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Synthesis of Wide-Angle Scanning Arrays Through Array Power Control

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ABSTRACT A new methodology for the synthesis of wide-angle scanning arrays is proposed. It is based on the formulation of the array design problem as a multi-objective one where, for each scan angle, both the radiated power density in the scan direction and the total reflected power are accounted for. A set of numerical results from full-wave simulated examples - dealing with different radiators, arrangements, frequencies, and number of elements - is reported to show the features of the proposed approach as well as to assess its potentialities. A widening of the field-of-view percentage of up to 22% with respect to standard scanning methods is demonstrated.

INDEX TERMS Wide-angle scanning, excitation design, array synthesis, active reflection coefficient stabilization, multi-objective evolutionary optimization.

I. INTRODUCTION AND MOTIVATION

WIDE-ANGLE scanning is one of the key requirements in modern antenna arrays for communications and sensing [1], [2]. Indeed, stable radiation performance over a wide range of steering angles are increasingly important in several applicative scenarios such as wireless coverage, automotive radars, satellite communications, and imagers [1], [2], [3]. Unfortunately, the gain/directivity of antenna arrays typically worsens as the scan angle increases, specifically in the presence of “blind spots” [1], [2], and it limits the array field-of-view (FoV) [1], [2]. These undesired phenomena are a direct consequence of the instability of the active return loss (ARL) versus the scan angle [4]. In the last decades, several approaches have been introduced in the state-of-the-art (SoA) literature to improve the ARL stability and to increase the

FoV of phased arrays by mitigating the inter-element coupling [3], [5], [6], [7], [8], [9], [10], [11], [12], [13]. More in detail, the exploitation of ad-hoc solutions such as the use of cavities [5], [6] or magnetic dipoles [7] has been proposed to individually insulate the array elements. An alternative strategy, not involving complex radiators, resorts to the introduction of decoupling feed networks [8], [9], [14], which synthesis and implementation become very challenging when a wideband working is required [8], [9], [14]. Wide-angle impedance matching (WAIM) layers coating the array and based either on standard [15] or artificial materials [3], [10], [11], [12], [13] has been also adopted to stabilize the ARL. Nonetheless, this latter solution - as well as the previous ones - needs non-trivial updates of the system hardware, as applied to existing architectures, and customized

designs/implementations strongly depend on the adopted antenna technology [3], [10], [11], [12], [13].

Otherwise, the *ARL* stabilization may be yielded without architectural modifications by considering a completely different perspective, as well. In fact, the *ARL* depends not only on the inter-element coupling, but also on the excitations applied to each elementary radiator of the array [4], [16]. While the synthesis of such excitations is typically aimed at controlling the scan direction, the sidelobe profile [17], and/or more advanced quality-of-service indicators [18], [19], there are not theoretical obstacles to even address the stabilization the *ARL* still fulfilling the user-desired pattern features when scanning.

Following this line of reasoning, an innovative methodology for the synthesis of wide-scanning arrays is proposed into the following. Such an approach is based on the formulation of the array synthesis as a multi-objective problem where, for each scan angle, both the radiated power density in the scan direction and the total reflected power are accounted for by including the finite array coupling in the design process (i.e., no periodicity approximation is employed).

Accordingly, the main innovative contributions of this work include (i) for the first time to the best of the authors' knowledge, the synthesis of the array excitations to minimize the reflected power, while keeping suitable radiation performance, for widening the *FoV* in finite arrays, (ii) the full-wave numerical proof for both linear and planar arrays that it is possible to obtain wider scan angles without modifying the antenna architecture, hence implicitly the assessment of the applicability of the proposed concept to upgrade existing arrays, and (iii) the derivation of a general-purpose strategy for wide angle scanning that, unlike *SoA* methods [3], [5], [6], [7], [8], [9], [10], [11], [12], [13], can be adopted regardless of the underlying antenna technology and/or physical implementation details.

The outline of the paper is as follows. The mathematical formulation of the array synthesis problem at hand is formulated in Section II, where the proposed excitation-based approach to yield wide-angle scanning performance is presented, as well. A set of numerical results from full-wave simulated examples is reported to show the features of the proposed strategy as well as to assess its potentialities in different working conditions and also in comparison with state-of-the-art competitive techniques (Section III). Finally, some conclusions are drawn (Section IV).

II. MATHEMATICAL FORMULATION

Consider the scenario in Fig. 1 where an antenna array of N elements, each one described by the L geometrical parameters $\underline{g} = \{g_l; l = 1, \dots, L\}$ and operating at the carrier frequency f_0 , is fed by the set of N incident excitation coefficients $\underline{w}_q^+ \triangleq \{w_{nq}^+; n = 1, \dots, N\}$ to radiate a beam pattern steered at the q -th ($q = 1, \dots, Q$) scan angle (θ_q, φ_q) [4], [16], [20]. The corresponding reflected

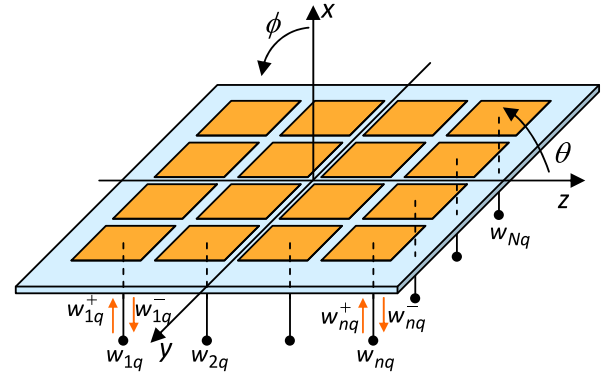


FIGURE 1. Problem geometry. Sketch of the reference phased array layout.

excitation coefficients, measured at the n -th ($n = 1, \dots, N$) element, of the array are defined as [16]

$$w_{nq}^- \triangleq w_{nq}^+ \times \Gamma_{nq} \quad (1)$$

where

$$\Gamma_{nq} \triangleq \sum_{m=1}^N S_{nm} \frac{w_{mq}^+}{w_{nq}^+} \quad (2)$$

it the n -th ($n = 1, \dots, N$) element active reflection coefficient at the q -th ($q = 1, \dots, Q$) scan angle, the corresponding *ARL* being $ARL_{nq} \triangleq |\Gamma_{nq}|^2$ ($n = 1, \dots, N$; $q = 1, \dots, Q$) [3], [4], [10], [11], [12], [13], and S_{nm} is the (n, m) -th ($n, m = 1, \dots, N$) entry of the $N \times N$ scattering matrix S [4], [16], [17], [20]. This latter univocally models the array inter-element coupling effects that, unlike Γ_{nq} (2) and w_{nq}^- (1), does not depend on the applied array excitations w_{nq}^+ .

Under the above assumptions, the total excitation at the n -th ($n = 1, \dots, N$) array element, w_{nq} ($w_{nq} \triangleq w_{nq}^+ + w_{nq}^-$) turns out to be [16], [20]

$$w_{nq} = (1 + \Gamma_{nq})w_{nq}^+, \quad (3)$$

then field radiated by the N -elements antenna array in the Fraunhofer region has the following expression [16], [17], [20]

$$\mathbf{E}(\mathbf{r}|\underline{w}_q) = \frac{jk_0 \exp(-jk_0 r)}{2\pi r} \sum_{n=1}^N w_{nq} \exp[-jk_0(\mathbf{r}_n \cdot \hat{\mathbf{r}})] \mathbf{E}_n(\theta, \varphi), \quad (4)$$

which holds true regardless of the array geometry (linear/planar) and element polarization (i.e., single/dual polarization). In (4), k_0 is the free-space wavenumber, $\mathbf{r} = \{r, \theta, \varphi\}$ is the spatial coordinate ($\hat{\mathbf{r}} = \frac{\mathbf{r}}{r}$), $\cdot$ stands for the scalar product, and $\mathbf{r}_n$ is the location of the n -th ($n = 1, \dots, N$) array element, which element pattern is $\mathbf{E}_n(\theta, \varphi)$ [17]. It is worth pointing out that in finite arrays, the element pattern is different for each n -th ($n = 1, \dots, N$) element [i.e., $\mathbf{E}_n(\theta, \varphi) \neq \mathbf{E}_m(\theta, \varphi)$, $n \neq m$ ($n, m =$

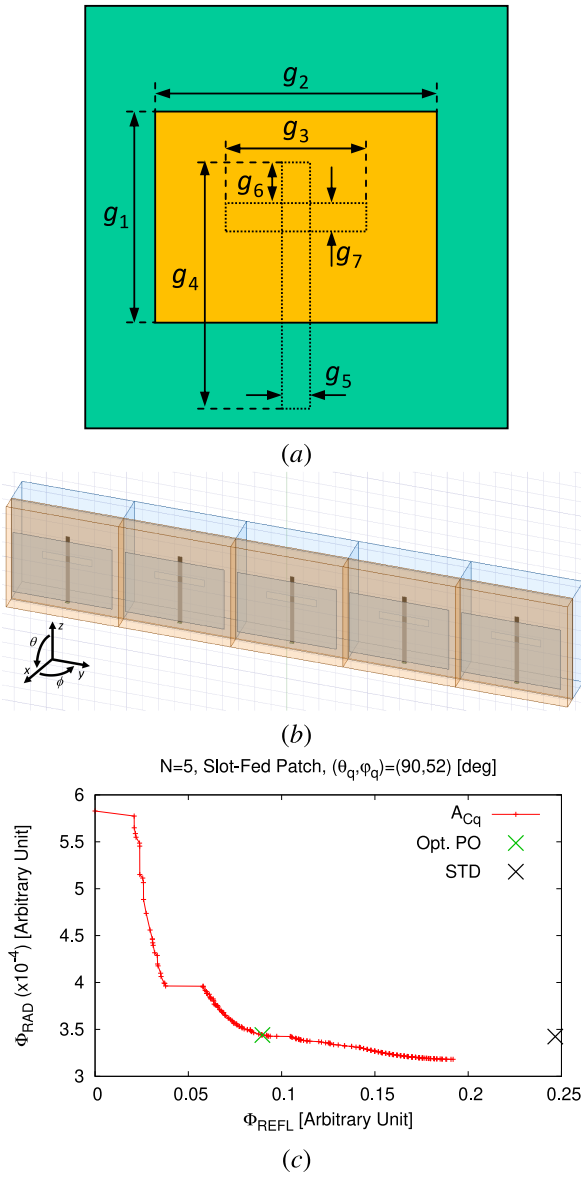


FIGURE 2. Slot-fed patch linear array ($N = 5$, $f = 2$ [GHz], $d = 0.47\lambda$). - Sketch of (a) the geometry of array element and (b) the HFSS model of the array, and plot of the Pareto-optimal solutions when $(\theta_q, \varphi_q) = (90, 52)$ [deg].

$1, \dots, N$], even though the centrally-located radiators of large arrays exhibit very similar angular responses [17].

The power density radiated along the direction (θ, φ) by the array steered towards (θ_q, φ_q) is obtained from (4) as follows

$$\psi_q(\theta, \varphi) = \lim_{r \rightarrow \infty} \frac{r^2}{2\nu} \left| \mathbf{E}(\mathbf{r} | \underline{w}_q) \right|^2 \quad (5)$$

and the total radiated power is $\mathbb{P}_q^{RAD} = \int_{\Omega} \psi_q(\theta, \varphi) d\Omega$, Ω being a far-field closed surface surrounding the array. Moreover, the total reflected power percentage is defined as

$$\zeta_q = \frac{\mathbb{P}_q^{REFL}}{\mathbb{P}_q^{IN}} \quad (6)$$

where $\mathbb{P}_q^{IN}$ and $\mathbb{P}_q^{REFL}$ are the input power ($\mathbb{P}_q^{IN} \triangleq \sum_{n=1}^N |w_{nq}^+|^2$) and the reflected power ($\mathbb{P}_q^{REFL} \triangleq \sum_{n=1}^N |w_{nq}^-|^2$) equal to $\mathbb{P}_q^{REFL} \triangleq \sum_{n=1}^N ARL_{nq} |w_{nq}^+|^2$ when the array is steered along (θ_q, φ_q) , respectively. According to these definitions, the *FoV* of an antenna array can be identified as the angular range Ω_{FoV} ($\Omega_{FoV} \subset \Omega_{TOT}$, $\Omega_{TOT} = \{(\theta_q, \varphi_q); q = 1, \dots, Q\}$) where the radiated power density in the scan direction (θ_q, φ_q) is greater than a threshold ψ_{TH} (i.e., $\psi_q(\theta_q, \varphi_q) \geq \psi_{TH}$) and the total reflected power percentage is smaller than a target value ζ_{TH} (i.e., $\zeta_q \leq \zeta_{TH}$). Mathematically, it turns out that $(\theta_q, \varphi_q) \in \Omega_{FoV}$ if

$$\begin{aligned} \psi_q(\theta_q, \varphi_q) &\geq \psi_{TH} \\ \zeta_q &\leq \zeta_{TH}. \end{aligned} \quad (7)$$

In order to achieve wide scan performance (i.e., a wide *FoV*), *SoA* methods in the reference literature [3], [5], [6], [7], [8], [9], [10], [11], [12], [13] require an ARL_{nq} almost independent on the angular direction (i.e., $ARL_{nq} \approx ARL_n$) and smaller than an application-defined threshold ARL_{TH} ($ARL_n \leq ARL_{TH}$) by enforcing the condition $|S_{nm}| \rightarrow 0$ ($n = 1, \dots, N; m = 1, \dots, M; m \neq n$) to yield

$$\Gamma_n(\theta_q, \varphi_q) \approx S_{nn}. \quad (8)$$

Therefore, the condition $ARL_n \leq ARL_{TH}$ is fulfilled by properly setting the descriptors of the elementary radiator, $\underline{g}$, so that $|S_{nn}|^2 \leq ARL_{TH}$, while the incident excitation coefficients $\underline{w}_q^+$ are optimized to fulfill the requirements on the radiated power density along (θ_q, φ_q) , $\psi_q(\theta_q, \varphi_q)$.

Otherwise, the proposed approach leverages the observation that, according to (2), for a given matrix $\mathcal{S}$ (i.e., resulting from the definition of the element geometries, physical properties, and positions), both the total reflected power percentage, ζ_q , and the radiation properties of the array $\mathbf{E}(\mathbf{r} | \underline{w}_q)$ (4) can be controlled by properly optimizing the excitations $\underline{w}_q^+$ ($q = 1, \dots, Q$) to maximize the *FoV* of the antenna array. Therefore, the design of the optimal excitations for the q -th ($q = 1, \dots, Q$) scan angle, $\underline{w}_q^+|_{opt}$, is mathematically formulated as a multi-objective constrained optimization where $\underline{w}_q^+|_{opt}$ is the set of excitations $\underline{w}_q^+$ that fulfils some feasibility conditions (i.e., $\underline{w}^+ \in \mathcal{W}$, $\mathcal{W}$ being a set of requirements on the excitation coefficients such as constant magnitude, phase quantization, etc.), while minimizing the $T = 2$ objective functions coding the *FoV* condition (7)

$$\underline{w}_q^+|_{opt} \triangleq \arg \left\{ \min_{\underline{w}_q^+} \left[\Phi_{REFL}(\underline{w}_q^+); \Phi_{RAD}(\underline{w}_q^+) \right] \right\}. \quad (9)$$

In (9), the first objective function Φ_{REFL} coincides with the total reflected power percentage when the array is steered along (θ_q, φ_q)

$$\Phi_{REFL}(\underline{w}_q^+) \triangleq \zeta_q, \quad (10)$$

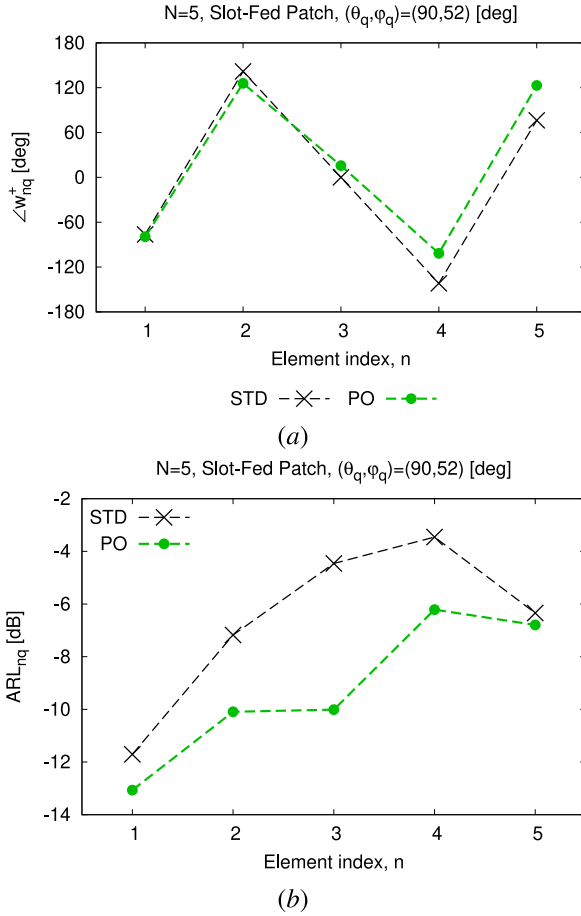


FIGURE 3. Slot-fed patch linear array ($N = 5$, $f = 2$ [GHz], $d = 0.47\lambda$, $q = 53$, $(\theta_q, \phi_q) = (90, 52)$ [deg]) - Plots of (a) $\angle w_{nq}^+$ and (b) ARL_{nq} , $n = 1, \dots, N$.

while the second one, Φ_{RAD} , is defined as the inverse of the power density radiated along the q -th ($q = 1, \dots, Q$) steering direction

$$\Phi_{RAD}(w_q^+) = \frac{1}{\psi_q(\theta_q, \phi_q)}. \quad (11)$$

The synthesis process at hand does not involve the geometrical descriptors of the array, $\underline{g}$, which are user-defined and *a-priori* set.

Since (9) codes a multi-objective optimization problem, which involves highly non-linear cost function terms, and the unknowns (i.e., magnitude/phase of the w^+ entries) are real-valued, a multi-objective steady-state evolutionary algorithm that uses ε -dominance archiving (ε -MOEA) to record a diverse set of Pareto optimal solutions has been adopted [22], [23], [24].

Unlike generational algorithms [25], the ε -MOEA evolves one solution at a time to ensure convergence and diversity throughout the optimization process [22], [23], [24]. For each q -th ($q = 1, \dots, Q$) scan angle, (θ_q, ϕ_q) , the following iterative procedure (c being the iteration index) is performed. At each c -th ($c = 1, \dots, C$) iteration, both a variable-sized Pareto front $\mathcal{A}_{cq}$ of A_{cq} solutions

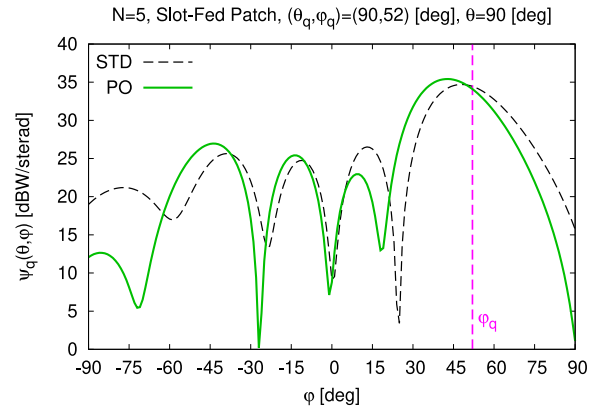


FIGURE 4. Slot-fed patch linear array ($N = 5$, $f = 2$ [GHz], $d = 0.47\lambda$, $q = 53$, $(\theta_q, \phi_q) = (90, 52)$ [deg], $\theta = 90$ [deg]) - Plot of $\psi_q(\theta, \phi)$.

$[\mathcal{A}_{cq} = \left\{ w_q^+ \right\}_c^{(a)} ; a = 1, \dots, A_{cq}]$ and a fixed size population $\mathcal{P}_{cq}$ of P individuals $[\mathcal{P}_{cq} = \left\{ w_q^+ \right\}_c^{(p)} ; p = 1, \dots, P]$ are evolved according to the ε -MOEA mutation, recombination, and selection operators with control parameters $\varepsilon = \{\varepsilon_1, \varepsilon_2\}$ [22], [23], [24]. At the convergence (i.e., $c = C$), a Pareto front of non-dominated excitation sets $\mathcal{A}_{Cq}$ is outputted and the user identifies the optimal one according to a criterion based on either functional (i.e., a set of conditions on Φ_{RAD} and/or Φ_{REFL}) or non-functional (e.g., costs, maintenance, etc...) constraints. Finally, the FoV Ω_{FoV} of the synthesized array (i.e., $\left. w_q^+ \right|_{opt} = \left\{ w_q^+ \right\}_{opt} ; q = 1, \dots, Q_{FoV}$) turns out to be the set of Q_{FoV} ($Q_{FoV} \leq Q$) adjacent scan directions for which (7) holds true.

III. NUMERICAL RESULTS

This section is aimed at both illustrating the features of the proposed method and assessing its effectiveness also in comparison with state-of-the-art solutions.

The first numerical example deals with a linear array of $N = 5$ equally-spaced by $d = 0.47\lambda$ cavity-backed slot-fed microstrip antennas [Fig. 2(a)] printed on a Polyflon CuFlon substrate with thickness 9×10^{-3} [m]. Each radiator [Fig. 2(b)] has been modeled in Ansys HFSS [27] by choosing the following $L = 7$ descriptors to operate at $f = 2$ GHz: $g_1 = 4.13 \times 10^{-2}$ [m], $g_2 = 6.18 \times 10^{-2}$ [m], $g_3 = 3.09 \times 10^{-2}$ [m], $g_4 = 4.58 \times 10^{-2}$ [m], $g_5 = 2.54 \times 10^{-3}$ [m], $g_6 = 1.27 \times 10^{-2}$ [m], and $g_7 = 4.63 \times 10^{-3}$ [m]. Moreover, Ω_{FoV} has been defined by setting in (7) $\zeta_{TH} = 10$ % and choosing ψ_{TH} so that the maximum array scan loss is -6 [dB] (i.e., $\psi_{TH} = -6 + \psi_1(90, 0)$ [dB]). As for the array synthesis method, the requirement that no magnitude tapering is allowed (i.e., $|w_{nq}^+| = 1$, $n = 1, \dots, N$ and $q = 1, \dots, Q$), so that only the phases of the excitations are optimized (PO), has been encoded into the feasibility set $\mathcal{W}$ and the criterion for selecting the optimal trade-off array

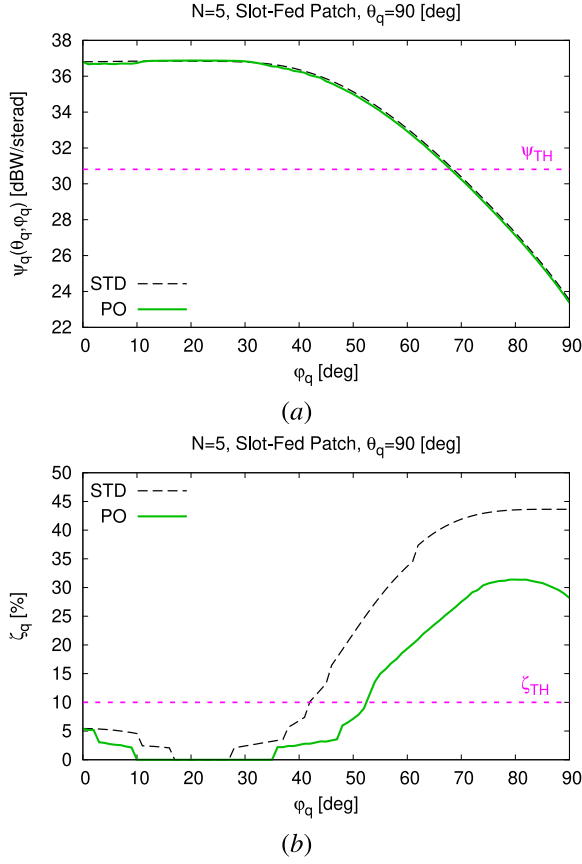


FIGURE 5. Slot-fed patch linear array ($N = 5$, $f = 2$ [GHz], $d = 0.47\lambda$, $\theta_q = 90$ [deg]) - Plots of (a) $\psi_q(\theta_q, \varphi_q)$ and (b) ζ_q versus φ_q ($q = 1, \dots, Q$; $Q = 91$).

(i.e., $\underline{W}^+_{opt}$) has been defined by setting $\Phi_{RAD}^{PO} \equiv \Phi_{RAD}^{STD}$, Φ_{RAD}^{STD} being (11) of the standard approach (STD) where the excitation phases are computed with the analytic linear phase shifting rule [17]. Finally, the ε -MOEA optimizer has been configured according to reference guidelines [22], [23], [24] by choosing the following setup: $\{\varepsilon_1, \varepsilon_2\} = \{5 \times 10^{-3}, 2.5 \times 10^{-2}\}$, $P = 50$, and $C = 1000$.

To illustrate the features and the behavior of the proposed method for a generic q -th ($q = 1, \dots, Q$) scan angle, Fig. 2(c) shows the Pareto front $\mathcal{A}_{Cq}$ when $(\theta_q, \varphi_q) = (90, 52)$ [deg] ($\rightarrow q = 53$). As it can be observed the PO Pareto front¹ dominates the STD solution and the optimal trade-off array [i.e., the “green” cross in Fig. 2(c)] reduces the total reflected power percentage from $\zeta_q^{STD} = 24.6$ % (i.e., $\Phi_{REFL}^{STD} = 0.246$) down to $\zeta_q^{PO} = 8.9$ % (i.e., $\Phi_{REFL}^{PO} = 0.089$), while keeping the same radiation performance (i.e., $\psi_q^{PO}(\theta_q, \varphi_q) \Big|_{q=53} = \psi_q^{STD}(\theta_q, \varphi_q) \Big|_{q=53} = 34.6$ [dBW/sterad] $\rightarrow \Phi_{RAD} = 3.43 \times 10^{-4}$). The optimized excitations, which are reported in Fig. 3(a), are very similar,

¹The overall Pareto front $\mathcal{A}_{Cq}$ shows that, for a given array physical architecture, better radiation performance corresponds to worse ARL and vice versa (i.e., a low total reflected power does not correspond to a high power density in the target steering direction, in general, even though it may yield to an increased total radiated power).

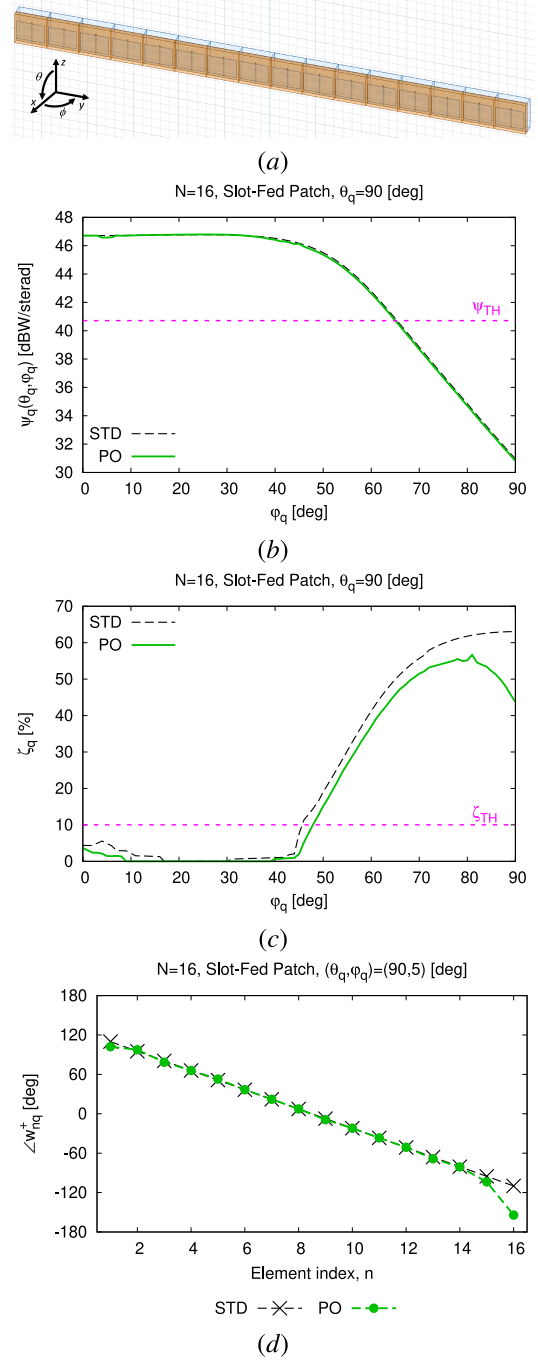


FIGURE 6. Slot-fed patch linear array ($N = 16$, $f = 2$ [GHz], $d = 0.47\lambda$, $\theta_q = 90$ [deg]) - Sketch of (a) the HFSS model of the array, plots of (b) $\psi_q(\theta_q, \varphi_q)$, and (c) ζ_q versus φ_q ($q = 1, \dots, Q$; $Q = 91$), and (d) behaviour of $\angle w_{nq}$ when $q = 6$, $(\theta_q, \varphi_q) = (90, 5)$ [deg].

which implies an analogous architectural complexity, while the corresponding ARL values are mitigated with respect to the STD approach [$ARL_{nq}^{PO} < ARL_{nq}^{STD}$, $n = 1, \dots, N$ - Fig. 3(b)]. This latter outcome confirms that it is possible to indirectly control the ARL of a finite array by minimizing the total reflected power without reducing the radiated power along the steering direction (i.e., dashed magenta line in Fig. 4), although slightly modifying the sidelobe profile and

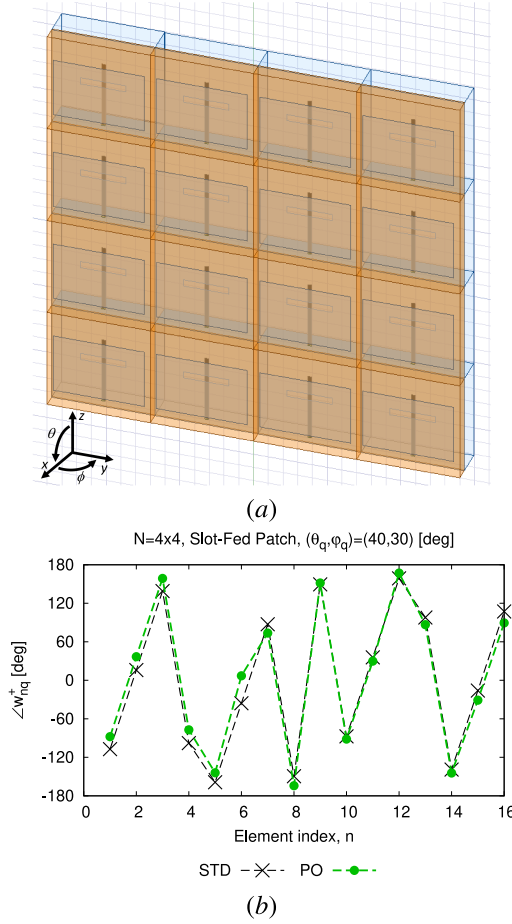


FIGURE 7. Slot-fed patch planar array ($N = 4 \times 4$, $f = 2$ [GHz], $d_x = d_y = 0.47\lambda$) - Sketch of the HFSS model of the array (a) and behaviour of $\angle w_{nq}^+$ when $q = 211$, $(\theta_q, \varphi_q) = (40, 30)$ [deg] (b).

main beam peak location (Fig. 4). As for the computation time and resources for optimizing such a layout, less than 3 minutes were required exploiting a non-optimized MATLAB implementation of the method running on a single-core 1.60 GHz CPU-laptop using less than 100 MByte of RAM. Moreover, it is worth remarking that such resources scale linearly with the array size, hence turning out efficient even for large arrangements.

The resume of the *PO* performance for the first test case is given in Fig. 5 where the plots of the radiated power density, $\psi_q(\theta_q, \varphi_q)$ ($q = 1, \dots, Q$; $Q = 91$) [Fig. 5(a)], and the total reflected power percentage, ζ_q ($q = 1, \dots, Q$; $Q = 91$) [Fig. 5(b)], are reported. The scan width of the finite array is enlarged of about 22 [deg] (i.e., -42 [deg] $\leq \Omega_{FoV}^{STD} \leq 42$ [deg] vs. -53 [deg] $\leq \Omega_{FoV}^{PO} \leq 53$ [deg]), which corresponds to an increment of the field of view percentage α_{FoV} ($\alpha_{FoV} \triangleq \frac{Q_{FoV}}{Q}$) from $\alpha_{FoV}^{STD} \approx 46$ % up to $\alpha_{FoV}^{PO} \approx 58$ % ($\rightarrow \Delta\alpha_{FoV} \approx 12$ %) and a maximum reduction of the reflected power within Ω_{FoV}^{PO} that amounts to $\Delta\zeta_{max}^{STD} = 15.7$ % [$\Delta\zeta_{max}^{STD} \triangleq \max_{q=1, \dots, Q_{FoV}} (\Delta\zeta_q^{STD})$ being $\Delta\zeta_q^{STD} \triangleq \zeta_q^{STD} - \zeta_q^{PO}$].

The applicability of the considered approach to wider linear arrays is assessed next by considering a larger

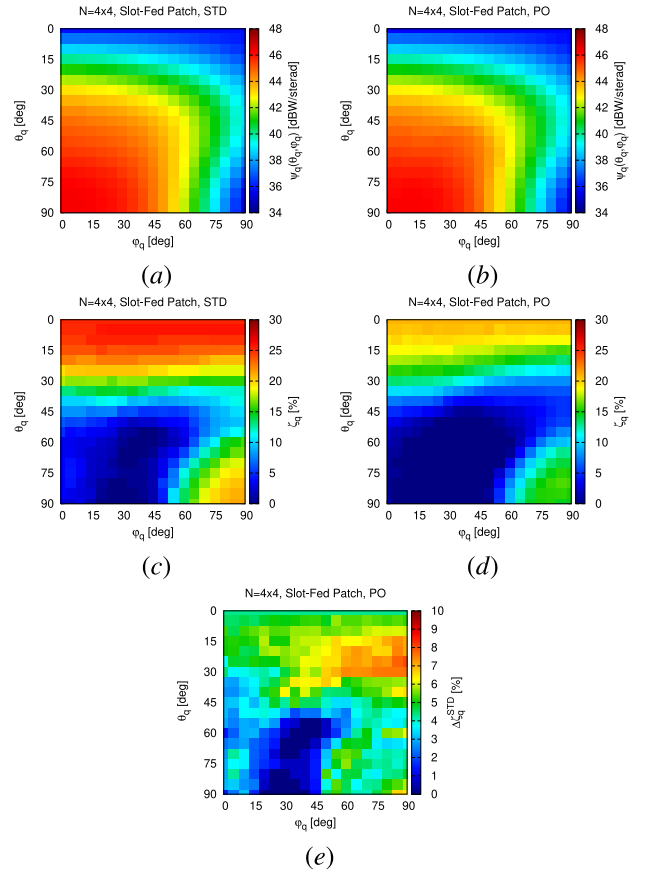


FIGURE 8. Slot-fed patch planar array ($N = 4 \times 4$, $f = 2$ [GHz], $d_x = d_y = 0.47\lambda$) - Plots of (a)(b) $\psi_q(\theta_q, \varphi_q)$, (c)(d) ζ_q , and (e) $\Delta\zeta_q^{STD}$ ($q = 1, \dots, Q$; $Q = 361$).

arrangement with $N = 16$ elements [Fig. 6(a)]. While there is still a widening of the *FoV* [Figs. 6(b)-6(c)], the improvement with respect to the *STD* layout is here smaller ($\Delta\alpha_{FoV} \approx 4$ % $\rightarrow \Delta\Omega_{FoV} \approx 8$ [deg]) since the efficiency of the proposed method in balancing the input power among the array elements is vanifed by the uniform coupling among the majority of them. This result suggests that (i) the margin of improvement enabled by the proposed method is influenced by the number and arrangement of the antenna elements, and that (ii) the mutual interactions in larger arrangements mitigate the achievable advantages. The plots of the optimized excitations of a representative $N = 16$ element solution (i.e., $q = 6$, $(\theta_q, \varphi_q) = (90, 5)$ [deg]) is provided in Fig. 6(d). for completeness.

The third test case is concerned with planar architectures. Towards this end, a planar $N = 4 \times 4$ arrangement of radiators as in Fig. 2(a) has been considered [Fig. 7(a)] - $d_x = d_y = 0.47\lambda$ and the results from the *STD* and the *PO* syntheses are summarized in Fig. 8 by showing the color maps of $\psi_q(\theta_q, \varphi_q)$ ($q = 1, \dots, Q$; $Q = 361$) [Figs. 8(a)-8(b)] and ζ_q ($q = 1, \dots, Q$; $Q = 361$) [Figs. 8(c)-8(d)], together with the optimized excitations of a representative solution [i.e., $q = 211$, $(\theta_q, \varphi_q) = (40, 30)$ [deg] - Fig. 7(b)]. As expected, Figs. 8(a)-8(b) are almost identical by definition since the

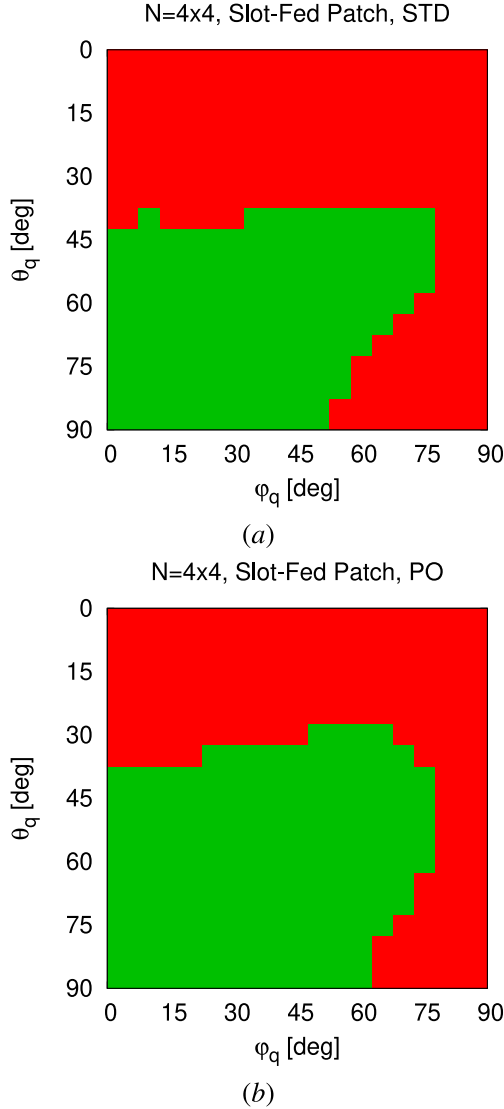


FIGURE 9. Slot-fed patch planar array ($N = 4 \times 4$, $f = 2$ [GHz], $d_x = d_y = 0.47\lambda$) - Plot of Ω_{FoV} (in green) of (a) the STD and (b) the PO arrays.

criterion for selecting the “optimal” trade-off solution forces the condition $\Phi_{RAD}^{PO} \approx \Phi_{RAD}^{STD}$. Otherwise, the advantage of the PO approach is evident when comparing Fig. 8(c) and Fig. 8(d) and it is further highlighted in Fig. 8(e) where the reduction of the total reflected power, $\Delta\zeta_q^{STD}$ ($q = 1, \dots, Q$; $Q = 361$) is reported, $\Delta\zeta_{max}$ being equal to $\Delta\zeta_{max} \approx 8\%$. Therefore, the main outcome (i.e., the scanning performance of the finite array) is that the PO widens the field of view percentage from $\alpha_{FoV}^{STD} \approx 40\%$ [Fig. 9(a)] up to $\alpha_{FoV}^{PO} \approx 49\%$ [Fig. 9(b)] ($\rightarrow \Delta\alpha_{FoV} \approx 9\%$).

The final numerical experiment is aimed at further assessing the PO-based solutions also against recently introduced wide scan layouts [28].²

²Reference [28] has been selected for the comparisons because it demonstrates excellent wide-scan performance and it provides detailed information concerning the array architecture and the optimized excitation coefficients.

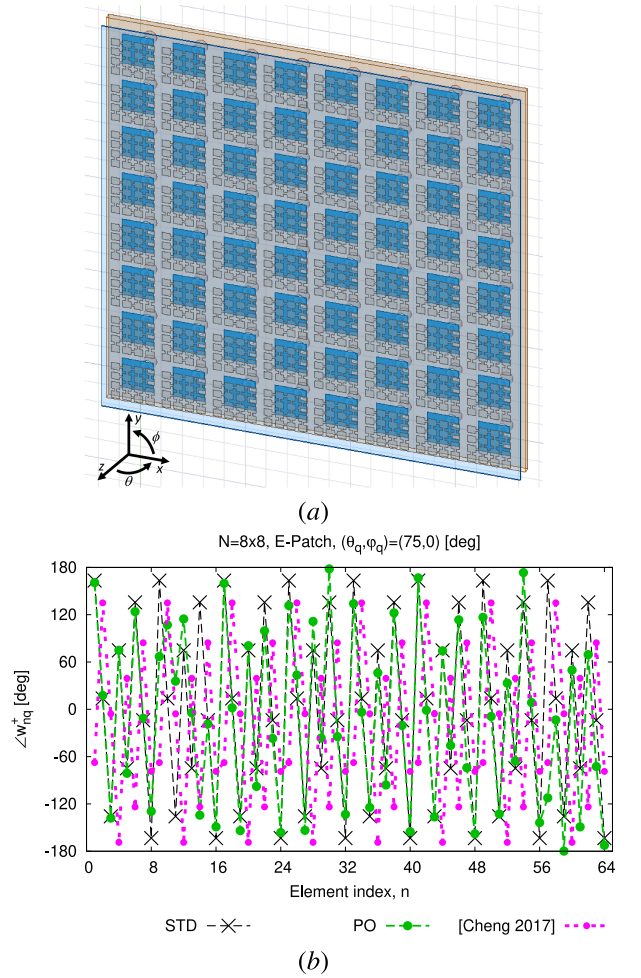


FIGURE 10. Pin-fed E-patch planar array ($N = 8 \times 8$, $f = 5.2$ [GHz], $d_x = 0.42\lambda$, $d_y = 0.43\lambda$) - Sketch of the HFSS model of the array (a) and behaviour of $\angle w_{nq}^+$ when $q = 34$, $(\theta_q, \phi_q) = (75, 0)$ [deg] (b).

Accordingly, a planar arrangement of $N = 8 \times 8$ ($d_x = 0.42\lambda$, $d_y = 0.43\lambda$) identical parasitic pixel layer-based E-shaped radiators [28] operating at $f = 5.2$ GHz has been numerically modeled in Ansys HFSS [Fig. 10(a)].

For comparison purposes, the plots of both the radiated power density and the total reflected power along the cut $\phi_q = 0$ [deg] are shown in Fig. 11 where the values in Table 1 [28] are reported, as well, along with $\angle w_{nq}^+$ of representative solutions [i.e., $q = 34$, $(\theta_q, \phi_q) = (75, 0)$ [deg] - Fig. 10(b)]. Once again the PO array outperforms the STD one yielding on the $\phi_q = 0$ [deg] plane an improvement of the FoV of about $\Delta\alpha_{FoV} \approx 13.2\%$ ($\alpha_{FoV}^{STD} \approx 35.4\%$ vs. $\alpha_{FoV}^{PO} \approx 48.6\%$ [Figs. 11(a)-11(b)]), while it turns out to be $\Delta\alpha_{FoV} \approx 22.6\%$ in the whole angular range [Fig. 12(a) vs. Fig. 12(b)].

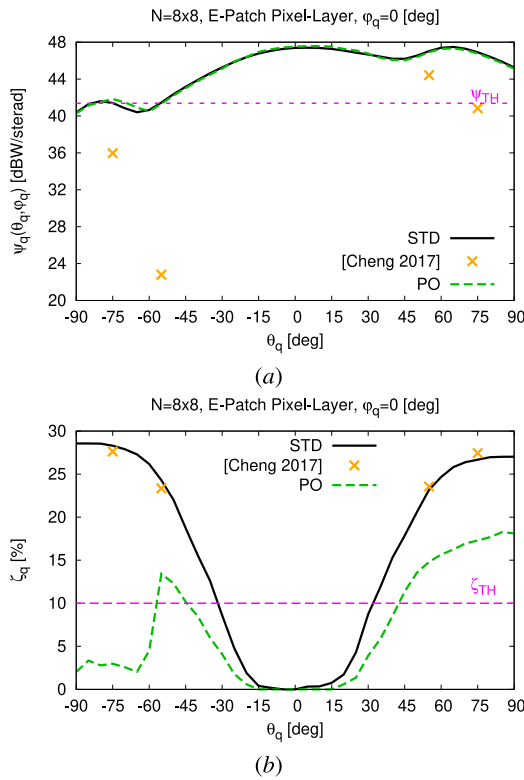
The PO layout is also better than that in [28] in terms of $\psi_q(\theta_q, 0)$ and ζ_q ($q = 1, \dots, Q$; $Q = 37$) at the available angular samples (i.e., $\theta_4 = -75$ [deg], $\theta_8 = -55$ [deg], $\theta_{30} = 55$ [deg], and $\theta_{34} = 75$ [deg]) since 2.49 [dB] $\leq \Delta\psi_q^{[Cheng2017]}(\theta_q, 0) \leq$

TABLE 1. Pin-fed E-patch planar array ($N = 8 \times 8$, $f = 5.2$ [GHz], $d_x = 0.42\lambda$, $d_y = 0.43\lambda$, $\varphi_q = 0$ [deg]) - Performance indexes.

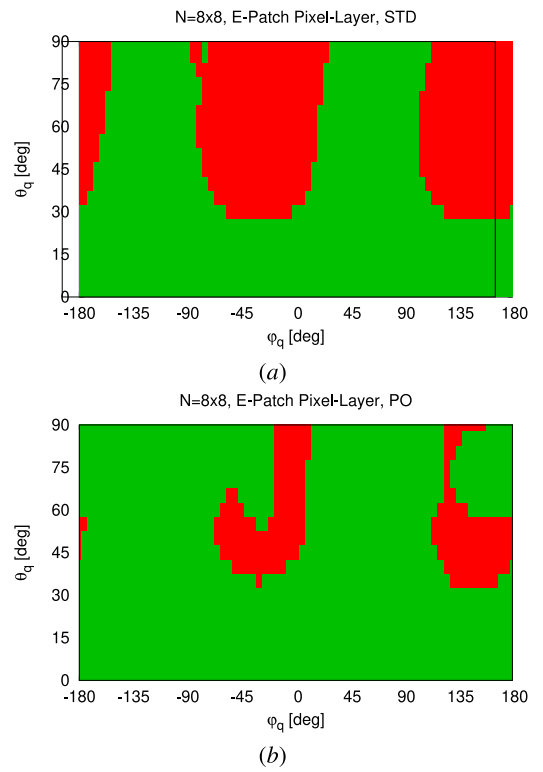
q	(θ_q, φ_q)	$\psi_q^{PO}(\theta_q, \varphi_q)$	$\psi_q^{[Cheng 2017]}(\theta_q, \varphi_q)$	$\psi_q^{STD}(\theta_q, \varphi_q)$	ζ_q^{PO}	$\zeta_q^{[Cheng 2017]}$	ζ_q^{STD}
4	$(-75, 0)$	41.83	35.97	41.40	2.98	27.62	28.30
8	$(-55, 0)$	41.32	22.79	41.45	13.42	23.35	24.30
30	$(55, 0)$	46.91	44.42	47.02	14.80	23.35	23.16
34	$(75, 0)$	46.76	40.83	46.91	17.28	27.43	26.68

TABLE 2. Pin-fed E-patch planar array ($N = 8 \times 8$, $f = 5.2$ [GHz], $d_x = 0.42\lambda$, $d_y = 0.43\lambda$, $\varphi_q = 0$ [deg]) - Performance indexes in [28].

q	(θ_q, φ_q)	SLL_q^{PO}	$SLL_q^{[Cheng 2017]}$	SLL_q^{STD}	G_q^{PO}	$G_q^{[Cheng 2017]}$	G_q^{STD}	$\Delta\Theta_q^{PO}$	$\Delta\Theta_q^{[Cheng 2017]}$	$\Delta\Theta_q^{STD}$
4	$(-75, 0)$	-2.12	-5.78	-5.38	6.57	6.56	7.38	3.55	4.03	3.75
8	$(-55, 0)$	-3.82	-5.68	-5.44	6.54	6.04	7.04	3.71	1.86	2.26
30	$(55, 0)$	-8.83	-10.06	-9.85	12.16	11.78	12.47	2.77	2.17	2.22
34	$(75, 0)$	-9.06	-9.29	-9.61	12.22	11.79	12.62	5.06	0.94	4.94

**FIGURE 11.** Pin-fed E-patch planar array ($N = 8 \times 8$, $f = 5.2$ [GHz], $d_x = 0.42\lambda$, $d_y = 0.43\lambda$, $\varphi_q = 0$ [deg]) - Plots of (a) $\psi_q(\theta_q, \varphi_q)$ and (b) ζ_q versus θ_q ($q = 1, \dots, Q$; $Q = 37$).

18.35 [dB] ($\Delta\psi_q^{[Cheng 2017]}(\theta_q, \varphi_q) \triangleq \psi_q^{PO}(\theta_q, \varphi_q) - \psi_q^{[Cheng 2017]}(\theta_q, \varphi_q)$ [dB]) and $8.8\% \leq \Delta\zeta_q^{[Cheng 2017]} \leq 24.64\%$ in Table 1. For fairness, the comparison has been also carried out by considering the figures-of-merit optimized in [28] and the outcomes are summarized in Table 2 where SLL_q , G_q , and $\Delta\Theta_q$ are the normalized sidelobe level ($SLL_q \triangleq \frac{\max_{(\theta, \varphi) \notin ML} \psi_q(\theta, \varphi)}{\max_{(\theta, \varphi) \in ML} \psi_q(\theta, \varphi)}$, ML being the main-lobe region), the array gain ($G_q \triangleq 4\pi \frac{\max_{(\theta, \varphi)} \psi_q(\theta, \varphi)}{\mathbb{P}_{IN}^q}$), and the scan angle error ($\Delta\Theta_q \triangleq \theta_q - \arg\{\max_{(\theta, \varphi)} \psi_q(\theta, \varphi)\}$), respectively. As it can be noticed (Table 2), the values for the two layouts are quite similar (i.e., -3.66 [dB] $\leq \Delta SLL_q^{[Cheng 2017]} \leq$

**FIGURE 12.** Pin-fed E-patch planar array ($N = 8 \times 8$, $f = 5.2$ [GHz], $d_x = 0.42\lambda$, $d_y = 0.43\lambda$) - Plot of Ω^{FOV} (in green) of (a) the STD and (b) the PO arrays.

-0.23 [dB], 0.01 [dB] $\leq \Delta G_q^{[Cheng 2017]} \leq 0.50$ [dB], and -0.48 [deg] $\leq \Delta\Theta_q^{[Cheng 2017]} \leq 4.12$ [deg]). These results remark that, owing to the nature of the design problem, the proposed trade-off solution exhibit improved performance in terms of radiated power density and total reflected power with respect to the technique in [28] (bold values in Table 1), while yielding sub-optimal, although very close, sidelobe and scan angle performance (e.g., bold values in Table 2).

Finally, Fig. 13 gives the color level maps of the ARL of the $N = 8 \times 8$ array elements when $(\theta_q, \varphi_q) = (-75, 0)$ [deg] (i.e., $q = 4$), ARL_{nq} ($n = 1, \dots, N$), for the PO [Fig. 13(a)], the SoA [28] [Fig. 13(b)], and the STD [Fig. 13(c)] layouts. As it can be observed, generally

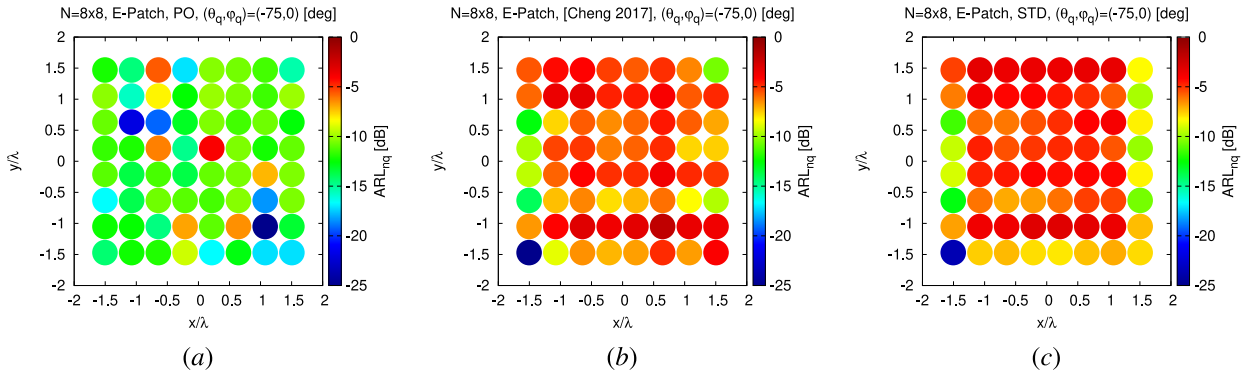


FIGURE 13. Pin-fed E-patch planar array ($N = 8 \times 8$, $f = 5.2$ [GHz], $d_x = 0.42\lambda$, $d_y = 0.43\lambda$, $q = 4$, $(\theta_q, \varphi_q) = (-75, 0)$ [deg]) - Plots of ARL_{nq} , $n = 1, \dots, N$ of (a) PO, (b) SoA [28], and (c) the STD arrays.

$ARL_{nq}^{PO} < ARL_{TH}$ ($n = 1, \dots, N$; $q = 4$), $ARL_{TH} = -10$ [dB]. This outcome further assesses the capability of the PO-based synthesis to indirectly control the ARL of a finite array by minimizing the total reflected power without reducing the radiated power along the steering direction. Moreover, the condition $ARL_{nq}^{PO} < ARL_{nq}^{[Cheng 2017]}$ ($q = 4$) holds true in several n -th ($n = 1, \dots, N$) radiating elements [Fig. 13(a) vs. Fig. 13(b)], while $ARL_{nq}^{[Cheng 2017]} \approx ARL_{nq}^{STD}$ ($n = 1, \dots, N$; $q = 4$).

IV. CONCLUSION AND OBSERVATIONS

The synthesis of the excitations of antenna arrays to achieve wide-angle scanning features (i.e., a wide FoV) with minimum reflected power has been addressed. Towards this end, an innovative methodology has been proposed by formulating the array synthesis as a multi-objective problem where, for each scan angle, both the radiated power density in the scan direction and the total reflected power are accounted for by including the finite array coupling in the design process (i.e., no periodicity approximation is employed). The effectiveness of the approach has been assessed with a selected set of numerical test cases by also proving the possibility to indirectly control the ARL of finite arrays.

The main outcomes from the numerical validation, which has been carried out with a commercial and industry-standard full-wave software [27] to have a faithful modelling of all the electromagnetic interactions and distortions caused by non-idealities, can be summarized as follows:

- the implemented design concept outperforms the standard linear phase shifting technique [17] in yielding wide-angle scanning arrays regardless of the frequency, the radiating elements, and the array arrangement (i.e., linear or planar);
- the PO-based synthesis method allows one to enhance the FoV of finite antenna arrays without modifying the antenna architecture, thus it can be applied to upgrade existing arrays;
- the multi-objective design process provides the user with a multiplicity of trade-off Pareto-front solutions, thus the possibility of choosing an optimal one depending on an application-driven criterion based on either

functional or non-functional (e.g., costs, maintenance, etc. . .) constraints;

- the proposed scheme has non-negligible advantages over recent alternative approaches such as [28].

Future works, beyond the scope of this manuscript, will be aimed at extending the proposed methodology to conformal structures. Moreover, the realization and subsequent measurement of a reconfigurable phased array prototype for the validation of the presented results, currently impossible in our University owing to the lack of suitable resources and facilities, will be addressed in future contributions at least on a small-scale case to verify the accuracy of the simulation results and the practical feasibility of the proposed method.

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