alpha| - 1$ with smooth coefficients. Differentiating w.r.t $k_i, i = 1, \dots, 4, \beta_i$ -times, respectively, gives a polynomial in k of order $|\alpha| - |\beta| - 1$, with smooth coefficients as well. Hence we get, considering $K \subset D$ compact, that there exists $C_{\alpha, \beta, K}(a) \geq 0$ for which

$$\left| \partial_k^\beta \partial_x^\alpha b(x, k) \right| \leq C_{\alpha, \beta, K}(a) (1 + \langle k, k \rangle^{1/2})^{-1 + |\alpha| - |\beta|} \quad (86)$$

holds true, so that $b \in S_{1,1}^{-1}(D)$ by definition. \square

We are in place to establish the main result of this section proving and going beyond Fredenhagen's conjecture.

Theorem 6.2.5. *The operator $\delta^{(m)} - \delta^{(0)}$ is a pseudodifferential operator of class $L_{1,1}^0(D)$.*

Proof. It holds $Z \in L_{1,0}^1(D)$. As shown above $b \in S_{1,1}^{-1}(D \times \mathbb{R}^4)$, hence by (84) and the recalled composition property of pseudodifferential operators ((2) in Remark 6.2.1), we have that the operator in the right-hand side of (84) is a pseudodifferential operator of the class $L_{1,1}^0(D)$, since in our case $m_1 = -1$ $m_2 = 1$ and $\rho_1 = \rho_2 = 1$ and $\delta_1 = 1$ $\delta_2 = 0$. \square

7 Conclusions and Outlook

We have shown how to build up, in the symplectic space of KG solutions, the analytical form of the one-parameter group of symplectic isomorphisms corresponding to the modular automorphism group of von Neumann algebra of a free massive scalar theory in the vacuum of Minkowski spacetime localized in double cones. The same procedure has lead to an explicit formula for the generator of that group as well. Many different kind of techniques contributed to achieve this goal: Mainly ideas coming from the correspondence between the Cauchy and Goursat problems for Klein-Gordon partial differential operators, some ideas coming from the algebraic approach to quantum field theory, even in curved spacetime, relating theories on different regions of spacetime and ideas coming from the microlocal approach to distributions.

A remarkable result concerns the fact that the theory in the bulk and that on the boundary are unitarily equivalent when promoted to von Neumann algebras theories so that, the modular groups are identified through the same unitary intertwiner. In particular, referring to the theory defined on the boundary, the modular group preserves a conformal-geometric meaning regardless the presence of a mass for the theory in the bulk where, conversely, any geometric meaning of the modular group is absent (for nonvanishing mass).

One of the main results has however been the classification of the difference of the generators of the modular groups for the massive and massless cases as a pseudodifferential operator acting on the solution space of the Klein-Gordon equation. This achievement goes much beyond the conjecture stated in the literature as soon as it makes the statement more precise even concerning the type of the operator, as it has been established to be of order 0 and type 1, 1. This kind of pseudodifferential operators seems to appear rather frequently in the discussions around nonlinear hyperbolic equations, see e.g. [Hö96].

Two natural directions deserve to be further investigated. One regards the extension of some of our achievements to curved spacetimes. Switching on the curvature, at least dealing with geodesically convex neighborhoods, the geometric picture is not very far from that in Minkowski spacetime. At the level of Weyl algebras – and it can be extended to include the algebras of Wick polynomials – the main features of the bulk-boundary interplay for geodesic double cones survive the appearance of the curvature, as recently

established in [DPP10]. In that case, states strictly analogous to λ can be defined on the boundary algebra and the induced states in the bulk – corresponding to our state ω_m – turn out to be of Hadamard class. It would be interesting to focus on the features of modular groups of the respective von Neumann algebras.

The other possible generalization could obviously consist of relaxing the requirement of having a free quantum field. In that case, however, the symplectic structure, which plays the crucial role in building up the bulk-boundary correspondence would not exist, stopping the construction at the first step. It seems plausible that a suitable extension of the time-slice axiom to null Cauchy surfaces may help to cover the gap.

As recalled in the introduction, there are many other ideas related to the results presented in this paper. A particularly nice thing to do would be to try to push even further the intuition that the modular group is a form of dynamics, by going to precisely define, in the spirit of the paper, the idea of a “local dynamics” generated by these modular groups and study its properties. We have some interesting preliminary result, at the moment only at the level of symplectic spaces and we hope to report soon on them. Another appealing idea would be to try to study the properties of the group generated by all the (local) modular operators.

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A Proof of some propositions

Proof of Lemma 3.1.5. Here Σ_D is the restriction to D of the Cauchy surface for \mathbb{M} at $t = 1$ so that C is the boundary $\partial\Sigma_D$ of Σ_D in $\Sigma \equiv \mathbb{R}^3$. If $\phi \in \mathcal{S}_m(D)$, its Cauchy data on Σ are in $C_0^\infty(\Sigma_D)$ so that they also belong to the Schwartz space on \mathbb{R}^3 , hence we can define $\widehat{\phi}$ as in (24). Since $m > 0$, $\widehat{\phi}$ belongs again to the Schwartz space. Thus (22) and (23) hold true interpreting the integrals in the standard way, since the smooth function defined in the right-hand side of (22) must coincide to ϕ (more precisely, $L_{\mathbb{M}D}^m \phi$) everywhere on \mathbb{M} , in view of the standard uniqueness theorem for the smooth Cauchy problem for Klein-Gordon equation (e.g. Theorem 10.1.2 in [Wa84]); indeed both functions satisfy the Klein-Gordon equation in \mathbb{M} and match the same Cauchy data at $t = 0$ on the whole Σ . If $\phi \in \widetilde{\mathcal{S}}_m(D)$ the situation is more complicated. However, in view of the definition of $\widetilde{\mathcal{S}}_m(D)$, (a) $\phi(1, \cdot)$ and $\partial_t \phi(1, \cdot)$ are smooth in $\Sigma_D \cup \partial\Sigma_D$ (which has finite Lebesgue measure) vanishing outside it and (b) $\phi(1, \mathbf{x}) = 0$ if $\mathbf{x} \in \partial\Sigma_D$ and outside $\Sigma_D \cup \partial\Sigma_D$ on Σ . From (a) $\partial_t \phi(1, \cdot) \in L^2(\mathbb{R}^3, d\mathbf{x})$ is bounded and has compact support and thus its Fourier transform $\widehat{\partial_t \phi(1, \cdot)}(\mathbf{k})$ is bounded and $\widehat{\partial_t \phi(1, \cdot)}(\mathbf{k}), E(\mathbf{k})^{-1} \widehat{\partial_t \phi(1, \cdot)}(\mathbf{k}), E(\mathbf{k})^{-1/2} \widehat{\partial_t \phi(1, \cdot)}(\mathbf{k})$ belong to $L^2(\mathbb{R}^3, d\mathbf{k})$ (even for $m = 0$). From (a) and (b) $\phi(1, \cdot) \in H_0^1(\Sigma_D) \subset H^1(\mathbb{R}^3)$ by the Sobolev trace theorem (e.g. see Theorem 9.5.1 in [Au00]). Consequently $\widehat{\phi(1, \cdot)}(\mathbf{k}), E(\mathbf{k}) \widehat{\phi(1, \cdot)}(\mathbf{k})$ belongs

to $L^2(\mathbb{R}^3, d\mathbf{k})$. We conclude from (24) that, on a hand $\widehat{\phi}$ belongs to $L^2(\mathbb{R}^3, d\mathbf{k})$ when $\phi \in \widetilde{\mathcal{S}}_m(D)$, on the other hand the right-hand sides of (22) and (23), *henceforth indicated by ψ and $\partial_t\psi$ respectively*, are well defined *if interpreting the integrals in the sense of Fourier-Plancherel transform*. By direct inspection one sees that the right-hand side of (22), ψ , is a weak solution of Klein-Gordon equation in \mathbb{M} . Furthermore $\psi \in C([1, T]; H^1(\mathbb{R}^3)) \cap C^1([1, T]; H^0(\mathbb{R}^3))$, for every $T > 1$, and it admits as initial conditions for $t = 1$ just the functions $\phi(1, \cdot) \in H^1(\mathbb{R}^3)$ and $\partial_t\phi(1, \cdot) \in H^0(\mathbb{R}^3)$ (extended to the zero functions outside Σ_D). We want to prove that $\psi = \phi$ in D . First consider the set $[1, +\infty) \times \mathbb{R}^3$, that is the causal future of Σ_D . We can smoothly extend $\phi \in \widetilde{\mathcal{S}}_m(D)$ in the whole set $J^+(o)$ as the unique solution of the Goursat problem (Theorem 5.4.2 in [Fr75]) with characteristic datum $u^{-1}L_{V\widetilde{D}}^m\phi$ (smoothly) extended to the zero function outside $\partial J^+(o) \setminus V$. Then, we can further extend ϕ to the the half space $[1, +\infty) \times \mathbb{R}^3$, simply defining the extension as the null function outside $J^+(o)$. This extension, by construction, is a smooth solution of the Klein-Gordon equation in $(1, +\infty) \times \mathbb{R}^3$ away from $\partial J^+(o)$. However it is a weak solution of the Klein-Gordon equation in a neighborhood of $\partial J^+(o)$ in $(1, +\infty) \times \mathbb{R}^3$. For a test function $\psi \in C_0^\infty(\mathbb{M})$, it can be verified by direct inspection using the fact that the boundary term on $\partial J^+(o)$ arising by employing the integration by parts and Stokes-Poincaré's lemma is proportional to $\int_{\partial J^+(o) \cap ((1, +\infty) \times \mathbb{R}^3)} d\omega du (\psi \partial_u \phi - \phi \partial_u \psi) = 0$, since the (extended) function ϕ vanishes on $\partial J^+(o) \cap ((1, +\infty) \times \mathbb{R}^3) = \partial J^+(o) \setminus V$ with its derivative ∂_u (i.e. along the affine parameter u describing the null geodesics forming $\partial J^+(o)$). Summing up, the extended function ϕ on the whole $[1, +\infty) \times \mathbb{R}^3$ is a weak solution of Klein-Gordon equation in $(1, +\infty) \times \mathbb{R}^3$ that matches the initial conditions $\phi(1, \cdot) \in H^1(\mathbb{R}^3)$ and $\partial_t\phi(1, \cdot) \in H^0(\mathbb{R}^3)$ (extended to the zero functions outside Σ_D). By direct inspection one sees that the constructed extended ϕ is in the class $C([1, T]; H^1(\mathbb{R}^3)) \cap C^1([1, T]; H^0(\mathbb{R}^3))$ for every $T > 1$ (in particular using the fact that $\phi(t, \cdot) \in H_0^1(J^+(o) \cap \Sigma_t)$ for every $t \geq 1$, Σ_t being the Cauchy surface of \mathbb{M} at fixed t). In view of Theorem 3.2 in [So08], in $[1, +\infty) \times \mathbb{R}^3$, the extended ϕ must coincide to the other weak solution ψ which satisfies the same requirements and initial data. In particular, $\phi = \psi$ in $D \cap ([1, 2] \times \mathbb{R}^3)$ so that (22) holds true therein. To conclude we prove (22), i.e. that $\phi = \psi$, even in $D \cap ([0, 1] \times \mathbb{R}^3)$. If ψ were smooth in the open set $D \cap ([0, 1] \times \mathbb{R}^3)$, in view of the standard uniqueness theorem for the Cauchy problem in $D \cap ([0, 1] \times \mathbb{R}^3)$ (e.g. Theorem 10.1.2 in [Wa84]), we would have $\phi = \psi$ in $D \cap ([0, 1] \times \mathbb{R}^3)$. To prove that ψ is smooth therein we notice that if it were not the case for a point $p \in D \cap ([0, 1] \times \mathbb{R}^3)$, there would be a full complete null geodesic γ through p with γ included in the singular support of the distribution ψ , in view of the *Hörmander's theorem of propagation of singularities*, since ψ solves the Klein-Gordon equation. On the other hand γ would enter the region $D \cap ([1, 2] \times \mathbb{R}^3)$, where we already know that ψ is smooth and thus its singular support is empty, giving rise to a contradiction. Thus ψ is smooth in $D \cap ([0, 1] \times \mathbb{R}^3)$ as wanted and (22) holds in $D \cap ([0, 1] \times \mathbb{R}^3)$, too. \square

Proof of Lemma 4.2.2. In the hypotheses of the lemma, assuming in particular that the case (b) holds (the other is similar) we have that both $\overline{\Phi'(\omega, k)}$ and $k\overline{\Phi'(\omega, k)}$ define

functions in $L^2(\mathbb{S}^2 \times \mathbb{R}, d\omega dk)$ varying $(\omega, k) \in \mathbb{S}^2 \times \mathbb{R}$. Now defining $\theta(k) = 1$ for $k \geq 1$ and $\theta(k) = 0$ otherwise, we have:

$$\int_{\mathbb{S}^2 \times \mathbb{R}^+} d\omega dk \overline{2k\Phi(\omega, k)} \widehat{\Phi'(\omega, k)} = \lim_{\epsilon \rightarrow 0^+} \int_{\mathbb{S}^2 \times \mathbb{R}} d\omega dk \overline{2\Phi(\omega, k)} \left(e^{-\epsilon k} \theta(k) k \widehat{\Phi'(\omega, k)} \right).$$

Using the fact that $\mathbb{R} \ni k \mapsto e^{-\epsilon k} \theta(k)$ is $L^2(\mathbb{R}, du)$ and is the Fourier-Plancherel transform of $\mathbb{R} \ni u' \mapsto (2\pi i)^{-1}/(u' - i\epsilon)$ and making use of the convolution theorem for $L^2(\mathbb{R}, du)$ functions [FJ75] as well as Plancherel's theorem, the integral in the right-hand side can be re-written as:

$$\int_{\mathbb{S}^2 \times \mathbb{R}} d\omega dk \overline{2\Phi(\omega, k)} \left(e^{-\epsilon k} \theta(k) k \widehat{\Phi'(\omega, k)} \right) = \int_{\mathbb{S}^2 \times \mathbb{R}} d\omega du \overline{\Phi(\omega, u)} \int_{\mathbb{R}} du' \frac{\partial_u \Phi'(\omega, u')}{u - u' - i\epsilon}.$$

Integrating by parts in the last integral, taking into account that $\Phi'(\omega, u) \rightarrow 0$ as $u \rightarrow 0_+$, and explicitly writing the supports of the functions, we end up with:

$$\int_{\mathbb{S}^2 \times \mathbb{R}} d\omega dk \overline{2\Phi(\omega, k)} e^{-\epsilon k} \theta(k) k \widehat{\Phi'(\omega, k)} = -\frac{1}{\pi} \int_{\mathbb{S}^2 \times [0, 1]} d\omega du \int_{[0, 1]} du' \frac{\overline{\Phi(\omega, u)} \Phi'(\omega, u')}{(u - u' - i\epsilon)^2}. \quad (87)$$

Eventually, we notice that $\int_{\mathbb{S}^2 \times [0, 1]} d\omega du \int_{[0, 1]} du' \left| \frac{\overline{\Phi(\omega, u)} \Phi'(\omega, u')}{(u - u' - i\epsilon)^2} \right| < +\infty$ because the function $\mathbb{S}^2 \times [0, 1] \times [0, 1] \ni (\omega, u, u') \mapsto |\Phi'(\omega, u')|/(|u - u'|^2 + \epsilon^2)$ is bounded by construction ($\Phi'(\omega, u)$ vanishes uniformly in the angles as $u \rightarrow 0$) and $\mathbb{S}^2 \times [0, 1] \ni (\omega, u) \mapsto \Phi(\omega, u)$ is $L^1(\mathbb{S}^2 \times [0, 1], d\omega du)$ it being $L^2(\mathbb{S}^2 \times [0, 1], d\omega du)$. This implies that we can apply Fubini-Tonelli's theorem in the last iterated integral in (87) computing it in the measure product, proving (43) and concluding the proof. \square

Proof of Lemma 4.2.4. In this proof, for every $\phi \in \mathcal{S}_m(D)$, the same symbol is used for its unique smooth extension $L_{\mathbb{M}D}^m \phi \in \mathcal{S}_m(\mathbb{M})$ defined in the whole \mathbb{M} . The right-hand side in formula (24) is nothing but the integral over Σ of the 3-form

$$\eta_{\mathbf{k}, \phi} \doteq -i \frac{1}{3!} \sqrt{|g|} g^{ea} \left(\frac{e^{-ik_l x^l}}{(2\pi)^{3/2} \sqrt{2E}} \partial_e \phi - \phi \partial_e \frac{e^{-ik_l x^l}}{(2\pi)^{3/2} \sqrt{2E}} \right) \epsilon_{abcd} dx^b \wedge dx^c \wedge dx^d \quad (88)$$

where ϵ is completely antisymmetric and $\epsilon_{1234} = 1$. As $\eta_{\mathbf{k}, \phi}$ is constructed out of solutions of Klein-Gordon equation, it satisfies the conservation relation $d\eta_{\mathbf{k}, \phi} = 0$ (which corresponds to the conservation of the associated current $J_{\mathbf{k}, \phi} \doteq * \eta_{\mathbf{k}, \phi}$). Now we specialize Σ to the surface at $t = 1$, and consider $\phi \in \mathcal{S}_m(D)$ so that its Cauchy data are supported in the portion of Σ included in D . Stokes-Poincaré's theorem applied to $\eta_{\mathbf{k}, \phi}$ leads immediately to $\widehat{\phi}(\mathbf{k}) = \int_{\Sigma} \eta_{\mathbf{k}, \phi} = \int_V \eta_{\mathbf{k}, \phi}$, where both 3-surfaces have future-oriented normal vector. Exploiting the fact that, on V , $t = u = \|\mathbf{x}\|$ so that one can write $\mathbf{x} = u\omega$ for some vector $\omega \in \mathbb{S}^2$, the second integral gives (integrating by parts noticing that $L_{V D}^m(\phi)$ vanishes for $u = 0$ and $u \geq 1$)

$$\widehat{\phi}(\mathbf{k}) = \frac{-2i}{(2\pi)^{3/2} \sqrt{2E}} \int_{\mathbb{S}^2 \times [0, 1]} d\omega du u e^{-i(\mathbf{k} \cdot u\omega - Eu)} \partial_u L_{V D}^m(\phi)(\omega, u). \quad (89)$$

Now we can insert the right-hand side of (89), computed for ϕ_1 and ϕ_2 , in the left-hand side of (50) and try to interchange the order of integrations. This cannot be done directly, but an ϵ prescription needs. However, a straightforward application of Fubini-Tonelli's theorem and Lebesgue's dominated convergence theorem lead to, if $T \doteq \mathbb{S}^2 \times \mathbb{S}^2 \times [0, 1] \times [0, 1]$

$$\int_{\mathbb{R}^3} d\mathbf{k} \widehat{\phi_1}(\mathbf{k}) \widehat{\phi_2}(\mathbf{k}) = \lim_{\epsilon \rightarrow 0^+} \frac{2}{(2\pi)^3} \int_T d\omega d\omega' du du' (\partial_u \Phi_1(\omega, u)) (\partial_{u'} \Phi_2(\omega', u')) \times \int_{\mathbb{R}^3} d\mathbf{k} \frac{uu'}{E} e^{-i(\mathbf{k} \cdot (u'\omega' - u\omega) - E(u' - u + i\epsilon))}, \quad (90)$$

where $\Phi_1 \doteq L_{VD}^m(\phi_1)$ and $\Phi_2 \doteq L_{VD}^m(\phi_2)$. We start considering the case $m > 0$. The last integral in (90) can be explicitly computed passing in polar coordinates in the variable \mathbf{k} and exploiting the result 3.961 in [GR95]:

$$\int_{\mathbb{R}^3} d\mathbf{k} E^{-1} e^{-i(\mathbf{k} \cdot (u'\omega' - u\omega) - E(u' - u + i\epsilon))} = I(u, u', \omega, \omega', m, \epsilon)$$

where

$$I(u, u', \omega, \omega', m, \epsilon) \doteq \frac{4\pi m K_1 \left(m \sqrt{2uu'(1 - \cos \widehat{\omega\omega'})} + 2i\epsilon(u - u' - i\epsilon/2) \right)}{\sqrt{2uu'(1 - \cos \widehat{\omega\omega'})} + 2i\epsilon(u - u' - i\epsilon/2)}.$$

(From now on, considering the complex functions we shall encounter which admit the origin as branch point, we always assume that the cut stays along the negative real axis.) For $\epsilon > 0$ fixed, this function is smooth and exponentially vanishes as the argument of K_1 tends to $+\infty$. We have obtained that:

$$\int_{\mathbb{R}^3} d\mathbf{k} \widehat{\phi_1}(\mathbf{k}) \widehat{\phi_2}(\mathbf{k}) = J$$

where

$$J \doteq \lim_{\epsilon \rightarrow 0^+} \int_T d\omega d\omega' du du' uu' (\partial_u \Phi_1(\omega, u)) (\partial_{u'} \Phi_2(\omega', u')) I(u, u', \omega, \omega', m, \epsilon).$$

Our last step consists of working out the limit above explicitly showing that it can be re-arranged as the right-hand side of (41). First of all, we notice that the various integrations computed before taking the limit can be interchanged in view of Fubini-Tonelli theorem, as the integrand is continuous with compact support and thus they are integrable in the product measure. So we start performing the integral over the angles $(\omega, \omega') \in \mathbb{S}^2 \times \mathbb{S}^2$ decomposing it into two integrals computed over two corresponding regions. A_δ individuated by $1 - \cos \widehat{\omega\omega'} > \delta$ and B_δ individuated by $0 \leq 1 - \cos \widehat{\omega\omega'} \leq \delta$ for a fixed $\delta > 0$. Let us separately consider the two terms of the decomposition:

$$\lim_{\epsilon \rightarrow 0^+} \int_T = \lim_{\epsilon \rightarrow 0^+} \int_{A_\delta \times [0,1]^2} + \lim_{\epsilon \rightarrow 0^+} \int_{B_\delta \times [0,1]^2}$$

of the integral in the right-hand side of (A). The kernel containing the parameter ϵ is jointly continuous in all variables including $\epsilon \in [0, \epsilon_0]$ and thus it is ϵ -uniformly bounded in $A_\delta \times [0, 1]^2$. Using Lebesgue's theorem again:

$$\begin{aligned} \lim_{\epsilon \rightarrow 0^+} \int_{A_\delta \times [0, 1]^2} d\omega d\omega' dudu'uu' (\partial_u \Phi_1(\omega, u) \partial_{u'} \Phi_2(\omega', u')) I(u, u', \omega, \omega', m, \epsilon) \\ = \int_{A_\delta \times \mathbb{R}_+^2} d\omega d\omega' dudu'uu' (\partial_u \Phi_1(\omega, u) \partial_{u'} \Phi_2(\omega', u')) I(u, u', \omega, \omega', m, 0) \\ = \int_{A_\delta \times \mathbb{R}_+^2} d\omega d\omega' dv dv' vv' f_\lambda(v, \omega) g_\lambda(v', \omega') I(v, v', \omega, \omega', m, 0) \end{aligned}$$

where we have replaced the range $[0, 1]$ of u, u' with \mathbb{R}_+ since ψ and ψ' vanish for $u, u' \geq 1$ and, only in the last step, we have changed variables: $u \rightarrow v \doteq \lambda u$ and $u' \rightarrow v' \doteq u'/\lambda$, where $\lambda > 0$ is a fixed real and we have defined $f_\lambda(x, \omega) \doteq \partial_u \Phi_1(\omega, u)|_{u=x/\lambda}$ and $g_\lambda(x, \omega) \doteq \partial_{u'} \Phi_2(\omega, u)|_{u=\lambda x}$. Notice that the result cannot depend on λ , so we are allowed to take the limit as $\lambda \rightarrow +\infty$. Notice that $\mathbb{S}^2 \times \mathbb{S}^2 \times \mathbb{R}_+ \times \mathbb{R}_+ \ni (\omega, \omega', v, v') \mapsto f_\lambda(v, \omega) g_\lambda(v', \omega')$ is λ -uniformly bounded by construction because $\Phi = u\phi$ has bounded u -derivative as it can be proved by direct inspection starting from the fact that the functions $\phi \in \mathcal{S}_m(\mathbb{M})$ are smooth. Furthermore, point-wisely, $f_\lambda(v, \omega) g_\lambda(v', \omega') \rightarrow \partial_u \Phi_1(0, \omega) 0 = 0$ constantly as $\lambda \rightarrow +\infty$ because $\Phi_2(\omega, u)$ smoothly vanishes if $u > 1$. Finally the function

$$A_\delta \times \mathbb{R}_+ \times \mathbb{R}_+ \ni (\omega, \omega', v, v') \mapsto \left| vv' \frac{mK_1 \left(m\sqrt{2vv'(1 - \cos \widehat{\omega\omega'})} \right)}{\pi^2 \sqrt{2vv'(1 - \cos \widehat{\omega\omega'})}} \right|$$

is integrable. Summing up, Lebesgue's theorem of dominated convergence eventually proves that, for every $\delta > 0$:

$$\begin{aligned} \lim_{\epsilon \rightarrow 0^+} \int_{A_\delta \times [0, 1]^2} d\omega d\omega' dudu'uu' \partial_u \Phi_1(\omega, u) \partial_{u'} \Phi_2(\omega', u') I(u, u', \omega, \omega', m, \epsilon) \\ = \lim_{\lambda \rightarrow +\infty} \int_{A_\delta \times \mathbb{R}_+^2} d\omega d\omega' dv dv' v^2 v'^2 f_\lambda(v, \omega) g_\lambda(v', \omega') I(v, v', \omega, \omega', m, 0) = 0. \end{aligned}$$

We conclude that, for every $\delta > 0$, $\int d\mathbf{k} \widehat{\phi_1} \widehat{\phi_2}$ amounts to:

$$\lim_{\epsilon \rightarrow 0^+} \int_{B_\delta \times [0, 1]^2} d\omega d\omega' dudu'uu' (\partial_u \Phi_1(\omega, u) \partial_{u'} \Phi_2(\omega', u')) I(u, u', \omega, \omega', m, \epsilon).$$

To compute that limit, we notice that, while evaluating the integral above, the argument of K_1 varies, in fact, in a bounded set, since ϵ can be fixed in a interval $[0, \epsilon_0)$ and the supports of the $\partial_u \phi_i$ are compact, therefore we can exploit the expansion about $z = 0$:

$$\frac{mK_1(mz)}{z} = \frac{1}{z^2} - \frac{m^2}{4} \ln(z^2) + O(1). \quad (91)$$

The last two terms inserted in the integrand over $B_\delta \times [0, 1]^2$ give rise to functions which are integrable also for $\epsilon = 0$. In that case, using Lebesgue's theorem again, one sees that the limit as $\epsilon \rightarrow 0_+$ can be computed by direct substitution of ϵ with 0 in the integrand. Next, as $\delta > 0$ can be taken arbitrarily small and using the fact that the measure of $B_\delta \times [0, 1]^2$ vanishes as $\delta \rightarrow 0$, one easily concludes that the terms $-\frac{m^2}{4} \ln(z^2) + O(1)$ give no contribution to $\int d\mathbf{k} \widehat{\phi}_1 \widehat{\phi}_2$. We end up with, where we make explicit the factor u in the definition of $\Phi(\omega, u) = u\phi(\omega, u)$:

$$\int_{\mathbb{R}^3} d\mathbf{k} \widehat{\phi}_1(\mathbf{k}) \widehat{\phi}_2(\mathbf{k}) = \lim_{\epsilon \rightarrow 0_+} \int_{B_\delta \times [0, 1]^2} d\omega d\omega' dud u' \frac{uu' \partial_u u \phi_1(\omega, u) \partial_{u'} u' \phi_2(\omega', u')}{2\pi^2 \left(uu'(1 - \cos \widehat{\omega\omega'}) + i\epsilon(u - u' - i\epsilon/2) \right)}. \quad (92)$$

Integrating by parts in the variable u we obtain:

$$\begin{aligned} & \int_{B_\delta \times [0, 1]^2} d\omega d\omega' dud u' \frac{uu' \partial_u u \phi_1(\omega, u) \partial_{u'} u' \phi_2(\omega', u')}{2\pi^2 \left(uu'(1 - \cos \widehat{\omega\omega'}) + i\epsilon(u - u' - i\epsilon/2) \right)} = \\ & - \int_{B_\delta \times [0, 1]^2} d\omega d\omega' dud u' uu' \phi_1(\omega, u) (\partial_{u'} u' \phi_2(\omega', u')) \partial_u \frac{u}{2\pi^2 \left(uu'(1 - \cos \widehat{\omega\omega'}) + i\epsilon(u - u' - i\epsilon/2) \right)} \\ & = \int_{B_\delta \times [0, 1]^2} d\omega d\omega' dud u' \frac{\phi_1(\omega, u) (\partial_{u'} u' \phi_2(\omega', u')) (i\epsilon uu'^2 + \epsilon^2 uu'/2)}{2\pi^2 \left(uu'(1 - \cos \widehat{\omega\omega'}) + i\epsilon(u - u' - i\epsilon/2) \right)^2}. \end{aligned} \quad (93)$$

The part of integral proportional to ϵ^2 in the last line vanishes as $\epsilon \rightarrow 0_+$ as we go to prove. Below ω is individuated by its polar angles, $\varphi \in (-\pi, \pi)$, $\theta \in (0, \pi)$. Furthermore, we have fixed the polar axis z to be ω' performing the integration in $d\omega$ in the last three integrals, so that, in fact, θ and φ depends on ω' parametrically, too. We have, for some constant $C \geq 0$ depending on ϕ_1 and ϕ_2

$$\begin{aligned} & \left| \int_{B_\delta \times [0, 1]^2} d\omega d\omega' dud u' \phi_1(\omega, u) (\partial_{u'} u' \phi_2(\omega', u')) \frac{\epsilon^2 uu'}{2\pi^2 \left(uu'(1 - \cos \widehat{\omega\omega'}) + i\epsilon(u - u' - i\epsilon/2) \right)^2} \right| \\ & \leq C \int_{B_\delta \times [0, 1]^2} d\omega d\omega' dud u' \frac{\epsilon^2 uu'}{\left(uu'(1 - \cos \widehat{\omega\omega'}) + \epsilon^2 \right)^2 + \epsilon^2 (u - u')^2} \\ & = C \int_{\mathbb{S}^2 \times [0, 1]^2} d\omega' d\varphi \int_{1-\delta}^1 d \cos \theta \frac{\epsilon^2 uu'}{\left(uu'(1 - \cos \widehat{\omega\omega'}) + \epsilon^2 \right)^2 + \epsilon^2 (u - u')^2} \\ & = 2\pi C \int_{[0, 1]^2} dud u' \int_0^\delta d\mu \frac{\epsilon^2 uu'}{(uu'\mu + \epsilon^2)^2 + \epsilon^2 (u - u')^2} \\ & = 2\pi C \int_{[0, 1]^2} dud u' \frac{\epsilon}{|u - u'|} \arctan \left(\frac{|u - u'|}{\epsilon} \frac{2uu'\delta}{2|u - u'|^2 + 2uu'\delta + \epsilon^2} \right) \\ & \leq 2\pi C \int_{[0, 1]^2} dud u' \frac{\epsilon}{|u - u'|} \arctan \left(\frac{|u - u'|}{\epsilon} \right) \rightarrow 0. \end{aligned} \quad (94)$$

In the next to the last step we exploited the fact that $0 \leq \frac{2uu'\delta}{2|u-u'|^2+2uu'\delta+\epsilon^2} \leq 1$ and that the function $\mathbb{R}_+ \ni \arctan x$ is increasing. In the last step we used the fact that $\mathbb{R}_+ \ni x^{-1} \arctan x$ is bounded (so, the integrand above is bounded when $(u, u') \in [0, 1]^2$), and $\frac{\epsilon}{|u-u'|} \arctan\left(\frac{|u-u'|}{\epsilon}\right) \rightarrow 0$ pointwisely as $\epsilon \rightarrow 0_+$ so that we can exploit Lebesgue's dominated convergence theorem.

Let us come back to (93), where we drop the part of the integral proportional to ϵ^2 and we notice that;

$$\frac{uu'}{\left(uu'(1 - \cos \widehat{\omega\omega'}) + i\epsilon(u - u' - i\epsilon/2)\right)^2} = \partial_{\cos \widehat{\omega\omega'}} \left(\frac{1}{uu'(1 - \cos \widehat{\omega\omega'}) + i\epsilon(u - u' - i\epsilon/2)} \right)$$

Starting from (92), inserting this identity in (93), performing an integration by parts we get:

$$\begin{aligned} & \int_{\mathbb{R}^3} d\mathbf{k} \overline{\widehat{\phi}_1(\mathbf{k})} \widehat{\phi}_2(\mathbf{k}) = \\ & \frac{1}{2\pi^2} \lim_{\epsilon \rightarrow 0_+} \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' du' (\partial_{u'} u' \phi_2(\omega', u')) \int_0^{2\pi} d\varphi \int_{1-\delta}^1 d \cos \theta \partial_{\cos \theta} \left(\frac{i\epsilon u' \phi_1(u, \theta, \varphi)}{uu'(1 - \cos \theta) + i\epsilon(u - u' - i\epsilon/2)} \right) \\ & - \frac{1}{2\pi^2} \lim_{\epsilon \rightarrow 0_+} \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' du' (\partial_{u'} u' \phi_2(\omega', u')) \int_0^{2\pi} d\varphi \int_{1-\delta}^1 d \cos \theta \left(\frac{i\epsilon u' \partial_{\cos \theta} \phi_1(u, \theta, \varphi)}{uu'(1 - \cos \theta) + i\epsilon(u - u' - i\epsilon/2)} \right) \\ & = \frac{1}{2\pi^2} \lim_{\epsilon \rightarrow 0_+} \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' du' (\partial_{u'} u' \phi_2(\omega', u')) 2\pi \left(\frac{i\epsilon u' \phi_1(u, \omega')}{i\epsilon(u - u' - i\epsilon/2)} \right) \\ & - \frac{1}{2\pi^2} \lim_{\epsilon \rightarrow 0_+} \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' du' (\partial_{u'} u' \phi_2(\omega', u')) \int_0^{2\pi} d\varphi \left(\frac{i\epsilon u' \phi_1(u, 1 - \delta, \varphi)}{uu'\delta + i\epsilon(u - u' - i\epsilon/2)} \right) \\ & - \frac{1}{2\pi^2} \lim_{\epsilon \rightarrow 0_+} \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' du' (\partial_{u'} u' \phi_2(\omega', u')) \int_0^{2\pi} d\varphi \int_{1-\delta}^1 d \cos \theta \left(\frac{i\epsilon u' \partial_{\cos \theta} \phi_1(u, \theta, \varphi)}{uu'(1 - \cos \theta) + i\epsilon(u - u' - i\epsilon/2)} \right). \end{aligned} \quad (95)$$

As we prove shortly, the last two integrals vanish so that we get, integrating by parts in the remaining integral:

$$\int_{\mathbb{R}^3} d\mathbf{k} \overline{\widehat{\phi}_1(\mathbf{k})} \widehat{\phi}_2(\mathbf{k}) = \lim_{\epsilon \rightarrow 0_+} -\frac{1}{\pi} \int_{\mathbb{S}^2 \times [0,1] \times [0,1]} d\omega du' \frac{u \phi_1(\omega, u) u' \phi_2'(\omega, u')}{(u - u' - i\epsilon)^2}. \quad (96)$$

This is just what we need to conclude the proof of the theorem. To complete the proof we have to prove that the last two integrals in (95) vanish. Defining, if $\cos \theta_\delta = 1 - \delta$,

$$\Psi(u) \doteq \int_0^{2\pi} d\varphi \phi_1(u, \theta_\delta, \varphi),$$

the next to last integral in (95) can be re-arranged as:

$$i\epsilon \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' du' (\partial_{u'} u' \phi_2(\omega', u')) \Psi(u) \frac{u'}{u'\delta + i\epsilon} \partial_u \ln(uu'\delta + i\epsilon(u - u' - i\epsilon/2)).$$

In turn this integral decomposes into three parts if integrating by parts:

$$\begin{aligned}
& -i\epsilon \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' dud u' (\partial_{u'} u' \phi_2(\omega', u')) \partial_u \Psi(u) \frac{u'}{u'\delta + i\epsilon} \ln(uu'\delta + i\epsilon(u - u' - i\epsilon/2)) \\
& + i\epsilon \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' dud u' (\partial_{u'} u' \phi_2(\omega', u')) \Psi(0) \frac{u'}{u'\delta + i\epsilon} \ln(u' + i\epsilon/2) \\
& + i\epsilon \ln(-i\epsilon) \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' dud u' (\partial_{u'} u' \phi_2(\omega', u')) \Psi(0) \frac{u'}{u'\delta + i\epsilon}. \tag{97}
\end{aligned}$$

Then notice that the logarithm as well as the fraction $u'/(u'\delta - i\epsilon)$ are integrable for every values of $\epsilon \geq 0$, including $\epsilon = 0$. Using the fact that $|u'/(u'\delta - i\epsilon)|$ is bounded uniformly in ϵ and that $u'\phi_2(\omega', u')\partial_u \Psi(u)$ is bounded on the domain of integration, each integral above can be proved to vanish as $\epsilon \rightarrow 0_+$ due to the presence of the factor $i\epsilon$. Let us discuss in details the most difficult integral, that is the first one. Assuming $\delta, \epsilon > 0$ sufficiently small, taking the absolute value of the integrand, the first integral in (97) can be bounded by, where $K \geq 0$ is some constant depending on ϕ_1 and ϕ_2 ,

$$\begin{aligned}
& \epsilon K \int_{[0,1]^2} dud u' \sqrt{[\ln((uu'\delta + \epsilon^2/2)^2 + \epsilon^2(u - u')^2)]^2 + \arctan^2\left(\frac{\epsilon(u - u')}{uu'\delta + \epsilon^2/2}\right)} \\
& \leq \epsilon K \int_{[0,1]^2} dud u' \sqrt{[\ln(uu'\delta)]^2 + \pi^2/4} \rightarrow 0 \quad \text{as } \epsilon \rightarrow 0_+. \tag{98}
\end{aligned}$$

The remaining two integrals in (97) can be treated similarly proving much more straightforwardly that they vanish as well. Finally, turning our attention on (95) again, defining:

$$\frac{u}{\sqrt{1 - \cos \theta}} \Psi(u, \theta) \doteq \int_0^{2\pi} d\varphi \partial_{\cos \theta} \phi_1(u, \theta, \varphi),$$

where Ψ is bounded. The factor $u/\sqrt{1 - \cos \theta}$ arises from the fact that, in view of the restriction to V , $\phi|_V = \phi(u, u \cos \theta, u \sin \theta \cos \varphi, u \sin \theta \sin \varphi)$ and $d \sin \theta / d \cos \theta = \cos \theta / \sqrt{(1 - \cos \theta)(1 + \cos \theta)}$. Thus the last integral in (95) can be re-arranged as:

$$i\epsilon \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' dud u' (\partial_{u'} u' \phi_2(\omega', u')) \int_0^\delta \frac{d\mu}{\sqrt{\mu}} u \Psi(u, \theta) \frac{u'}{u'\mu + i\epsilon} \partial_u \ln(uu'\mu + i\epsilon(u - u' - i\epsilon/2)).$$

By integrating by parts it can be re-arranged to:

$$-i2\epsilon \int_{\mathbb{S}^2 \times [0,1]^2} d\omega' dud u' (\partial_{u'} u' \phi_2(\omega', u')) \int_0^{\sqrt{\delta}} d\sqrt{\mu} \frac{\partial_u u \Psi(u, \mu) u'}{u'(\sqrt{\mu})^2 + i\epsilon} \ln(u'\mu + i\epsilon(u - u' - i\epsilon/2)).$$

As before, assuming ϵ and δ small enough, this integral can be bounded by:

$$\begin{aligned}
& \epsilon K' \int_{[0,1]^2 \times [0, \sqrt{\delta}]} dud u' dx \sqrt{[\ln((uu'x + \epsilon^2/2)^2 + \epsilon^2(u - u')^2)]^2 + \arctan^2\left(\frac{\epsilon(u - u')}{uu'x + \epsilon^2/2}\right)} \\
& \leq \epsilon K' \int_{[0,1]^2 \times [0, \sqrt{\delta}]} dud u' dx \sqrt{[\ln(uu'x)]^2 + \pi^2/4} \rightarrow 0 \quad \text{as } \epsilon \rightarrow 0_+. \tag{99}
\end{aligned}$$

where $K' \geq 0$ is some constant depending on ϕ_1 and ϕ_2 .

The case $m = 0$ is very similar but much more simple. In fact, coming back to (90), making use of 3.895 in [GR95], (92) arises directly, the analog integral on A_δ being vanishing similarly. Afterwards the proof is the same as for $m > 0$. \square

Proof of Theorem 5.1.1. Replace the coordinate $u \in (0, 1)$ on V with the coordinate $\ell \in (-\infty, +\infty)$ such that $u = 1/(1 + e^{-\ell})$ and $X = \partial_\ell$ induces trivial displacements $\ell \rightarrow \ell + \tau$. The functions $\Phi \in \mathbf{S}(V)$ turns out to be smooth jointly in the variables (ω, ℓ) , they vanishes for ℓ sufficiently large and vanish with all of their derivatives, uniformly in ω , for $\ell \rightarrow -\infty$ with order $O(e^\ell)$. In particular $\Phi \in \mathbf{S}(V)$ belongs to the Schwartz space $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ of the complex smooth functions which vanishes as $|\ell| \rightarrow +\infty$, uniformly in ω and with all of their derivatives, faster than every inverse power of $|\ell|$. Therefore the ℓ -Fourier transform of $\Phi \in \mathbf{S}(V)$ is well defined and belongs to $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ again. Then the function $K'_\lambda : \mathbf{S}(V) \rightarrow L^2(\mathbb{S}^2 \times \mathbb{R}; m(h)d\omega dh)$ is well defined (notice that $|m(h)| \leq C(1 + |h|^2)$ for some constant $C \geq 0$ and all $\ell \in \mathbb{R}$). By direct inspection, using standard theorems of Fourier transform theory, and the fact that $\overline{\Phi(\omega, -h)} = \overline{\Phi(\omega, h)}$ if $\Phi \in \mathbf{S}(V)$ since it is real, one has that

$$\sigma(\Phi, \Phi') = -2Im\langle K'_\lambda \Phi, K'_\lambda \Phi' \rangle_{L^2(\mathbb{S}^2 \times \mathbb{R}; m(h)d\omega dh)} \quad \text{for all } \Phi, \Phi' \in \mathbf{S}(V). \quad (100)$$

Let us now prove that $\overline{K'_\lambda \mathbf{S}(V)} + iK'_\lambda \mathbf{S}(V) = L^2(\mathbb{S}^2 \times \mathbb{R}; m(h)d\omega dh)$. This is an immediate consequence of the following lemma, whose proof can be found in this appendix.

Lemma A.0.6. *Referring to coordinates $(\omega, \ell) \in \mathbb{S}^2 \times \mathbb{R}$ to define $C_0^\infty(\mathbb{S}^2 \times \mathbb{R})$, the space $K'_\lambda C_0^\infty(\mathbb{S}^2 \times \mathbb{R}) + iK'_\lambda C_0^\infty(\mathbb{S}^2 \times \mathbb{R})$ is dense in $L^2(\mathbb{S}^2 \times \mathbb{R}; d\mu_{\mathbb{S}^2} m(h)dh)$.*

If we define $\mu_{\lambda'}(\Phi, \Phi') \doteq Re\langle K'_\lambda \Phi, K'_\lambda \Phi' \rangle_{H'_\lambda}$ for $\Phi, \Phi' \in \mathbf{S}(V)$, it turns out that $\mu_{\lambda'}$ is a real scalar product (strictly positive) on $\mathbf{S}(V)$. Furthermore it satisfies $|\sigma_V(\Phi, \Phi')|^2 \leq 4\mu_{\lambda'}(\Phi, \Phi)\mu_{\lambda'}(\Phi', \Phi')$ for all $\Phi, \Phi' \in \mathbf{S}(V)$ in view of (100) and the Cauchy-Schwarz inequality for $\langle \cdot, \cdot \rangle_{H'}$. Therefore, due to Proposition 3.1 in [KW91], there is a quasifree state on $\mathscr{W}(V)$ (the unique satisfying the identity in (40) with μ_λ replaced for $\mu_{\lambda'}$) with one-particle space structure (K'_λ, H'_λ) . To conclude the proof of (a) we have to prove that (K'_λ, H'_λ) is unitarily equivalent to (K_λ, H_λ) (so that λ' coincides to λ). Changing variables in the two-point function of state λ and obtaining:

$$\lambda(\Phi, \Phi') = \lim_{\epsilon \rightarrow 0^+} -\frac{1}{4\pi} \int_{\mathbb{S}^2 \times \mathbb{R}^2} d\omega d\ell d\ell' \frac{\Phi(\omega, \ell)\Phi'(\omega, \ell')}{\left(\sinh\left(\frac{\ell-\ell'}{2}\right) - i\epsilon \cosh(\ell/2) \cosh(\ell'/2)\right)^2}, \quad (101)$$

we can exploit the following technical result proved below in this appendix.

Lemma A.0.7. *Let N be a smooth Riemannian manifold and let μ be the natural measure induced by the metric on N . If $f, h \in C_0^\infty(N)$ with $h > 0$, $df \neq 0$, $fdh - hdf \neq 0$ on N , then:*

$$\lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu = \lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{(f(q) \pm i\epsilon)^2} d\mu \quad \text{for every } g \in C_0^\infty(N), \quad (102)$$

and both limits do exist and are finite.

Assuming $N = \mathbb{S}^2 \times \mathbb{R}^2$ equipped with the natural product metric (referred to coordinates (ℓ, ℓ') on \mathbb{R}^2) and taking $f(\omega, \ell, \ell') = \sinh((\ell - \ell')/2)$, so that $df \neq 0$ everywhere on N , and $h(\omega, \ell, \ell') = \cosh(\ell/2) \cosh(\ell'/2)$ which is strictly positive and $f dh - h df \neq 0$ as it can be checked by direct inspection, Lemma A.0.7 and (101) entail:

$$\lambda(\Phi, \Phi') = \lim_{\epsilon \rightarrow 0_+} -\frac{1}{4\pi} \int_{\mathbb{S}^2 \times \mathbb{R}^2} d\omega d\ell d\ell' \frac{\Phi(\omega, \ell) \Phi'(\omega, \ell')}{(\sinh(\frac{\ell - \ell'}{2}) - i\epsilon)^2} \quad \text{if } \Phi, \Phi' \in C_0^\infty(\mathbb{S}^2 \times \mathbb{R}). \quad (103)$$

The ℓ -Fourier transform of the distribution $1/(\sinh(\ell/2) - i0_+)^2$ is $-2\sqrt{2\pi}m(h)$ with δ defined in (55), so that the convolution theorem applied to the right hand side of (103) proves that, if $\Phi, \Phi' \in C_0^\infty(\mathbb{S}^2 \times \mathbb{R})$:

$$\langle K_\lambda \Phi, K_\lambda \Phi' \rangle_{H_\lambda} = \int_{\mathbb{S}^2 \times \mathbb{R}^2} m(h) d\omega dh \overline{\widetilde{\Phi}(\omega, h)} \widetilde{\Phi}'(\omega, h) = \langle K'_\lambda \Phi, K'_\lambda \Phi' \rangle_{H'_\lambda}. \quad (104)$$

We have found that the \mathbb{R} -linear, bijective map $U_0 : K'_\lambda(C_0^\infty(\mathbb{S}^2 \times \mathbb{R})) \rightarrow K_\lambda(C_0^\infty(\mathbb{S}^2 \times \mathbb{R}))$ which associates² $K'_\lambda \Phi$ to $K_\lambda \Phi$ is isometric. On the other hand the following lemma (proved below in the appendix) holds.

Lemma A.0.8. *Referring to coordinates $(\omega, \ell) \in \mathbb{S}^2 \times \mathbb{R}$ to define $C_0^\infty(\mathbb{S}^2 \times \mathbb{R})$, the space $K_\lambda(C_0^\infty(\mathbb{S}^2 \times \mathbb{R})) + iK'_\lambda(C_0^\infty(\mathbb{S}^2 \times \mathbb{R}))$ is dense in H_λ .*

Since also $K'_\lambda C_0^\infty(\mathbb{S}^2 \times \mathbb{R}) + iK_\lambda C_0^\infty(\mathbb{S}^2 \times \mathbb{R})$ is dense in H'_λ (by Lemma A.0.6), Lemma A.1 in [KW91] implies that U_0 linearly and continuously extends to a unique (\mathbb{C} -linear) Hilbert space isomorphism $U : H'_\lambda \rightarrow H_\lambda$. By construction $UK'_\lambda = K_\lambda$, so that the two one-particle space structures (K'_λ, H'_λ) and (K_λ, H_λ) are unitarily equivalent.

Let us prove (a) and (b) in the thesis of Theorem 5.1.1. Since $X = \partial_\ell$, it holds $(\beta_\tau^X \Phi)(\omega, \ell) = \Phi(\omega, \ell - \tau)$. Thus, referring to the one-particle space structure (K'_λ, H'_λ) , the action of β_τ^X is equivalent to the appearance of a phase $e^{ih\tau}$ in front of the $K'_\lambda \Phi$ by standard properties of the Fourier transform. Defining $(V_\tau^{(\lambda')})\Psi(\omega, h) \doteq e^{ih\tau}\Psi(\omega, h)$, it holds $V_\tau^{(\lambda')} K'_\lambda \Phi = K'_\lambda \beta_\tau^X \Phi$. In view of Stone's theorem and a few of elementary observations, \hat{h} is the generator of $V^{(\lambda')}$. As $\mu_\lambda(\beta_\tau^X \Psi, \beta_\tau^X \Psi) = \langle V_t^{(\lambda')} K'_\lambda \Phi, V_t^{(\lambda')} K'_\lambda \Phi \rangle_{H'_\lambda} = \langle K'_\lambda \Phi, K'_\lambda \Phi \rangle_{H'_\lambda} = \mu_\lambda(\Phi, \Phi)$ because phases cancel, λ turns out to be α^X -invariant due to (40), (41) and (53). Passing to the unitarily equivalent one-particle space structure (K_λ, H_λ) , the result remains unchanged if replacing $V^{(\lambda')}$ with the corresponding unitary $V^{(\lambda)}$. Passing to the second-quantization $U_\tau^{(\lambda)}$ of $V_\tau^{(\lambda)}$ (assuming $U_\tau^{(\lambda)}\Psi_\lambda = \Psi_\lambda$) implements α^X as follows from (57) by standard procedures. To conclude the proof of (b), for $\beta = 2\pi$, define the antilinear map $j : K'_\lambda \mathcal{S}(V) \rightarrow \overline{K'_\lambda \mathcal{S}(V)}$ such that $(j\widetilde{\Phi})(\omega, h) \doteq -e^{-\beta h/2} \overline{\widetilde{\Phi}(\omega, -h)}$. The definition of j is well-posed since $\overline{\widetilde{\Phi}(\omega, -h)} = \widetilde{\Phi}(\omega, h)$ if $\Phi \in \mathcal{S}(V)$ because Φ is real-valued and $K'_\lambda(\mathcal{S}) \subset \text{Dom}(e^{-\frac{1}{2}\beta\hat{h}})$ by direct inspection. Notice that

²This map is well defined since $K_\omega \Phi = K_\omega \Phi_1$ implies $\Phi = \Phi_1$ as the \mathbb{R} -linear map $K_\omega : \mathcal{S} \rightarrow H_\omega$ is injective for every one-particle space structure (K_ω, H_ω) of every quasifree state ω of a Weyl algebra of a symplectic space (S, σ) , because $\text{Im}\langle K_\omega \Phi, K_\omega \Phi_1 \rangle$ is proportional to $\sigma(\Phi, \Phi_1)$ that is nondegenerate by hypotheses.

$jj = I$ and j commutes with $V_t^{(\lambda)}$. We can extend j to the whole H'_λ by antilinearity and continuity eventually obtaining an antilinear operator $j : H'_\lambda \rightarrow H'_\lambda$ with $jj = I$ and the following facts are verified: (i) $K'_\lambda(\mathbb{S}) \subset \text{Dom}\left(e^{-\frac{1}{2}\beta\hat{h}}\right)$, (ii) $[j, V_t^{(\lambda)}] = 0$ if $t \in \mathbb{R}$, (iii) $e^{-\frac{1}{2}\beta\hat{h}}K'_\lambda\psi = -jK'_\lambda\psi$ if $\psi \in K'_\lambda\mathbb{S}(V)$. By the standard procedure of second-quantization with Fock space \mathcal{H}_λ , j uniquely determines an antiunitary operator J leaving fixed Ψ_λ . By standard procedures using (i),(ii) and (iii) and referring to J and $U_t^{(\lambda)}$, one easily proves that the requirements (1), (2) and (3) in the (second equivalent) definition of KMS state presented at the end of this appendix are satisfied (see Sec. 3.2 in [KW91] and references therein for details) so that λ turns out to be KMS at the inverse temperature $\beta = 2\pi$ with respect to α^X . The remaining part of the statement (a) arises by the link between KMS condition and modular group outlined at the beginning of Section 5. (b) arises immediately defining $V_\tau^{(\lambda)}$ as the unitary corresponding to $V_\tau^{(\lambda)}$ in the structure (K_λ, H_λ) . \square

Proof of the Lemma A.0.6. $\mathbb{S}(V) + i\mathbb{S}(V) \supset C_0^\infty(\mathbb{S}^2 \times \mathbb{R}) + iC_0^\infty(\mathbb{S}^2 \times \mathbb{R}) \doteq C_0^\infty(\mathbb{S}^2 \times \mathbb{R}; \mathbb{C})$. Since the latter is dense in $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ and the Fourier transform is continuous with respect to the topology of $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ and leaves that space fixed, we conclude that $\overline{K'_\lambda\mathbb{S}(V) + iK'_\lambda\mathbb{S}(V)}$ (where the closure is that of the topology of $L^2(\mathbb{S}^2 \times \mathbb{R}; m(h)d\omega dh)$) includes the subspace $\mathcal{S}_0 \doteq K'_\lambda C_0^\infty(\mathbb{S}^2 \times \mathbb{R}) + iK'_\lambda C_0^\infty(\mathbb{S}^2 \times \mathbb{R})$, which is dense in $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ with respect the $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ topology (where, now $\mathbb{S}^2 \times \mathbb{R}$ is referred to the coordinates (ω, h)). Since the convergence in the $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ topology implies that in the topology of each $L^2(\mathbb{S}^2 \times \mathbb{R}; (1+|h|^n)d\omega dh)$ for $n = 0, 1, 2, \dots$, and $|m(h)| \leq C(1+h^2)$ for some $C \geq 0$ and all $h \in \mathbb{R}$, we conclude that \mathcal{S}_0 is dense in $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R})$ also referring to the topology of $L^2(\mathbb{S}^2 \times \mathbb{R}; m(h)d\omega dh)$. In particular it has to hold $\mathcal{S}(\mathbb{S}^2 \times \mathbb{R}) \subset \overline{K'_\lambda\mathbb{S}(V) + iK'_\lambda\mathbb{S}(V)}$ and thus $C_0^\infty(\mathbb{S}^2 \times \mathbb{R}; \mathbb{C}) \subset \overline{K'_\lambda\mathbb{S}(V) + iK'_\lambda\mathbb{S}(V)}$ where, again, $\mathbb{S}^2 \times \mathbb{R}$ is referred to the coordinates (ω, h) . To conclude, it is sufficient to establish that $C_0^\infty(\mathbb{S}^2 \times \mathbb{R}; \mathbb{C})$ is dense in $L^2(\mathbb{S}^2 \times \mathbb{R}; m(h)d\omega dh)$. This is a trivial consequence of the fact that $C_0^\infty(\mathbb{S}^2 \times (0, +\infty); \mathbb{C})$ is dense in $L^2(\mathbb{S}^2 \times (0, +\infty); d\omega dx)$, passing from the variable x to the variable h such that $x(h) = \int_0^h m(h')dh'$ and noticing that $g \in C_0^\infty(\mathbb{S}^2 \times (0, +\infty); \mathbb{C})$ if and only if $\mathbb{R} \ni h \mapsto g(x(h))$ belongs to $C_0^\infty(\mathbb{S}^2 \times \mathbb{R}; \mathbb{C})$. \square

Proof of Lemma A.0.7. With our hypothesis on f , all the nonempty sets of the form $\Sigma_c := \{x \in N \mid f(x) = c\}$ are embedded submanifolds of N . Thus, in a neighborhood Ω_p of every $p \in N$, there is a coordinate patch $\psi : \Omega_p \ni q \mapsto (y^1(q), y^2(q), \dots, y^n(q))$ satisfying $f \circ \psi^{-1} = y^1$ so that $\Sigma_c \cap \Omega_p$ is made of the points $q \in N$ with $y^1(q) = c$. We henceforth assume that $\Sigma_0 \neq \emptyset$ and we restrict to work in Ω_p . Distributions $\delta(f)$ and $\mathcal{P}(1/f)$ in $\mathcal{D}'(\Omega_p)$ can be defined as follows, for every $g \in C_0^\infty(\Omega_p)$:

$$\langle \delta(f), g \rangle \doteq \int_{\mathbb{R}^{n-1}} g(0, y^2, \dots, y^n) \frac{d\mu_{\Sigma_c}}{|df|},$$

$$\left\langle \mathcal{P}\left(\frac{1}{f}\right), g \right\rangle \doteq \lim_{\epsilon \rightarrow 0^+} \left(\int_{-\infty}^{-\epsilon} \frac{dy^1}{y^1} + \int_{\epsilon}^{+\infty} \frac{dy^1}{y^1} \right) \int_{\mathbb{R}^{n-1}} g(y^1, \dots, y^n) \frac{d\mu_{\Sigma_c}}{|df|},$$

where μ_{Σ_0} is the measure induced on Σ_0 by the metric of N , and $|df|$ the norm of df referred to the metric. The reader can verify that the definitions are well-posed, reduce to the standard ones for $N = \mathbb{R}$ (equipped with the natural Euclidean metric) and are independent from the used coordinate patch on Ω_p verifying the hypotheses. Finally, working in the said coordinate patch and through a very straightforward generalization of the proof of Sochockij's formulas as in [V179], one extends these formulas to our case. More precisely, it turns out that: *Assuming $df \neq 0$ on N , in a sufficiently small open neighborhood Ω_p of any point $p \in N$:*

$$\lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{f(q) \pm i\epsilon h(q)} d\mu = \left\langle \mathcal{P}\left(\frac{1}{f}\right), g \right\rangle \mp i\pi \langle \delta(f), g \rangle \quad \text{for every } g \in C_0^\infty(\Omega_p), \quad (105)$$

where the assigned function $h : N \rightarrow \mathbb{R}$ is every smooth function that is strictly positive on Ω_p . In particular one can choose $h \equiv 1$. The obtained result entails that, in weak sense:

$$\frac{1}{f(q) \pm i\epsilon h(q)} \rightarrow \mathcal{P}(1/f) \mp i\pi \delta(f) \leftarrow \frac{1}{f(q) \pm i\epsilon} \quad \text{as } \epsilon \rightarrow 0^+.$$

As weakly convergent sequences of distributions can be derived term-by-term obtaining again a sequence that weakly converges to a distribution, using again coordinates y^1, \dots, y^n , taking the y^1 derivative – taking into account some further smooth factors arising from the density function in the integration measure – and coming back to general coordinates on Ω_p , we conclude that, if f and h satisfy our hypotheses and $g \in C_0^\infty(\Omega_p)$:

$$\begin{aligned} \lim_{\epsilon \rightarrow 0^+} \left(\int_N \frac{g(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu \pm i\epsilon \int_N \frac{g(q) \partial_{y^1} h(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu \right) \\ = \lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{(f(q) \pm i\epsilon)^2} d\mu \in \mathbb{C}. \end{aligned} \quad (106)$$

Notice that the result above not only states that the two sides coincide, but it even states that both limit exist and are finite. The second integral in the left-hand side can be re-arranged as:

$$\int_N \frac{g(q) (\partial_{y^1} h(q)) / h(q)^{1/2}}{((f(q)/h(q)) \pm i\epsilon)^2} d\mu$$

The existence of the limit in the right-hand side in (106) for $f(q)$ replaced by $f(q)/h(q)$, and $g(q)$ replaced by $g(q) (\partial_{y^1} h(q)) / h(q)^{1/2}$ implies that $\lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q) \partial_{y^1} h(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu$ exists in \mathbb{C} , provided that $d(f/g) \neq 0$, i.e. $hdf - fdh \neq 0$, shrinking the neighborhood Ω_p if necessary. Consequently, under these hypotheses, the second term in the left-hand side of (106) vanishes as $\epsilon \rightarrow 0^+$ in view of the overall factor ϵ . We conclude that: *if $df \neq 0$, $h > 0$, $hdf - fdh \neq 0$ on N then, in a sufficiently small neighborhood Ω_p of every $p \in N$, one has:*

$$\lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu = \lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{(f(q) \pm i\epsilon)^2} d\mu \in \mathbb{C}, \quad \text{for every } g \in C_0^\infty(\Omega_p) \quad (107)$$

Now, with the given hypotheses on f and h on N , we extend the result to the whole manifold N . Take $g \in C_0^\infty(N)$. For every $p \in \text{supp}(g)$ there is a neighborhood Ω_p where

(107) can be applied. As $\text{supp}(g)$ is compact, we can extract a finite open covering $\{\Omega_{p_i}\}_{i=1,\dots,N}$ of $\text{supp}(g)$ and construct a partition of the unit $\{\psi_i\}_{i=1,\dots,N}$ subordinate to that finite covering. Notice that in each Ω_{p_i} the identity (107) holds because we can construct the relevant coordinate systems therein. Then, since the sum over i is finite so that it can be interchanged with the integral symbol, and all the limits are finite, we achieve the thesis:

$$\begin{aligned} \lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu &= \lim_{\epsilon \rightarrow 0^+} \sum_{i=1}^N \int_{\Omega_{p_i}} \frac{g(q)\psi_i(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu \\ &= \sum_{i=1}^N \lim_{\epsilon \rightarrow 0^+} \int_{\Omega_{p_i}} \frac{g(q)\psi_i(q)}{(f(q) \pm i\epsilon h(q))^2} d\mu = \sum_{i=1}^N \lim_{\epsilon \rightarrow 0^+} \int_{\Omega_{p_i}} \frac{g(q)\psi_i(q)}{(f(q) \pm i\epsilon)^2} d\mu \\ &= \lim_{\epsilon \rightarrow 0^+} \sum_{i=1}^N \int_{\Omega_{p_i}} \frac{g(q)\psi_i(q)}{(f(q) \pm i\epsilon)^2} d\mu = \lim_{\epsilon \rightarrow 0^+} \int_N \frac{g(q)}{(f(q) \pm i\epsilon)^2} d\mu. \end{aligned}$$

□

Proof of Lemma A.0.8. Let $\Phi \in \mathcal{S}(V)$, then Φ is smooth in coordinates $(\omega, u) \in \mathbb{S}^2 \times \mathbb{R}_+$, approaching $u = 0$ it vanishes with order $O(u)$ uniformly in the angles, Φ vanishes also if $u > u_0 > 0$ for some $u_0 < 1$, finally $\partial_u \Phi_u$ is bounded. Define $\chi : \mathbb{R} \rightarrow [0, 1]$ such that $\chi \in C^\infty(\mathbb{R})$ and $\chi(u) = 1$ for $u \geq 2$ whereas $\chi(u) = 0$ for $u \leq 1$. Next define $\chi_n(u) \doteq \chi(nu)$ for $n = 1, 2, 3, \dots$. By direct inspection one sees that the sequence of functions $\Phi_n \doteq \chi_n \cdot \Phi$ satisfies the following as $n \rightarrow +\infty$: (a) $|\Phi_n| \leq |\Phi|$ and $\Phi_n \rightarrow \Phi$ pointwisely, (b) $\Phi_n \rightarrow \Phi$ in the sense of $L^2(\mathbb{S}^2 \times \mathbb{R}_+, d\omega du)$, (c) $\partial_u \Phi_n \rightarrow \partial_u \Phi$ in the sense of $L^2(\mathbb{S}^2 \times \mathbb{R}_+, d\omega du)$. Passing to the u -Fourier transforms, from (b) and (c) and Cauchy-Schwarz inequality, one has that $\mathcal{K}_\lambda \Phi_n \rightarrow \mathcal{K}_\lambda \Phi$ in the topology of $L^2(\mathbb{S}^2 \times \mathbb{R}_+, k^n d\omega dk)$ with $n = 0, 1, 2$. This entails the thesis immediately because, by construction, $\Phi_n \in C_0^\infty(\mathbb{S}^2 \times \mathbb{R})$, where now \mathbb{R} is referred to the coordinate ℓ . □

Proof of Lemma 5.2.1. We adopt the same conventions and notations as in the proof of Lemma 3.1.5. For the fixed Φ and ϕ as in the hypotheses, the sequence $\{\phi_n\}_{n \in \mathbb{N}} \subset \mathcal{S}_m(D)$ is individuated by the sequence of corresponding smooth compactly-supported Cauchy data on Σ_D , $\{(\phi_n(1, \cdot), \partial_t \phi_n(1, \cdot))\}_{n \in \mathbb{N}}$ with: (a) $\phi_n(1, \cdot) \rightarrow \phi(1, \cdot)$ in the topology of $H^1(\Sigma_D)$ and (b) $\partial_t \phi_n(1, \cdot) \rightarrow \partial_t \phi(1, \cdot)$ in the topology of $H^0(\Sigma) \equiv L^2(\Sigma_D, dx)$. The former sequence exists because $\phi(1, \cdot) \in H_0^1(\Sigma_M)$ and $C_0^\infty(\Sigma)$ is dense in $H_0^1(\Sigma_D)$. The latter sequence exists because $C_0^\infty(\Sigma_D)$ is dense in $H^0(\Sigma_D)$. Passing to the \mathbf{x} Fourier-Plancherel transform, the fact that $\{(\phi_n(1, \cdot), \partial_t \phi_n(1, \cdot))\}_{n \in \mathbb{N}}$ converges to the pair $(\phi(1, \cdot), \partial_t \phi(1, \cdot))$ in the said topologies entails that $\hat{\phi}_n \rightarrow \hat{\phi}$ in the topology of $L^2(\mathbb{R}^3, d\mathbf{k})$ as it follows from (24). Indeed, denoting by \mathcal{F} the standard Fourier-Plancherel transformation on \mathbb{R}^3 , (24) and the Cauchy-Schwarz identity imply that the nonnegative number $\|\hat{\phi}_n - \hat{\phi}\|_{L^2(\mathbb{R}^3, d\mathbf{k})}^2/2$ is bounded by

$$\int_{\mathbb{R}^3} \sqrt{m^2 + \mathbf{k}^2} |\mathcal{F}(\phi_n(1, \cdot))(\mathbf{k}) - \mathcal{F}(\phi(1, \cdot))(\mathbf{k})|^2 d\mathbf{k} + \int_{\mathbb{R}^3} \frac{|\mathcal{F}(\partial_t \phi_n(1, \cdot))(\mathbf{k}) - \mathcal{F}(\partial_t \phi(1, \cdot))(\mathbf{k})|^2}{\sqrt{m^2 + \mathbf{k}^2}} d\mathbf{k}$$

$+ \|\phi_n(1, \cdot) - \phi(1, \cdot)\|_{L^2(\Sigma_D, dx)} \|\partial_t \phi_n(1, \cdot) - \partial_t \phi(1, \cdot)\|_{L^2(\Sigma_D, dx)} \rightarrow 0$ for $m > 0$ as $n \rightarrow +\infty$,

because $(\phi_n(1, \cdot), \partial_t \phi_n(1, \cdot)) \rightarrow (\phi(1, \cdot), \partial_t \phi(1, \cdot))$ in the said Sobolev topologies. If $m = 0$, the only apparent problem concerns the second integral. The obstruction can be avoided using a sequence of *uniformly bounded* smooth compactly supported (in Σ_D) functions $\partial_t \phi_n(1, \cdot)$ converging almost everywhere to $\partial_t \phi(1, \cdot)$ (and thus also in $L^2(\Sigma, dx)$ since Σ_D has finite measure). It exists as a consequence of Luzin's and Stone-Weierstrass theorems, and the sequence of the Fourier transforms $\mathcal{F}(\partial_t \phi_n(1, \cdot))$ converges to $\mathcal{F}(\partial_t \phi(1, \cdot))$ both pointwisely (due to Lebesgue's dominated convergence theorem) and in the $L^2(\mathbb{R}^3, dk)$ sense. Since $\mathcal{F}(\partial_t \phi(1, \cdot))$ and the $\mathcal{F}(\partial_t \phi_n(1, \cdot))$ are uniformly bounded by a constant independent from n and $1/\|\mathbf{k}\|$ being $d\mathbf{k}$ -integrable, one easily proves that the second integral above vanishes for $n \rightarrow +\infty$.

We have so far obtained that, for $m \geq 0$, $\tilde{\mathcal{K}}_m L_{\tilde{D}D}^m(\phi_n) \rightarrow \tilde{\mathcal{K}}_m \phi$. To prove that the other convergence property in the thesis holds, we notice that the thesis of Lemma 4.2.4, in particular the identity $\int_{\mathbb{R}^3} |\hat{\phi}(\mathbf{k})|^2 d\mathbf{k} = \langle \mathcal{K}_\lambda \Phi, \mathcal{K}_\lambda \Phi \rangle_{\mathcal{H}_\lambda}$ is valid for a generic smooth-in- \bar{D} function ϕ satisfying the Klein-Gordon equation in M , such that $\hat{\phi}$ defined as in (24) is $L^2(\mathbb{R}^3, d\mathbf{k})$ and $\Phi = L_{V\tilde{D}}^m \phi$ defines an element of $\mathcal{S}(V)$ as in our case, since just those hypotheses were used in the proof. Therefore, we can re-adapt that proof to our case replacing ϕ for $\phi - \phi_n$ and Φ for $\Phi - L_{V\tilde{D}}^m(\phi_n)$, obtaining $\|\mathcal{K}_\lambda \Phi - \mathcal{K}_\lambda L_{V\tilde{D}}^m(\phi_n)\|_{\mathcal{H}_\lambda}^2 = \|\hat{\phi}_n - \hat{\phi}\|_{L^2(\mathbb{R}^3, d\mathbf{k})}^2 \rightarrow 0$ for $n \rightarrow +\infty$. The last statement in the lemma is trivially true noticing that $\tilde{\mathcal{H}}_m \supset \mathcal{H}_m$ are closed subspaces of $L^2(\mathbb{R}^3, d\mathbf{k})$ and thus $\mathcal{K}_m = \tilde{\mathcal{K}}_m L_{\tilde{D}D}^m$ is valid by the definitions of $\tilde{\mathcal{K}}_m$ and \mathcal{K}_m . Since convergence property $\tilde{\mathcal{K}}_m L_{\tilde{D}D}^m \phi_n \rightarrow \tilde{\mathcal{K}}_m \phi$ shows that \mathcal{H}_m is dense in $\tilde{\mathcal{H}}_m$, we conclude that $\tilde{\mathcal{H}}_m = \mathcal{H}_m$. \square

On KMS states

Comparing Definition 5.3.1 and Proposition 5.3.7 in [BR96II] we say that a state ω on a C^* -algebra \mathcal{A} is **KMS at inverse temperature** $\beta \in \mathbb{R}$ with respect to a one-parameter group of $*$ -automorphisms $\{\alpha_t\}_{t \in \mathbb{R}}$ if the function $\mathbb{R} \ni t \mapsto F_{A,B}^{(\omega)}(t) \doteq \omega(A\alpha_t(B))$ satisfies the following three requirements for every pair $A, B \in \mathcal{A}$. **(a)** It extends to a continuous complex function $F_{A,B}^{(\omega)} = F_{A,B}^{(\omega)}(z)$ on $\overline{D_\beta} \doteq \{z \in \mathbb{C} \mid 0 \leq \text{Im}z \leq \beta\}$ if $\beta \geq 0$, or $\overline{D_\beta} \doteq \{z \in \mathbb{C} \mid \beta \leq \text{Im}z \leq 0\}$ if $\beta \leq 0$; **(b)** the extension is analytic in the interior of $\overline{D_\beta}$; **(c)** the **KMS condition** holds: $F_{A,B}^{(\omega)}(t + i\beta) = \omega(\alpha_t(B)A)$ for all $t \in \mathbb{R}$. Another definition is the following. We say that a state ω on a C^* -algebra \mathcal{A} is **KMS at inverse temperature** $\beta \in \mathbb{R}$ with respect to a one-parameter group of $*$ -automorphisms $\{\alpha_t\}_{t \in \mathbb{R}}$ if its GNS triple $(\mathcal{H}_\omega, \pi_\omega, \Psi_\omega)$ satisfies the following three requirements. **(1)** ω is α -invariant and the unique unitary group $\mathbb{R} \ni t \mapsto U_t$ which leaves Ψ_ω invariant implementing α (i.e. $\pi_\omega(\alpha_t(A)) = U_t \pi_\omega(A) U_t^*$ if $A \in \mathcal{A}$ and $t \in \mathbb{R}$) is strongly continuous. **(2)** $\pi_\omega(\mathcal{A}) \Psi_\omega \subset \text{Dom}(e^{-\beta H/2})$ where $e^{itH} = U_t$, for $t \in \mathbb{R}$, with H self-adjoint in \mathcal{H}_ω . **(3)** There exists $J : \mathcal{H}_\omega \rightarrow \mathcal{H}_\omega$ antilinear with $JJ = I$ such that: $J e^{-itH} = e^{-itH} J$, for all $t \in \mathbb{R}$, and $e^{-\beta H/2} \pi_\omega(A) \Psi_\omega = J \pi_\omega(A^*) \Psi_\omega$ for all $A \in \mathcal{A}$.

These two definitions are equivalent as we prove here. A state satisfying (a), (b), (c)

is $\{\alpha_t\}_{t \in \mathbb{R}}$ -invariant [BR96II] and fulfils the conditions (1), (2) and (3) due to Theorem 6.1 in [Hu72]. Conversely, consider a state ω on \mathcal{A} fulfilling the conditions (1), (2) and (3). When A and B are entire analytic elements of \mathcal{A} (see [BR96II]), $\mathbb{R} \ni t \mapsto F_{A,B}^{(\omega)}(t)$ uniquely extends to an analytic function on the whole \mathbb{C} and thus (a) and (b) are true. (1), (2), (3) and $e^{zH}\Psi_\omega = \Psi_\omega$, for all $z \in \overline{D_\beta}$ (following from (2) and (3)) entail (c), too: $\omega(\alpha_t(B)A) = \langle \Psi_\omega, U_t \pi_\omega(B) U_t^* \pi_\omega(A) \Psi_\omega \rangle = \langle \pi_\omega(B^*) \Psi_\omega, U_t^* \pi_\omega(A) \Psi_\omega \rangle = \langle J U_t^* \pi_\omega(A) \Psi_\omega, J \pi_\omega(B^*) \Psi_\omega \rangle = \langle U_t^* e^{-\beta H/2} \pi_\omega(A^*) \Psi_\omega, e^{-\beta H/2} \pi_\omega(B) \Psi_\omega \rangle = \langle \Psi_\omega, \pi_\omega(A) e^{i(t+i\beta)H} \pi_\omega(B) e^{-i(t+i\beta)H} \Psi_\omega \rangle = F_{A,B}^{(\omega)}(t+i\beta)$. The validity of conditions (a), (b) and (c) for entire analytic elements $A, B \in \mathcal{A}$ implies the validity for all $A, B \in \mathcal{A}$, as established in [BR96II] (compare Definition 5.3.1 and Proposition 5.3.7 therein).

References

- [Au00] J.-P. Aubin, *Applied Functional Analysis*. Second edition, Wiley-Interscience, New York (2000),
- [BGP07] C. Bär, N. Ginoux, F. Pfäffle, *Wave equations on Lorentzian manifolds and quantization* ESI Lectures in Mathematics and Physics, the European Mathematical Society Publishing (2007)
- [BJL02] H. Baumgärtel, M. Jurke, and F. Lledó, *Twisted duality of the CAR-Algebra*, J. Math. Phys. **43**, (2002), 41584179.
- [BMRW09] M. Bischoff, D. Meise, K.-H. Rehren, I. Wagner, *Conformal quantum field theory in various dimensions* arXiv:0908.3391v1 [math-ph]
- [BW75] J. Bisognano and E. H. Wichmann, *On the duality condition for a Hermitian quantum field*, J. Math. Phys. **16** (1975) 985-1007.
- [BW76] J. Bisognano and E. H. Wichmann, *On the duality condition for quantum fields*, J. Math. Phys. **17** (1976) 303321.
- [Bo92] H. J. Borchers, *The CPT-theorem in two-dimensional theories of local observables*, Commun. Math. Phys. **143**, 315-332 (1992).
- [Bo00] H. J. Borchers, *On revolutionizing quantum field theory with Tomita's modular theory*, J. Math. Phys. **41**, (2000), no. 6, 3604-3673.
- [BY99] H. J. Borchers and J. Yngvason, *Modular groups of quantum fields in thermal states*, J. Math. Phys. **40**, (1999), no. 2, 601-624.
- [BR96I] O. Bratteli, D. W. Robinson, *Operator Algebras and Quantum Statistical Mechanics*. Vol.1, Springer Berlin, Germany (1996)
- [BR96II] O. Bratteli, D. W. Robinson, *Operator Algebras and Quantum Statistical Mechanics*. Vol. 2, Springer Berlin, Germany (1996)

- [BGL93] R. Brunetti, D. Guido and R. Longo, *Modular structure and duality in conformal quantum field theory*, Comm. Math. Phys. **156**, (1993), 201-219.
- [BGL95] R. Brunetti, D. Guido and R. Longo, *Group cohomology, modular theory and space-time symmetries*, Rev. Math. Phys. **7**, (1995), 57-71.
- [BGL02] R. Brunetti, D. Guido and R. Longo, *Modular localization and Wigner particles*, Rev. Math. Phys. **14**, (2002), 759-785.
- [Bu77] D. Buchholz, *On the structure of local quantum fields with non-trivial interaction*, in Proc. Intern. Conf. Operator Algebras, Ideals, and Their Applications in Physics, ed. H. Baumgärtel (Teubner, 1977), 146153.
- [BDFS] D. Buchholz, O. Dreyer, M. Florig and S. Summers, *Geometric modular action and spacetime symmetry groups*, Rev. Math. Phys. **12**, (2000), 475-560.
- [BMS] D. Buchholz, J. Mund and S. Summers, *Covariant and quasi-covariant quantum dynamics in Robertson-Walker spacetimes*, Classical Quantum Gravity **19**, (2002), 6417-6434.
- [CH09] H. Casini and M. Huerta, *Reduced density matrix and internal dynamics for multi- component regions*, Class. Quant. Grav. **26**, (2009) 185005
- [Co74] A. Connes, *Une classification des facteurs de type III*, Ann. Ec. Norm. Sup., **6**, (1974) 415-445.
- [Du73] J. J. Duistermaat, *Fourier Integral Operators*. Courant Institute of Mathematical Sciences, New York (1973)
- [DMP06] C. Dappiaggi, V. Moretti and N. Pinamonti, *Rigorous steps towards holography in asymptotically flat spacetimes*, Rev. Math. Phys. **18**, (2006), 349 [gr-qc/0506069]
- [DMP09] C. Dappiaggi, V. Moretti and N. Pinamonti, *Cosmological horizons and reconstruction of quantum field theories*, Commun. Math. Phys. **285** (2009), 1129 arXiv:0712.1770 [gr-qc]
- [DMP09b] C. Dappiaggi, V. Moretti and N. Pinamonti, *Distinguished quantum states in a class of cosmological spacetimes and their Hadamard property*, J. Math. Phys. **50** (2009), 062304 arXiv:0812.4033 [gr-qc]
- [DPP10] C. Dappiaggi, N. Pinamonti and M. Porrmann, *Local causal structures, Hadamard states and the principle of local covariance in quantum field theory*. In print in Commun. Math. Phys. [arXiv:1001.0858]
- [DMS97] P. Di Francesco, P. Mathieu and David Sénéchal, *Conformal Field Theory*. Springer, Berlin (1997)

- [FG89] F. Figliolini and D. Guido, *The Tomita operator for the free scalar field*, Ann. Inst. Henri Poincaré Phys. Theor. **51**, (1989) 419435.
- [Fr75] F. G. Friedlander, *The Wave Equation on a Curved Space-Time*, Cambridge University Press, London (1975)
- [FJ75] F. G. Friedlander and M. Joshi, *Introduction to the Theory of Distributions*, 2nd edition Cambridge University Press, London (2003)
- [GR95] I. S. Gradshteyn, I. M. Ryzhik, *Table of Integrals, Series, and Products*, Fifth Edition, Academic Press, (1995).
- [GS94] A. Grigis and J. Sjöstrand, *Microlocal Analysis for Differential Operators*, London Mathematical Society, Lecture Notes Series, vol. 196, Cambridge University Press, (1994).
- [Gu08] D. Guido, *Modular theory for the von Neumann algebras of local quantum physics*, arXiv:0812.1511.
- [Ha92] R. Haag, *Local Quantum Physics: Fields, Particles, Algebras*, Second Revised and Enlarged Edition, Springer (1992).
- [Ho00] S. Hollands, *Aspects of Quantum Field Theory in Curved Spacetime*. Ph.D.thesis, University of York, 2000, unpublished.
- [Hö94] L. Hörmander, *The Analysis of Linear Partial Differential Operators III Pseudo-Differential Operators*, Springer-Verlag, Berlin (1994)
- [Hö96] L. Hörmander, *Lectures on nonlinear hyperbolic differential equations*, Mathématiques & Applications, Vol. 26, Springer-Verlag, Berlin (1996)
- [HL95] P. D. Hislop and R. Longo, *Modular Structure of the Local Algebras Associated with the Free Massless Scalar Field Theory*, Commun. Math. Phys. **84**, 71 (1982).
- [Hu72] N.M. Hugenholtz in *Mathematics of Contemporary Physics*, R.F. Streater Editor, Academic Press, London, (1972)
- [Ki96] S. Kichenassamy, *Nonlinear Wave Equations*, Marcel Dekker, Inc., New York (1996).
- [KW91] B. S. Kay and R. M. Wald, *Theorems on the uniqueness and thermal properties of stationary, nonsingular, quasifree states on space-times with a bifurcate Killing horizon*, Phys. Rept. **207** (1991), 49.
- [Lle09] F. Lledó, *Modular theory by example*, [arXiv:0901.1004]
- [MLR10] P. Martinetti, R. Longo and H. Rehren, *Geometric modular action for disjoint intervals and boundary conformal field theory*, Rev. Math. Phys. **22**, (2010), 331-354.

- [Mo06] V. Moretti, *Uniqueness theorem for BMS-invariant states of scalar QFT on the null boundary of asymptotically flat spacetimes and bulk-boundary observable algebra correspondence*, Commun. Math. Phys. **268** (2006), 727 [arXiv:gr-qc/0512049].
- [Mo08] V. Moretti, *Quantum out-states states holographically induced by asymptotic flatness: Invariance under spacetime symmetries, energy positivity and Hadamard property*, Commun. Math. Phys. **279** (2008), 31 [arXiv:gr-qc/0610143].
- [RS80] M. Reed, B. Simon, *Methods of Modern Mathematical Physics. I Functional Analysis*. Revised and enlarged edition, Academic Press, San Diego, (1980)
- [RS80II] M. Reed, B. Simon, *Methods of Modern Mathematical Physics. II Fourier Analysis and Self-Adjointness of operators*, Revised and enlarged edition, Academic Press, San Diego, (1980)
- [Sa06] T. Saffary, *On the generator of massive modular groups*, Lett. Math. Phys. **77**, (2006) 235-248.
- [Sa05] T. Saffary, *Modular action on the massive algebras*, Ph.D. Thesis, Universität Hamburg, (2005), <http://unith.desy.de/research/aqft/doctoral-theses/>
- [Sch03] B. Schroer, *Lightfront holography and area density of entropy associated with localization on wedge horizons*, Int. J. Mod. Phys. A **18** (2003) 16711696.
- [Sch08] B. Schroer, *Area density of localization entropy: I. The case of wedge localization*, Class. Quant. Grav. **23** (2006) 52275248, and Addendum *ibid*, **24** (2007) 42394249.
- [SW00] B. Schroer and H. W. Wiesbrock, *Modular theory and geometry*, Rev. Math. Phys. **12** (2000) 139-158
- [Se82] G. L. Sewell, *Quantum field on manifolds: PCT and gravitationally induced thermal states*, Ann. Phys. **141** (1982) 201.
- [Sc08] M. Schottenloher, *A Mathematical Introduction to Conformal Field Theory*, second edition Springer, (2008).
- [So08] C. D. Sogge, *Lectures on Non-Linear Wave Equations*, second Edition International Press, Boston, (2008).
- [Su05] S. J. Summers, *Tomita-Takesaki modular theory*, arXiv:math-ph/0511034.
- [Ta01] M. Takesaki, *Theory of Operator Algebras II*, Encyclopedia of Mathematical Sciences, Vol. 125, Springer, New-York, (2001).

- [Ta96] M. E. Taylor, *Partial Differential Equations*, Applied Mathematical Sciences, vol. 116, Springer, New York, (1996).
- [Vl79] V. S. Vladimirov, *Generalized Functions in Mathematical Physics*, MIR, Moscow, (1979).
- [Wa84] R. M. Wald, *General Relativity*, Chicago University Press, (1984).
- [Wa94] R. M. Wald, *Quantum Field Theory in Curved Space-Time and Black Hole Thermodynamics*, The University of Chicago Press (1994).
- [Wi93] H. W. Wiesbrock, *Half-sided modular inclusions of von Neumann algebras*, Commun. Math. Phys. **157**, (1993) 8392
- [Wi93b] H. W. Wiesbrock, *Half-sided modular inclusions of von Neumann algebras*, Erratum, Commun. Math. Phys. **184**, (1997) 683685
- [Wi97] H. W. Wiesbrock, *Symmetries and modular intersections of von Neumann algebras*, Lett. Math. Phys. **39**, (1997) 203212